

***SUMMER SCHOOL ON PARTICLE PHYSICS***

**16 June - 4 July 2003**

**NEUTRINO PHYSICS**

**Lectures IV & V**

**S. PARKE  
Fermilab  
Batavia, IL  
U.S.A.**



## References:

Past ICTP Summer School Lectures:

Lisi — 2001 (ICTP web)  
Akmedov — 1999 hep-ph/0001264  
Pakvasa — 1997 hep-ph/9804426

+ references therein

## Recent Reviews

Pakvasa + Valle hep-ph/0301601  
Gonzalez + Garcia + Nir hep-ph/0202058

## MSW + Solar $\nu$ s

Smirnov — hep-ph/0305106

Parker — SLAC Summer School 1986

Fermilab - Conf - 86-131-T

available on SPIRES

click Fermilab-Library-Server:

⚡  
little known

Problem:

2

Kamiland:

Expected:  $87 \pm 6$

Observed: 54

Prob  $< 0.05\%$

$$\frac{\text{Diff:}}{\sqrt{\text{Exp}}} = \frac{87 - 54}{\sqrt{87}} = 3.5 = \frac{33}{9.3}$$

K2K:

Expected:  $80 \pm 6$

Observed: 56

Prob:  $\sim 1.3\%$

$$\frac{\text{Diff:}}{\sqrt{\text{Exp}}} = \frac{80 - 56}{\sqrt{80}} = \frac{24}{8.9} = 2.7$$

Redo with Poisson ?

Statistics

Problem:

3

For 3 flavors:

SHOW

$$2 \sum_{i>j} \text{Im}(U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j}) \sin 2\Delta_{ij}$$

$$= \pm \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \\ \sin \Delta_{21} \sin \Delta_{32} \sin \Delta_{31}$$

where  $\Delta_{ij} = \frac{\delta m_{ij}^2 L}{4E}$

SAME for all  $\alpha \neq \beta$  (if  $\alpha \neq \beta$ )

- or + for  $\mu \rightarrow e, \tau \rightarrow \mu, e \rightarrow \tau$   
+ or - for  $e \rightarrow \mu, \mu \rightarrow \tau, \tau \rightarrow e$

Hint:  $\Delta_{31} - \Delta_{32} - \Delta_{21} \equiv 0$

# $\theta_{13}$ and Beyond

4

Reactors:  $\bar{\nu}_e \rightarrow \bar{\nu}_e$  disappearance  
ala Chooz (KamLAND etc)

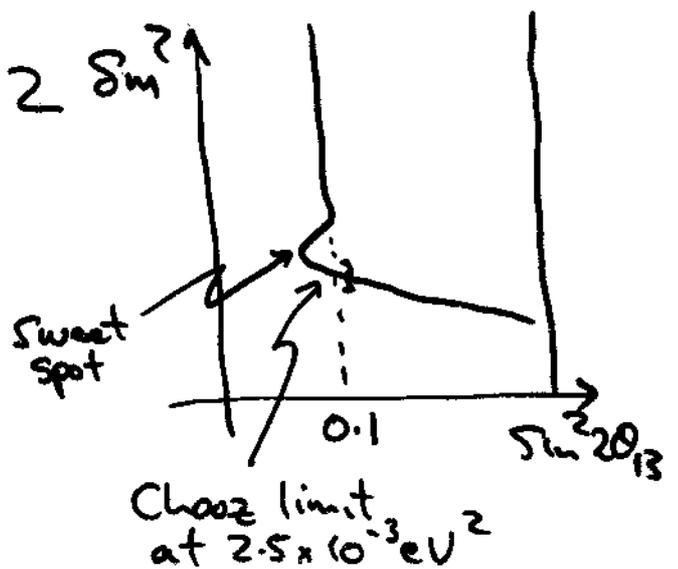
$$P_{\bar{\nu}_e \rightarrow \bar{\nu}_e} \approx 1 - \sin^2 2\theta_{13} \sin^2 \Delta_{31}$$
$$\Delta_{31} = \frac{\delta m_{31}^2 L}{4E}$$

Currently from Chooz

$$\sin^2 2\theta_{13} < 0.05 \text{ at any } \delta m_{31}^2$$

at  $2.5 \times 10^{-3} \text{ eV}^2$ :  $\sin^2 2\theta_{13} < 0.01$   
SK  $\delta m_{mat}^2$

Could get a factor of 2  $\delta m^2$   
by working at a better  
distance



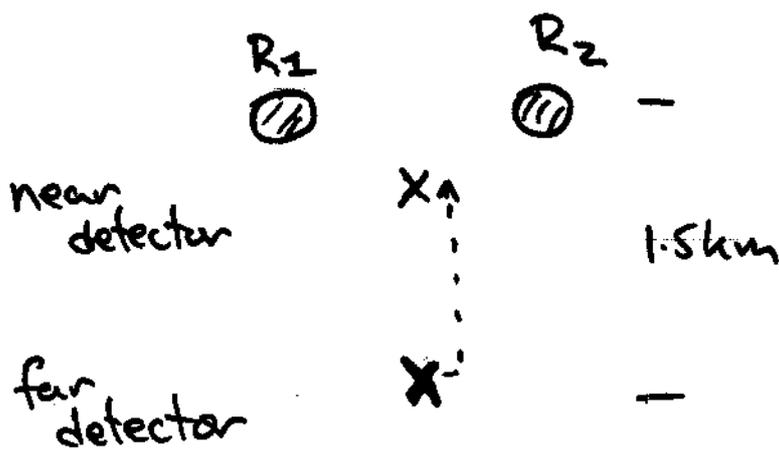
Oscillation Peak (or dip in disappearance) <sup>5</sup>  
occurs at

$$L = \frac{4\pi E}{28m_{atm}^2} = \frac{\pi E}{2(1.27)8m_{atm}^2} = \frac{\pi 3 \times 10^{-3}}{2.54 \times 2.5 \times 10^{-3}} \text{ km}$$

$$\approx 1.5 \text{ km}$$

Also reduce systematic errors

→ 1%



Differential Total Flux measurement

maybe able to get to  $\sin^2 2\theta_{13} \sim 0.01$  Shape  
Lindner et al  
'03

- to do better than Chooz, relatively easy
- to get to 0.01 HARD, DIFFERENT, ...

$$P_{\nu_e \rightarrow \nu_e} = \left| U_{e3} U_{e3}^* e^{-i m_3 L / 2E} + U_{e2} U_{e2}^* e^{-i m_2 L / 2E} + U_{e1} U_{e1}^* e^{-i m_1 L / 2E} \right|^2 \quad 6$$

$$U_{e1} U_{e1}^* = 1 - U_{e2} U_{e2}^* - U_{e3} U_{e3}^* \quad \text{unitarity}$$

$$P_{\nu_e \rightarrow \nu_e} = \left| 1 - 2i s_{13}^2 e^{-i \Delta_{31}} \sin \Delta_{31} - 2i c_{13}^2 s_{12}^2 e^{-i \Delta_{21}} \sin \Delta_{21} \right|^2$$

Square

$$= 1 - \sin^2 2\theta_{13} \left( \sin^2 \Delta_{31} + 2 s_{12}^2 \sin \Delta_{31} \cos \Delta_{32} \sin \Delta_{21} + s_{12}^4 \sin^2 \Delta_{21} \right) - c_{13}^2 \sin^2 2\theta_{12} \sin^2 \Delta_{21}$$

wiggles

At distances, intermediate between (30km) Chooz and KamLAND, there are wiggles which, for  $\sin^2 2\theta_{13}$  not too small, could be used to determine  $\sin^2 \theta_{12}$ ,  $\sin^2 \theta_{13}$ . see Petcov et al June 2003



clinton



Old UCI Page

# The President's View

Announcement *(excerpted from remarks at the MIT commencement, June 6, 1998)*

Paper  
(submitted to  
Phys.Rev.Lett)

Clinton on  
Neutrinos

FAQ

Glossary

Links

[W]e must help you to ensure that America continues to lead the revolution in science and technology. Growth is a prerequisite for opportunity, and scientific research is a basic prerequisite for growth. Just yesterday in Japan, physicists announced a discovery that tiny neutrinos have mass. Now, that may not mean much to most Americans, but it may change our most fundamental theories -- from the nature of the smallest subatomic particles to how the universe itself works, and indeed how it expands.



This discovery was made. in Japan, yes, but it had the support of the investment of the U.S. Department of Energy. This discovery calls into question the decision made in Washington a couple of years ago to disband the Super-conducting Supercollider, and it reaffirms the importance of the work now being done at the Fermi National Acceleration Facility in Illinois.

President Clinton addresses the graduating class at MIT

The larger issue is that these kinds of findings have implications that are not limited to the laboratory. They affect the whole of society -- not only our economy, but our very view of life, our understanding of our relations with others, and our place in time.



The full text is of the President's address is also available.

