

New Ideas in Randall-Sundrum Phenomenology