$\theta_{13}$  and

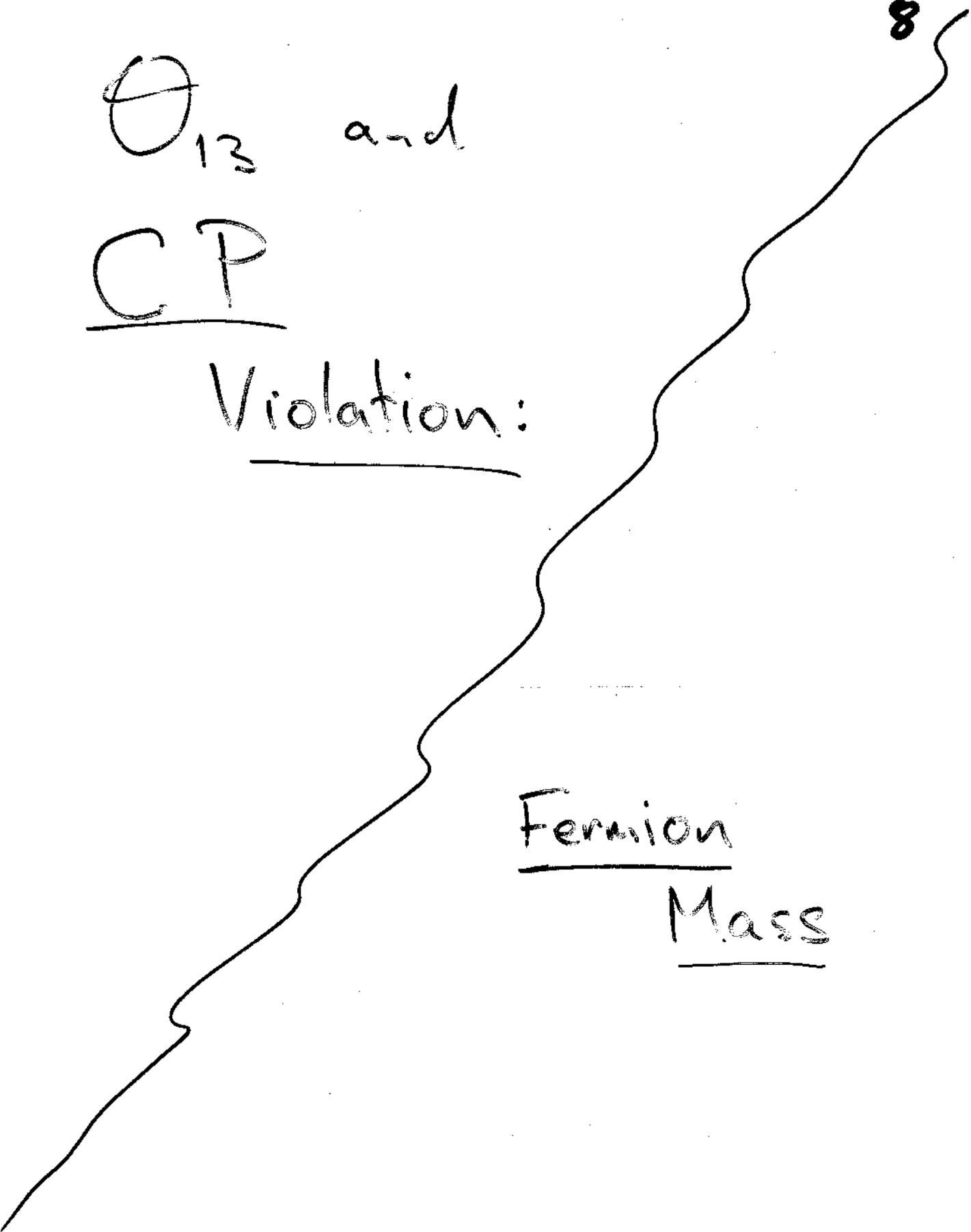
CP

Violation:

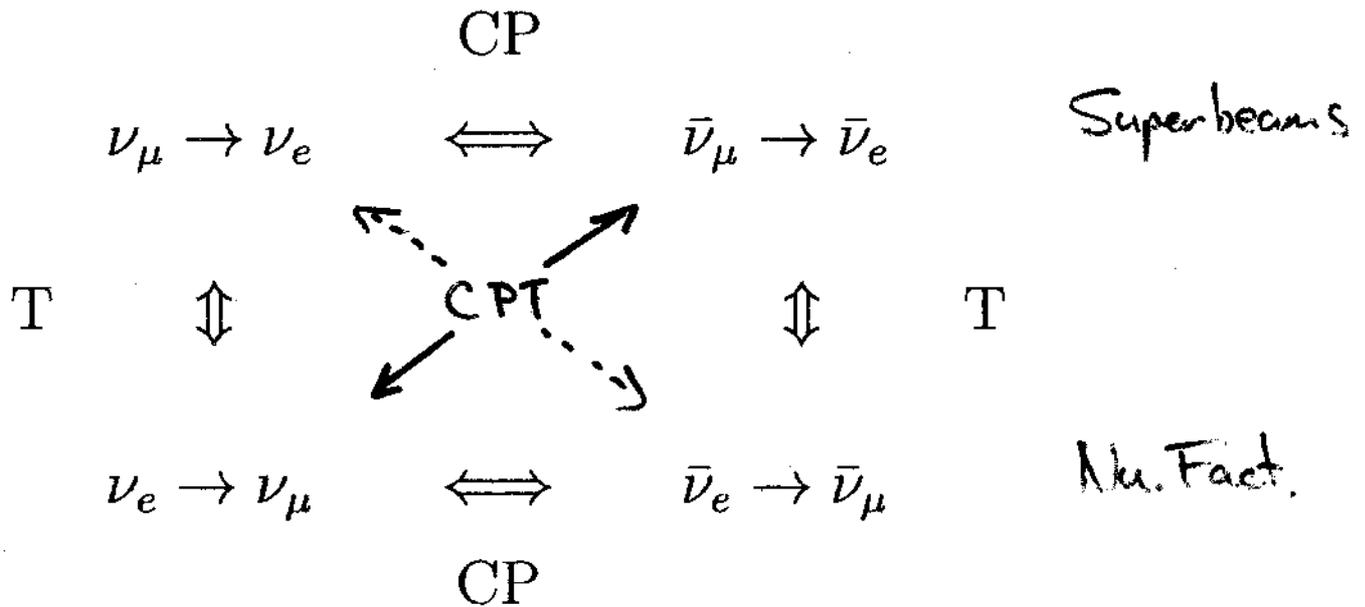
Fermion

Mass

8



## Leptonic CP and T Violation in Oscillations



IN GENERAL (in vacuum):

CP Violation:

$$\alpha \neq \beta \quad P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

T Violation:

$$\alpha \neq \beta \quad P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\nu_\beta \rightarrow \nu_\alpha)$$

and  $P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \neq P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$

CPT Violation:

$$\text{any } \alpha, \beta \quad P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$$

$$\underline{\underline{P_{\nu_{\mu} \rightarrow \nu_e}}}$$

$$P_{\mu \rightarrow e} = \left| U_{e3} U_{\mu 3}^* e^{-i m_3^2 / 2L} \right.$$

$$+ U_{e2} U_{\mu 2}^* e^{-i m_2^2 / 2L}$$

$$\left. + U_{e1} U_{\mu 1}^* e^{-i m_1^2 / 2L} \right|$$

multiple by  $e^{+i m_1^2 / 2L}$

$$\text{then } U_{e1} U_{\mu 1}^* = -U_{e3} U_{\mu 3}^* - U_{e2} U_{\mu 2}^* \quad \Delta_{ij} = \frac{\delta m_{ij}^2 L}{4E}$$

$$= \left| 2U_{e3} U_{\mu 3}^* e^{-i \Delta_{31}} \sin \Delta_{31} \right.$$

$$\left. + 2U_{e2} U_{\mu 2}^* e^{-i \Delta_{21}} \sin \Delta_{21} \right|^2$$

multiple by  $e^{+i \Delta_{32}}$

$$= \left| e^{-i \Delta_{32}} 2U_{e3} U_{\mu 3}^* \sin \Delta_{31} \right.$$

$$\left. + 2U_{e2} U_{\mu 2}^* \sin \Delta_{21} \right|^2$$

$$2U_{e3}U_{\mu 3}^* = 2S_{13}C_{13} \cdot S_{23} e^{-i\delta} \quad ||$$

$$2U_{e2}U_{\mu 2}^* \approx 2C_{13}S_{12}(C_{23}C_{12} - S_{13}S_{12}S_{23}e^{-i\delta})$$

$$P_{\mu \rightarrow e} = \left| e^{-i(\Delta_{32} + \delta)} \sqrt{P_{\text{atm}}} + \sqrt{P_{\odot}} \right|^2$$

with  $P_{\text{atm}} \equiv \left( S_{23} \sin^2 2\theta_{13} \sin \Delta_{31} \right)^2$

$$P_{\odot} = \left( C_{13} C_{23} \sin 2\theta_{12} \sin \Delta_{21} \right)^2$$

Oscillation Max:  $\Delta_{31} \approx \Delta_{32} = \frac{\pi}{2}$

Max: CP Violations:  $\delta = \frac{\pi}{2}$

$$P_{\mu \rightarrow e} = \left| \sqrt{P_{\text{atm}}} \pm \sqrt{P_{\odot}} \right|^2$$

(+  $\nu$ ) (-  $\bar{\nu}$ )

$$\text{Asym: } \frac{P - \bar{P}}{P + \bar{P}} = \frac{2\sqrt{P_{\text{atm}}}\sqrt{P_{\odot}}}{(P_{\text{atm}} + P_{\odot})}$$

# $P(\nu_\mu \rightarrow \nu_e)$

## Why Everybody is Excited!

• Maximum Allowed Asymmetry ( $\delta = \frac{\pi}{2}$  or  $\frac{3\pi}{2}$ ) for  $\nu_\mu \rightarrow \nu_e$  at first Oscillation Maximum in vac:

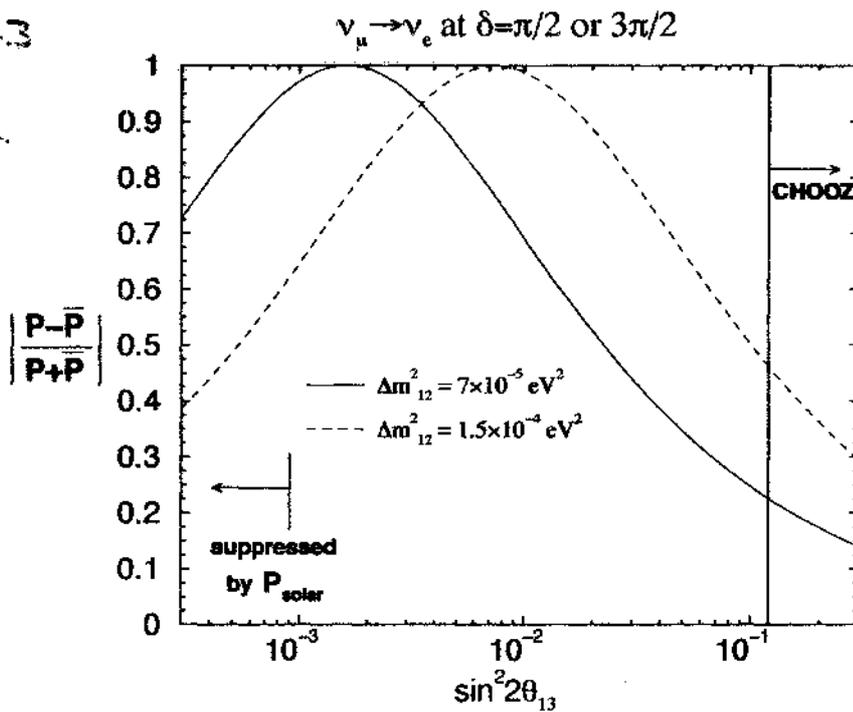
•  $P, \bar{P} = |a_{\mu \rightarrow e}^{atm} + a_{\mu \rightarrow e}^\odot|^2 \approx (\sin \theta_{23} \sin 2\theta_{13} \pm \sqrt{P_\odot})^2$

•  $|P - \bar{P}| \approx 4\sqrt{P_\odot} \sin \theta_{23} \sin 2\theta_{13}$

•  $P + \bar{P} \approx 2 \sin^2 \theta_{23} \sin^2 2\theta_{13} + 2P_\odot$

easily generalized  
 •  $180^\circ \neq \frac{\pi}{2}$   
 • off O.M.  
 • matter

$a_{\mu \rightarrow e}^{atm} \propto U_{e3}$   
 $a_{\mu \rightarrow e}^\odot \propto U_{e2}$



Asymmetry  
 LARGE:  $\checkmark$

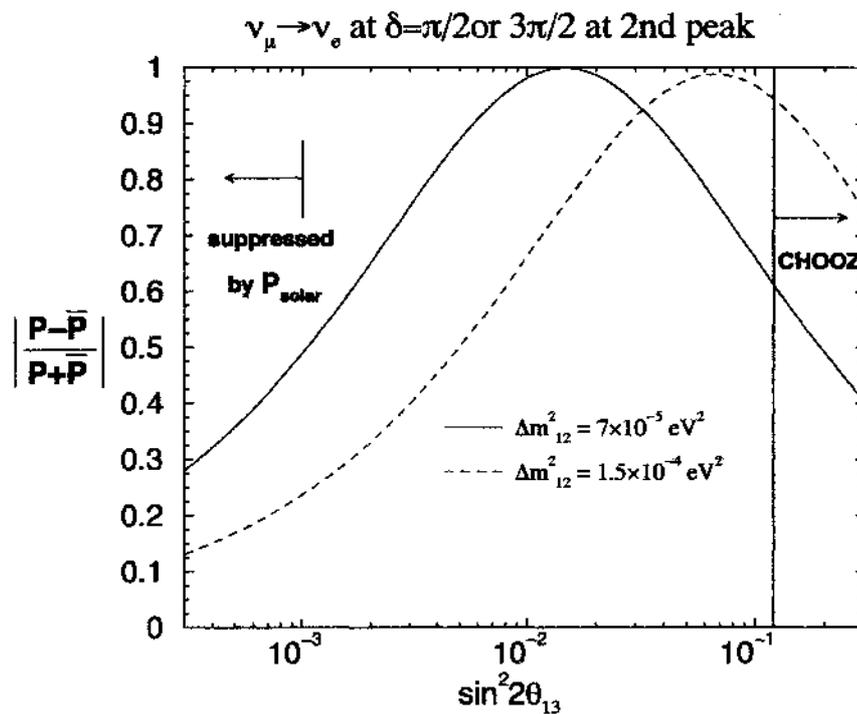
• Peak occurs at

$$\sin^2 2\theta_{13} \approx \frac{\sin^2 2\theta_{12}}{\tan^2 \theta_{23}} \left[ \frac{\pi}{2} \frac{\delta m_{12}^2}{\delta m_{13}^2} \right]^2$$

at OM  $\sqrt{P_\odot} = \cos \theta_{13} \cos \theta_{23} \sin 2\theta_{12} \sin \left( \frac{\pi}{2} \frac{\delta m_{12}^2}{\delta m_{13}^2} \right)$

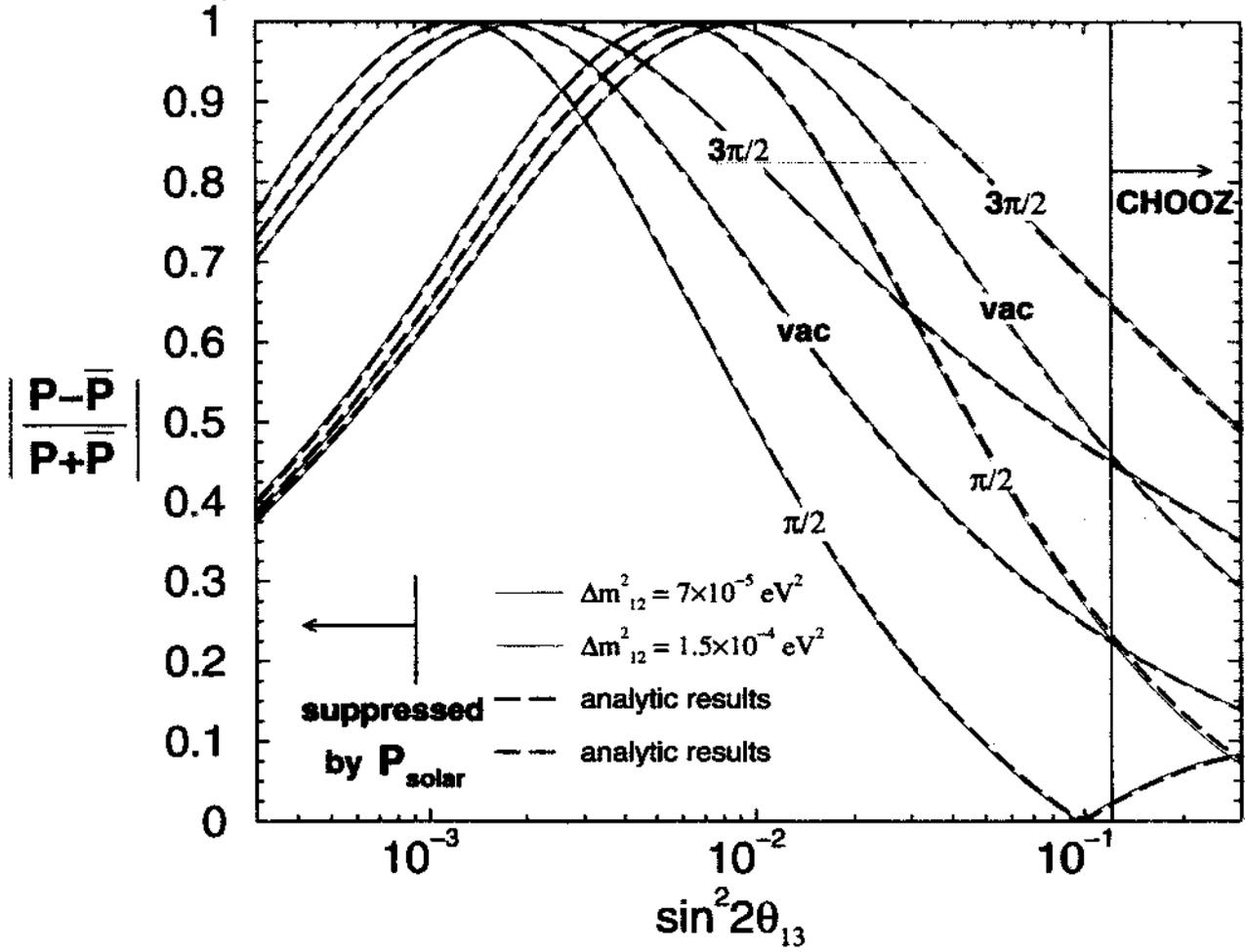
• For BK

## 2nd Peak



- Peak occurs at  $\sin^2 2\theta_{13} \approx \frac{\sin^2 2\theta_{12}}{\tan^2 \theta_{23}} \left[ \frac{3\pi}{2} \frac{\delta m^2_{12}}{\delta m^2_{13}} \right]^2$

$\nu_\mu \rightarrow \nu_e$  at  $\delta = \pi/2$  or  $3\pi/2$  in matter for  $L = 732$  km



$$P_{\mu \rightarrow e} = P_{\text{atm}} + 2\sqrt{P_{\text{atm}}}\sqrt{P_{\odot}} \cos(\Delta_{32} \pm \delta) + P_{\odot} \quad 15$$

$$\cos(\Delta_{32} + \delta) = \cos \Delta_{32} \cos \delta + \frac{\sin \Delta_{32} \sin \delta}{\leftarrow}$$

CP violating term is

$$\sin \delta \cdot \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13}$$

$$* \sin \Delta_{21} \sin \Delta_{32} \sin \Delta_{31}$$

# FULL EXPRESSION:

16

$$P_{\mu \rightarrow e} = P_{\otimes} - 2 S_{23}^2 S_{2(13)}^2 S_{12}^2 \sin \Delta_{31} \cos \Delta_{32} \underbrace{\sin \Delta_{21}}_{\text{small}} \\ + P_{\odot} (1 - 2 S_{13} \tan \theta_{23} \tan \theta_{12} \cos \delta + S_{13}^2 \tan^2 \theta_{23} \tan^2 \theta_{12}) \\ + 2 \sqrt{P_{\otimes} P_{\odot}} \cos (\Delta_{32} \mp \delta)$$

$$P_{\otimes} = S_{23}^2 S_{2(13)}^2 \sin^2 \Delta_{31}$$

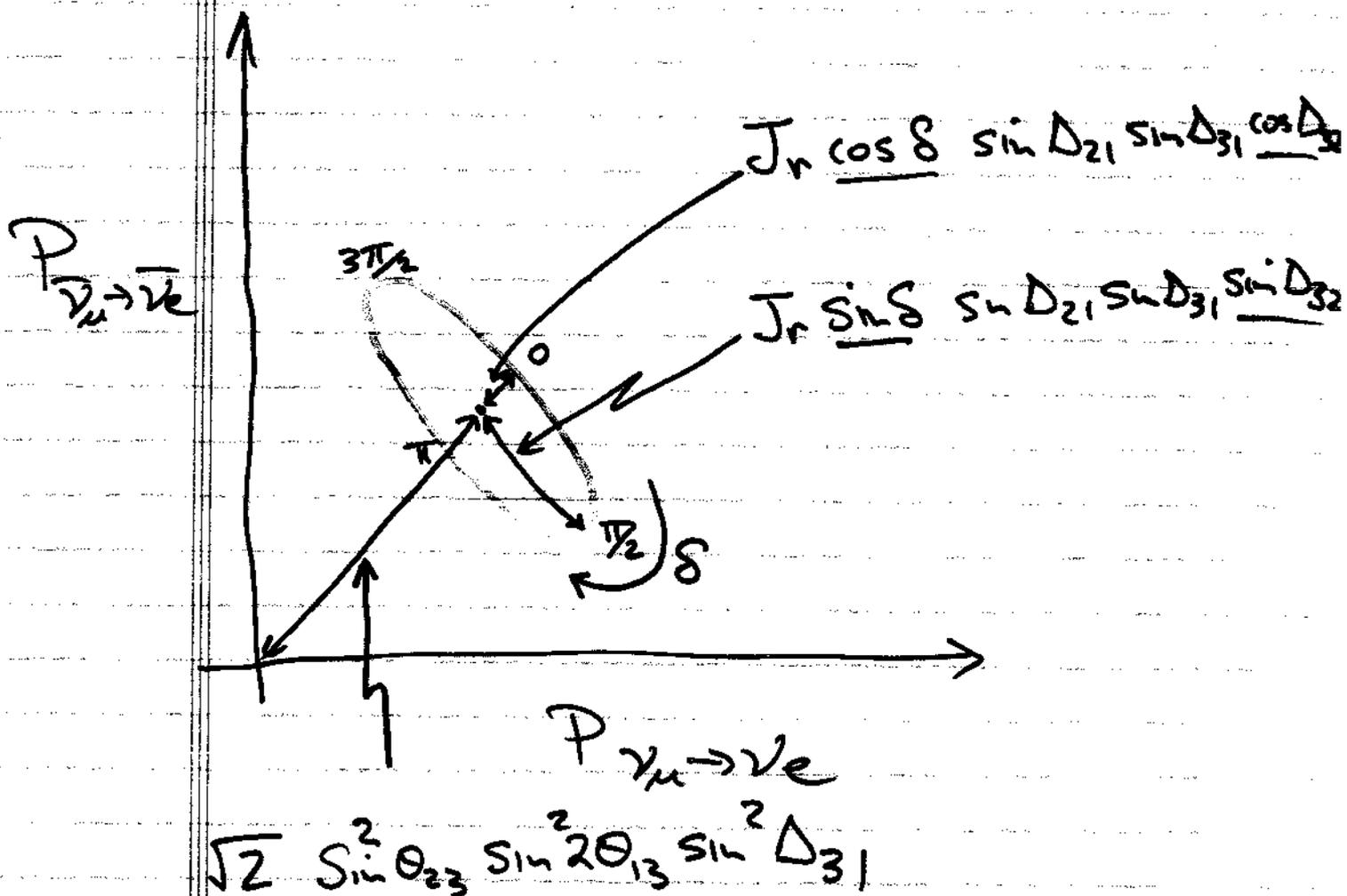
$$P_{\odot} = C_{13}^2 C_{23}^2 S_{2(12)}^2 \sin^2 \Delta_{21}$$

$P_{\odot}$  is irrelevant unless  $S_{13}$  is small  
where  $(1 - 2 S_{13} \dots) \approx 1$

$$\frac{\Delta_{21}}{\Delta_{31}} = \frac{\sin^2 \Delta_{21}}{\sin^2 \Delta_{31}} \approx \frac{1}{30} \quad \text{also } \sin \Delta_{31} \approx 1 \\ \cos \Delta_{32} \approx 0$$

# Bi-Probability Plot

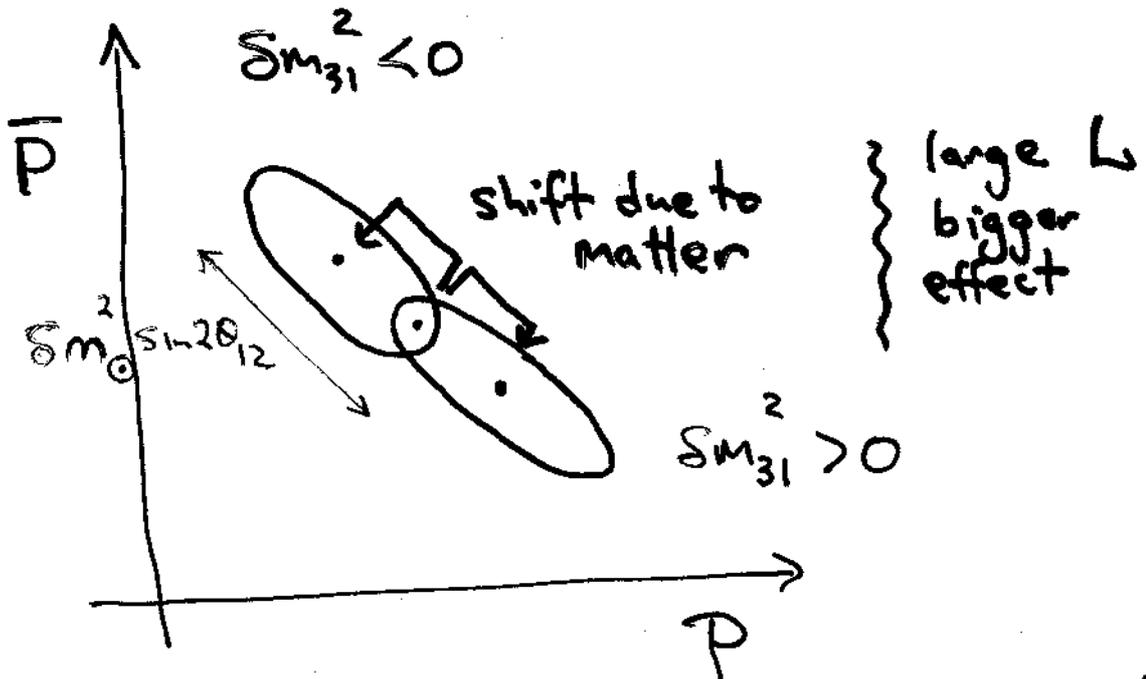
(P, P̄) varying only  $\delta_{CP}$



IF SMA or LOW solar sol<sup>n</sup> ellipse collapses to point.

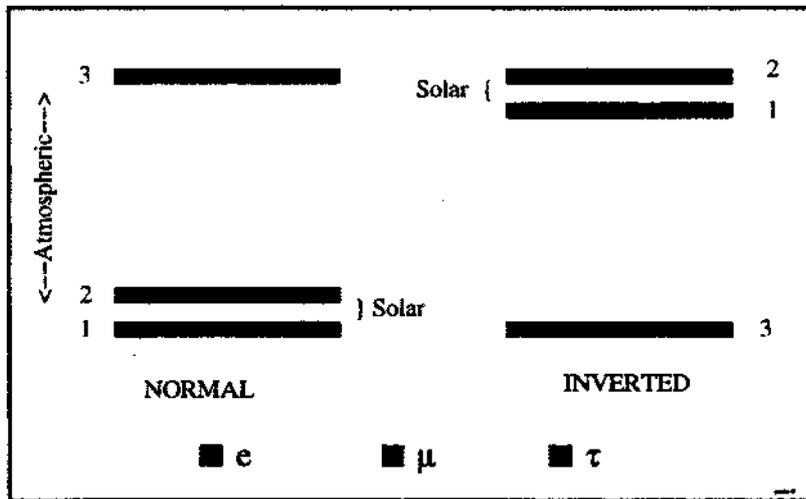
LMA - Waw  $\frac{DD}{000}$

# Matter Effects:



$$P_{\text{mat}}^{\text{center}} \approx \left(1 \pm \frac{E_\nu}{6\text{GeV}}\right) P_{\text{vac}}^{\text{center}}$$

O.M.  
for small  $E_\nu$



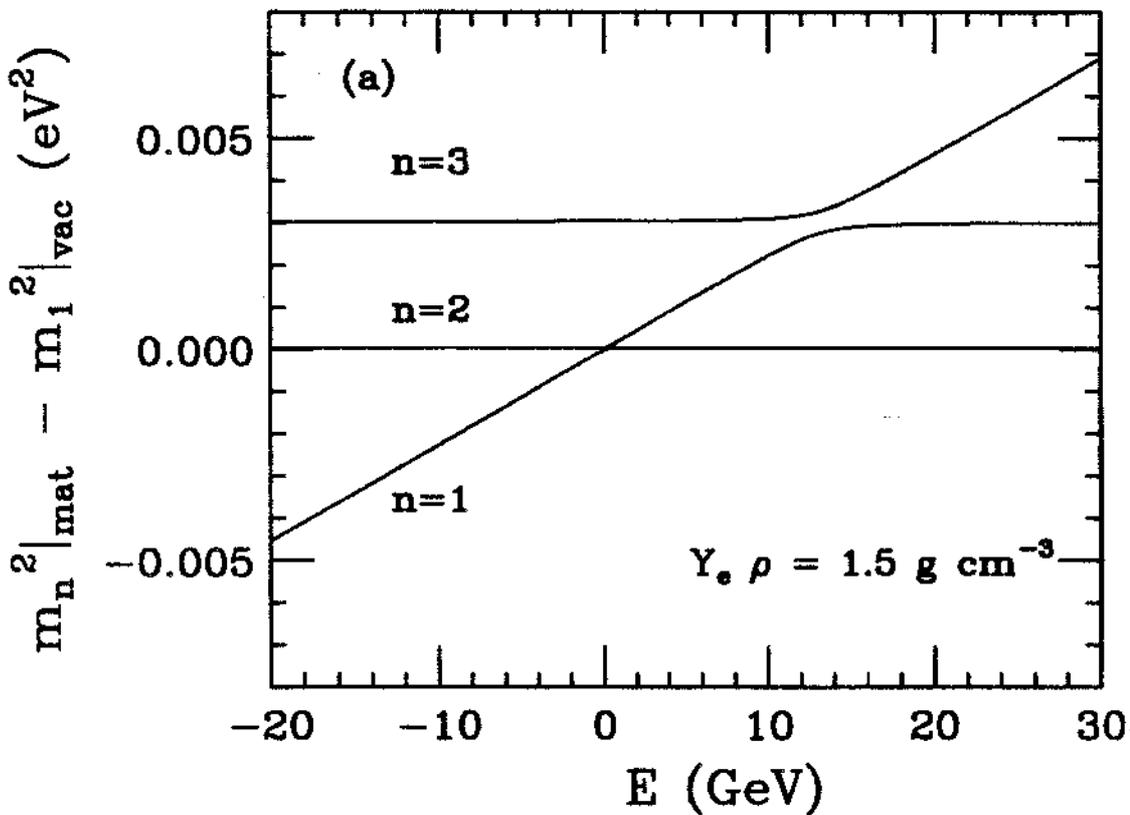
$$\delta m_{31}^2 > 0$$

$$\delta m_{31}^2 < 0$$

$$\left( \sin 2\theta_{13} \delta M_{31}^2 \right)_{\text{matter}} = \left( \sin 2\theta_{13} \delta M_{31}^2 \right)_{\text{vac}} \quad 19$$

and

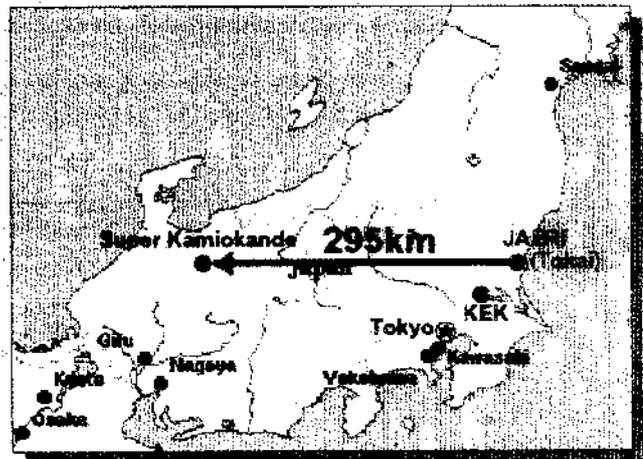
$$\left( \sin 2\theta_{12} \delta M_{21}^2 \right)_{\text{MATTER}} = \left( \sin 2\theta_{12} \delta M_{21}^2 \right)_{\text{vac}}$$



← antineutrinos | → neutrinos

## JHF → Super-Kamiokande

- ✓ 295 km baseline
- ✓ Super-Kamiokande:
  - 22.5 kton fiducial
  - Excellent  $e/\mu$  ID
  - Additional  $\pi^0/e$  ID
- ✓ Hyper-Kamiokande
  - 20× fiducial mass of SuperK
- ✓ Matter effects small
- ✓ Study using fully simulated and reconstructed data



Requires New Beamline:

~~<http://www.off-axis.fnal.gov/>~~

LOI: hep-ex/0106019

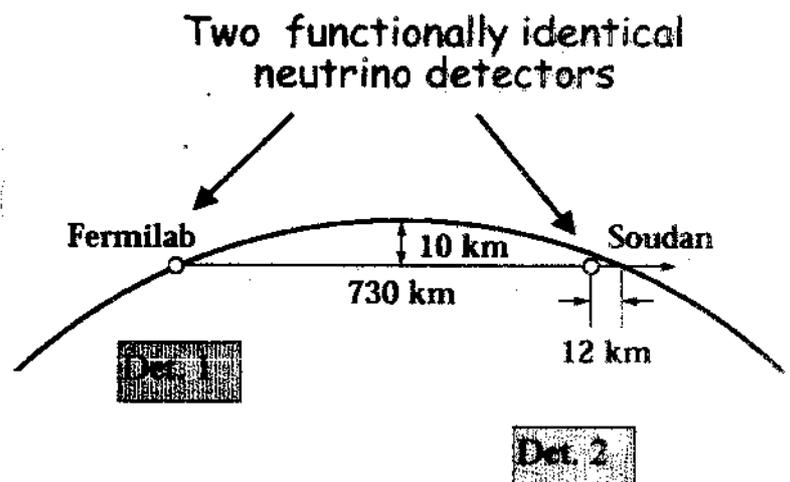
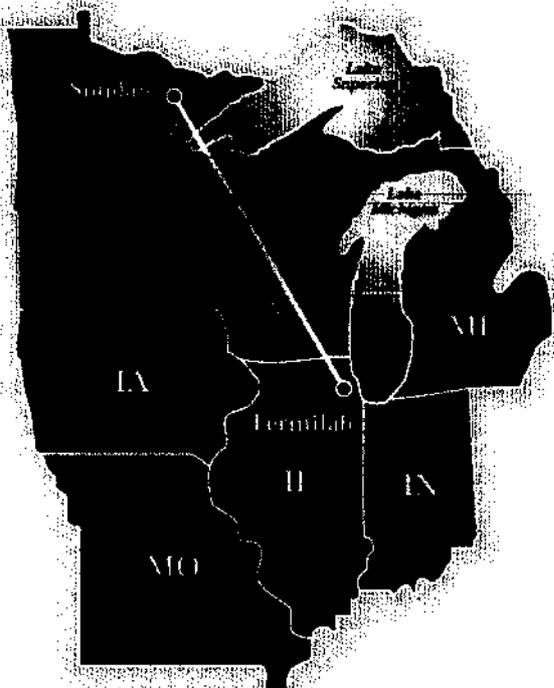
<http://www-nu.kek.jp/jhfnu/>

$$\bar{E}_\nu = 0.6 \text{ to } 1.0 \text{ GeV}$$

~20% spread

$$L = 295 \text{ km}$$

# The NUMI Beamline



New Detector Required:

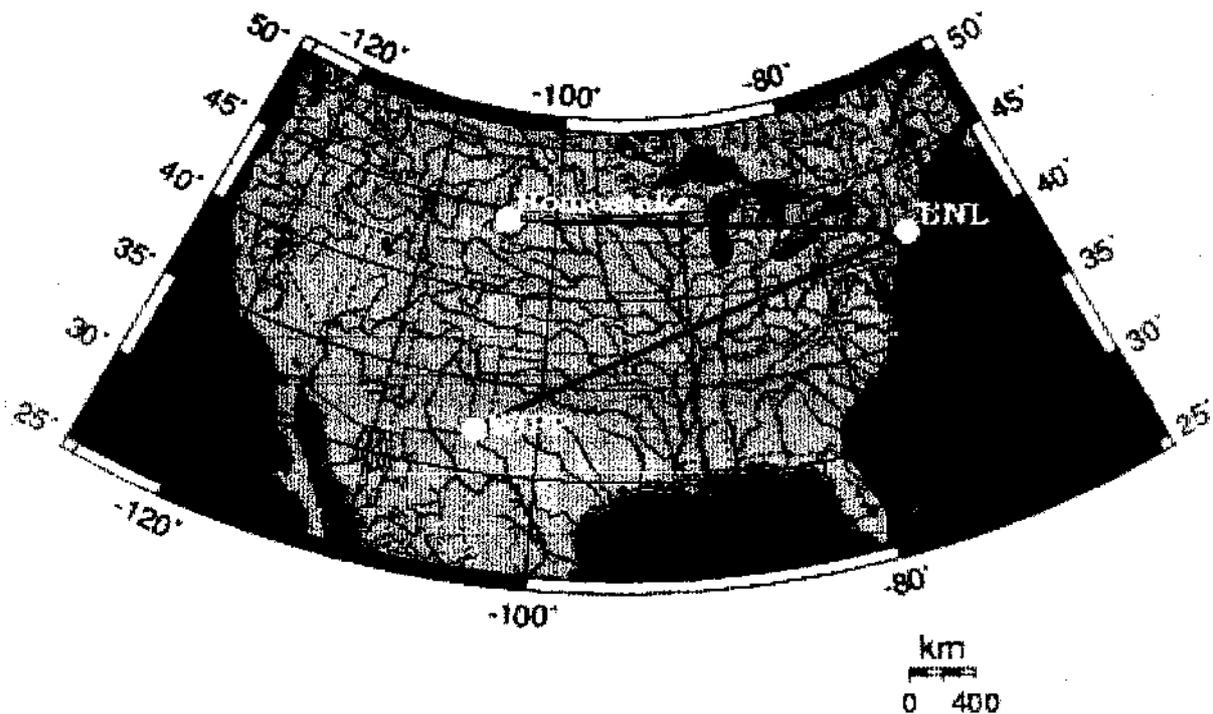
<http://www-off-axis.fnal.gov/>

LOI: hep-ex/0210005

$$E \sim 2 \text{ GeV}$$

$$L \sim 732 \text{ km} \\ \pm 200 \text{ km}$$

# Brookhaven to Homestake OR WIPP



$L = \underline{2540 \text{ km}}$  or 2880 km

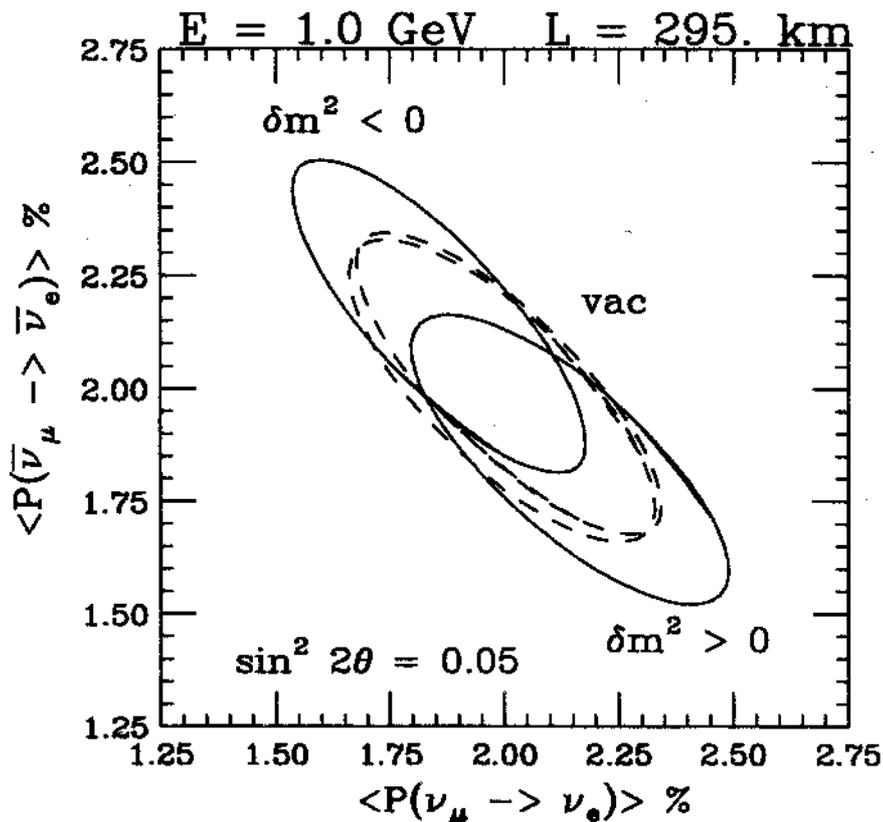
*2<sup>nd</sup> peak 2 GeV:*

New Beamline, New Detector:

<http://www.neutrino.bnl.gov/>

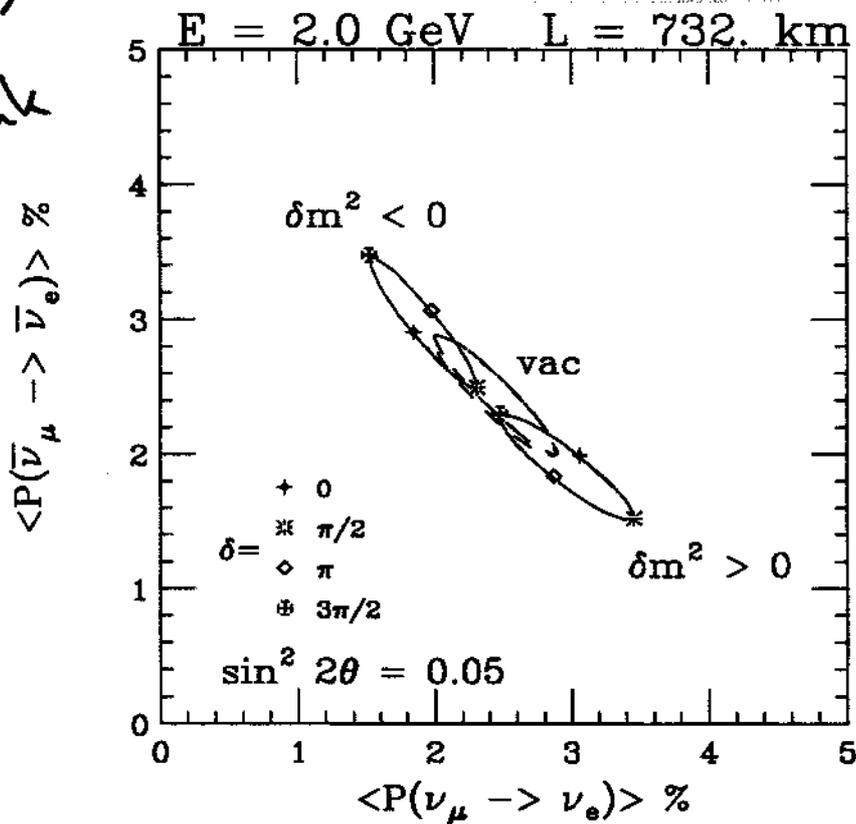
LOI: hep-ex/0205040

JHF  $\rightarrow$  SK  
1<sup>st</sup> peak



20% Spread

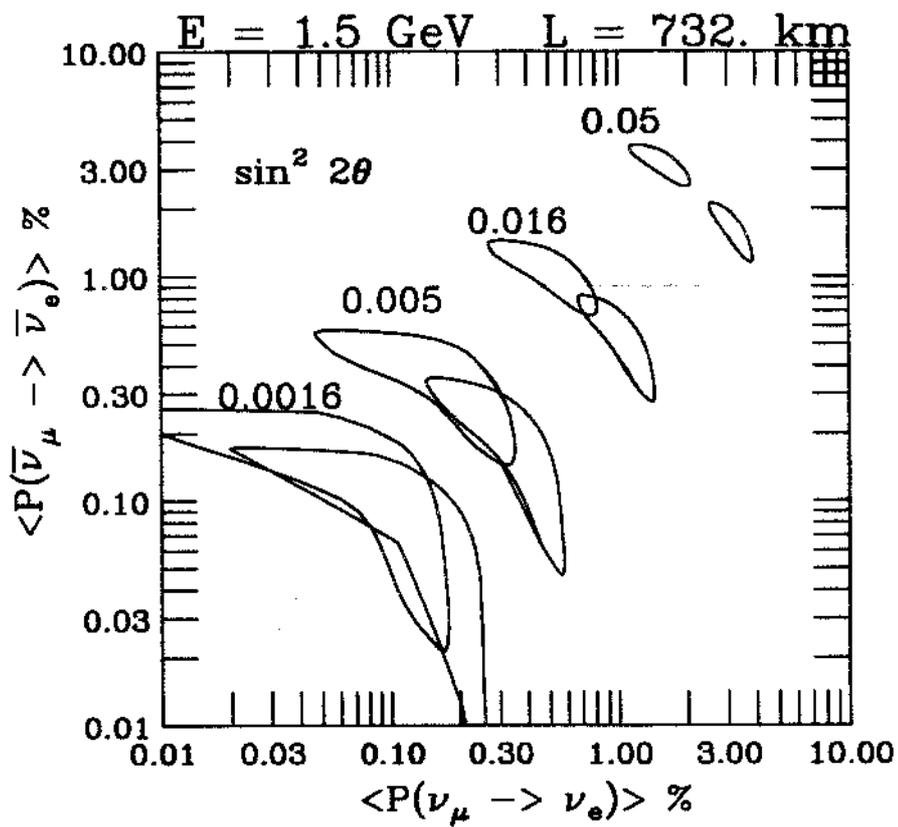
NuMI  $\rightarrow$   
1<sup>st</sup> peak



20% Spread.

Matter Effects  
separate  
 $\delta m^2 > 0$   
from  
 $\delta m^2 < 0$

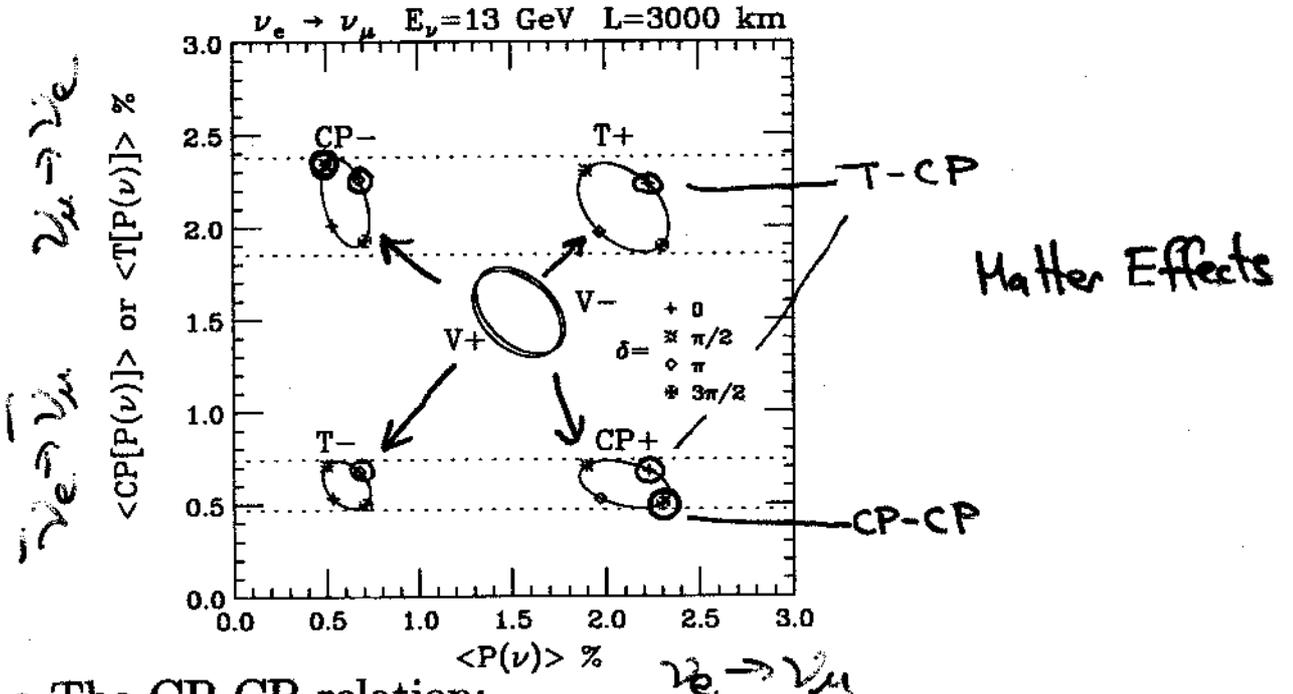
As  $\sin^2 2\theta_{13}$  Varies:



NuMI Off-Axis LOI

hep-ex/0210005

Anatomy of the Bi-Probability Plot:



- The CP-CP relation:

$$\begin{aligned}
 & P(\nu_e \rightarrow \nu_\mu; \Delta m_{31}^2, \Delta m_{21}^2, \delta, a) \\
 &= P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, -\Delta m_{21}^2, \delta, a) \\
 &\approx P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, +\Delta m_{21}^2, \pi + \delta, a)
 \end{aligned}$$

← evolution eqn.

- The T-CP relation:

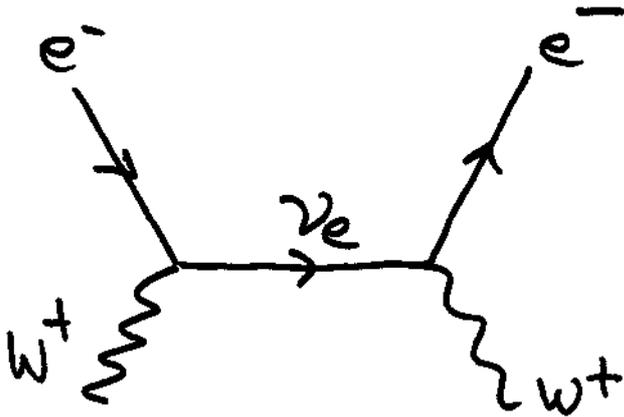
assuming symmetry matter distribution

$$\begin{aligned}
 & P(\nu_\mu \rightarrow \nu_e; \Delta m_{31}^2, \Delta m_{21}^2, \delta, a) \\
 &= P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, -\Delta m_{21}^2, -\delta, a) \\
 &\approx P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu; -\Delta m_{31}^2, +\Delta m_{21}^2, \pi - \delta, a)
 \end{aligned}$$

- $\approx$  trade sign of  $\delta m_{12}^2$  for shift by  $\pi$  of  $\delta$ :

$$(\dots) + \delta m_{12}^2 [(\dots) \cos \delta + (\dots) \sin \delta]$$

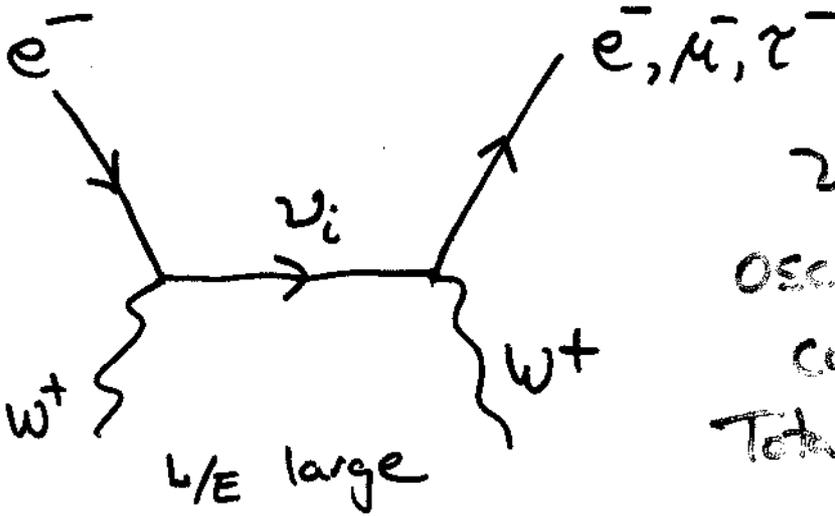
↑ Small correction here



Days of  
Weinberg + Salam 21

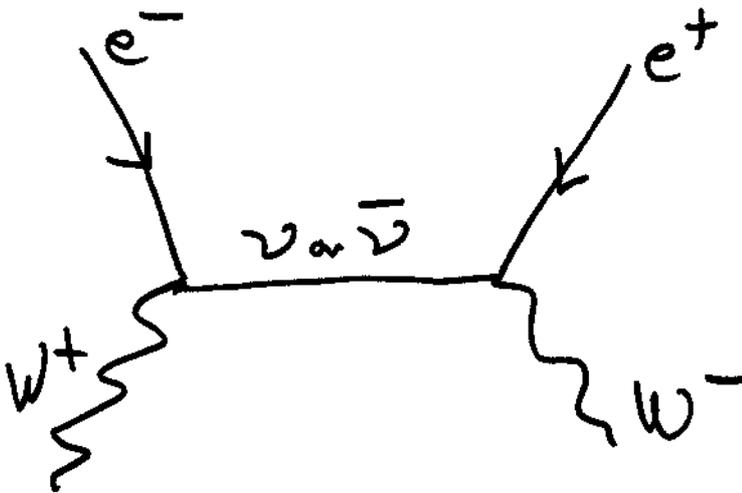
S.M.

flavor diagonal  
conserves  
 $L_e, L_\mu, L_\tau$



$\nu$  MASS  
oscillations ~~(flavor)~~  
conserves  
Total Lepton #

$L/E$  large



Majorana Mass  
Lepton # Violating  
(highly suppressed)  
Why ???

# Fermion Mass: (cartoon:) 27

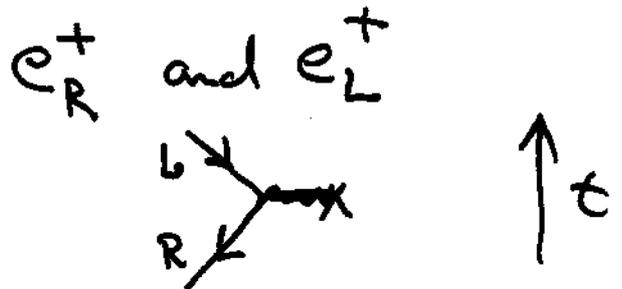
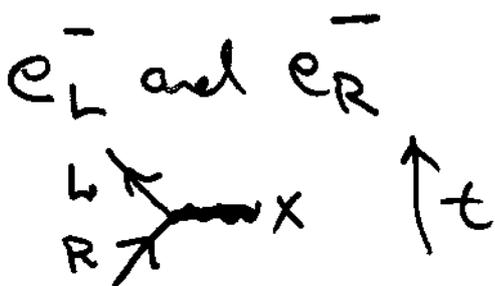
- Consider a massless electron:

$$\begin{array}{cc}
 e_L^- & e_R^- \\
 e_R^+ & e_L^+
 \end{array}$$

- Four States:  $2e^- + 2e^+$

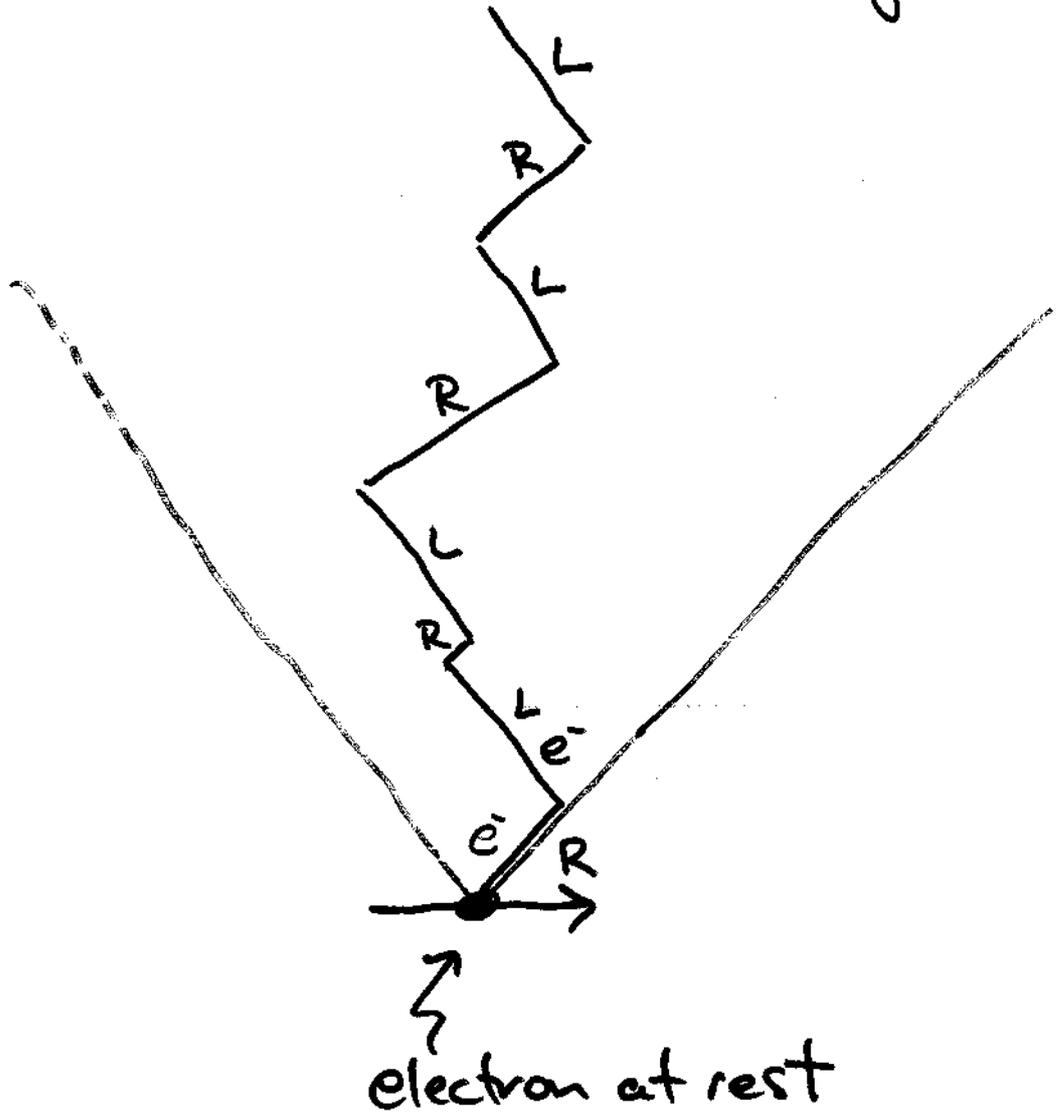


- Mass is a coupling between Left and Right. (chiral states)



CANNOT couple  
 $e^-_L$  to  $e^+_R$

because of electric charge conservation.



# Massive Fermion with given Spin: <sup>29</sup>

$$P^2 = M^2 \quad S^2 = -1 \quad P \cdot S = 0$$

e.g. (at rest  $P = (M, 0, 0, 0)$   $S = (0, 0, 0, 1)$   
Spin in z direction

$$P = \frac{P+MS}{2} + \frac{P-MS}{2}$$

$$\frac{P+MS}{2}$$

$$\text{and } \frac{P-MS}{2}$$

Massless

$$\begin{aligned} \left( \frac{P \pm MS}{2} \right)^2 &= \frac{P^2 \pm 2P \cdot S + M^2 S^2}{4} \\ &= \frac{M^2 \pm 2M \cdot 0 + M^2(-1)}{4} \\ &= 0 \end{aligned}$$

Sol<sup>n</sup> Dirac Eqn  
Massive, Spin

massless  
spinors

$$U(P, S) = \frac{1+\gamma_5}{2} U\left(\frac{P+MS}{2}\right)$$

$$+ e^{2i\phi} \frac{1-\gamma_5}{2} U\left(\frac{P-MS}{2}\right)$$

chiral  
projections

$$u(p, s) = \frac{1 + \gamma_5 \not{s}}{2} u(p)$$

$$= \frac{1 + \gamma_5}{2} \underbrace{\frac{1 + \not{s}}{2} u(p)}_{Q_+} + \frac{1 - \gamma_5}{2} \underbrace{\frac{1 + \not{s}}{2} u(p)}_{Q_-}$$

$$Q_+ \overline{Q_+} = \frac{p \pm Ms}{2}$$

$$\text{and } \frac{p \pm Ms}{2} Q_{\pm} = 0$$

$$\text{Therefore } Q_{\pm} = u\left(\frac{p \pm Ms}{2}\right)$$

Relative phase determined by

$$\overline{u(p, s)} u(p, s) = 2M$$

MASSIVE PARTICLE AT REST,  
Spin in z direction:

31

$$P = \frac{P+MS}{2} + \frac{P-MS}{2}$$

$$(M, 0, 0) = \frac{M}{2} (1, 0, 0, 1) + \frac{M}{2} (1, 0, 0, -1)$$

BIG BOOST IN z-DIRECTION  $\begin{pmatrix} \gamma & \beta\gamma \\ \beta\gamma & \gamma \end{pmatrix}$

$$\gamma M (1, 0, 0, \beta) = \frac{M}{2} (\gamma + \beta\gamma) (1, 0, 0, 1) + \frac{M}{2} (\gamma - \beta\gamma) (1, 0, 0, -1)$$

as  $\beta \rightarrow 1$   $(\gamma + \beta\gamma) \rightarrow 2\gamma$   $(\gamma - \beta\gamma) \rightarrow \frac{1}{2}\gamma$

where  $\gamma = E/M$

$$E (1, 0, 0, \beta) \approx E (1, 0, 0, 1) + \frac{M^2}{E} (1, 0, 0, -1)$$

$$U(P, S) = \frac{1+\gamma_5}{2} U\left(\frac{P+MS}{2}\right) + \frac{1-\gamma_5}{2} U\left(\frac{P-MS}{2}\right)$$

HELICITY  
STATE

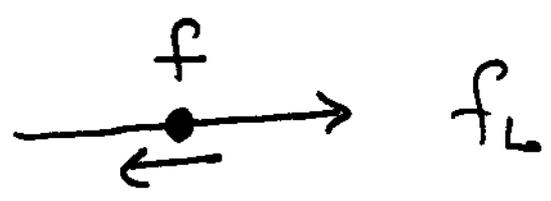
RIGHT

LEFT

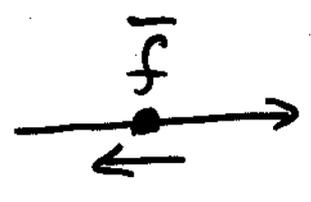
CHIRAL PROJECTIONS

HELICITY  $\neq$  CHIRALITY  
 for massive particles

Massless Fermion:



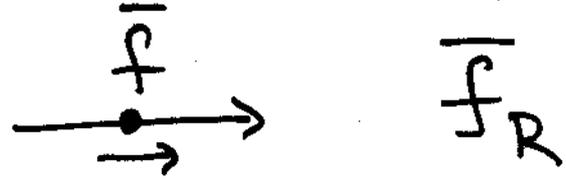
Charge Conjugation (C)



Parity (P)



Time Reversal (T)

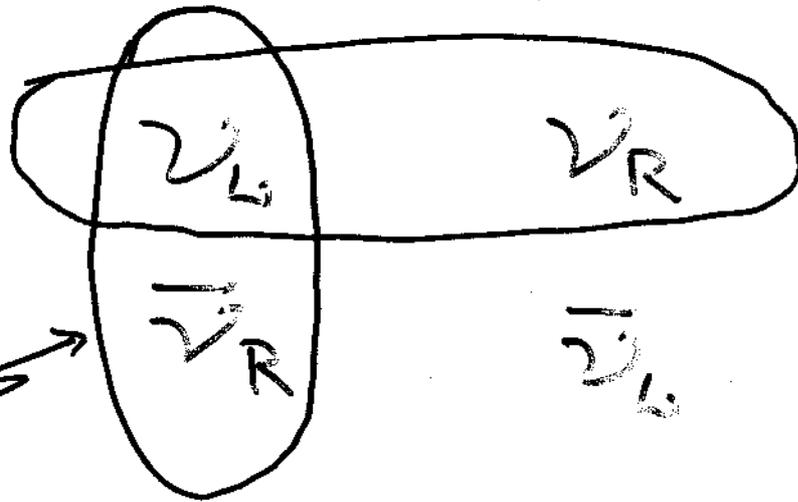


CPT  $f_L \leftrightarrow \bar{f}_R$

At a minimum by CPT

$f_L, \bar{f}_R$

For Neutral Fermions (neutrinos) <sup>3:</sup>



Dirac Mass.

required by weak interaction

Majorana Mass

(Violates  $b \neq$ )

See-Saw Mechanism gives

$$\begin{pmatrix} \nu_L \\ \nu_R \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} N_R \\ N_L \end{pmatrix}$$

light  $\frac{m^2}{M}$   $M$

Neutrinos being Majorana

is the minimal solution

to giving  $\nu$  mass.

$\nu = \bar{\nu}$       two states  $\nu_L$  and  $\nu_R$

---

If Dirac then there is a global symmetry (lepton #) which needs to be explained.

$\nu \neq \bar{\nu}$  and four states

weak  $\rightarrow$   $\nu_L$  and  $\bar{\nu}_R$   
 $\nu_R$  and  $\bar{\nu}_L$

Suppression factor:

$$\left( \frac{\overline{M}_\nu}{E} \right)^2$$

chirality

≠

helicity

for massive particles

$\overline{M}_\nu \sim 1 \text{ eV}$      $E = 1 \text{ GeV}$      $\left( \frac{M}{E} \right)^2 \approx \underline{\underline{10^{-20}}}$

BUT

⊗  $2\nu\beta\beta$  decay



Same Nucleus also has  $2\nu\beta\beta$  decay

→ end point:

## WHAT WE DON'T KNOW:

36

- Majorana OR Dirac
- Absolute mass of lightest neutrino.  
(except  $< \sim 1\text{eV}$ )
- Size of  $\theta_{13}$ : ( $\nu_e$  in the "3" state.)  
 $\sin^2 \theta_{13} < 0.03$
- Is  $\theta_{23} = \text{or } \lesseqgtr \frac{\pi}{4}$  the  $\mu \leftrightarrow \tau$  symmetric point.  
(maximal mixing)  
 $0.35 < \sin^2 \theta_{23} < 0.65$
- Sign of  $\Delta m_{21}^2$  (  $\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \text{ or } \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array}$  )  
type of spectrum
- phase  $\delta \ll \text{if } \neq 0$  leads to CP violation
- Number of light Neutrinos: 3 or are there more than 3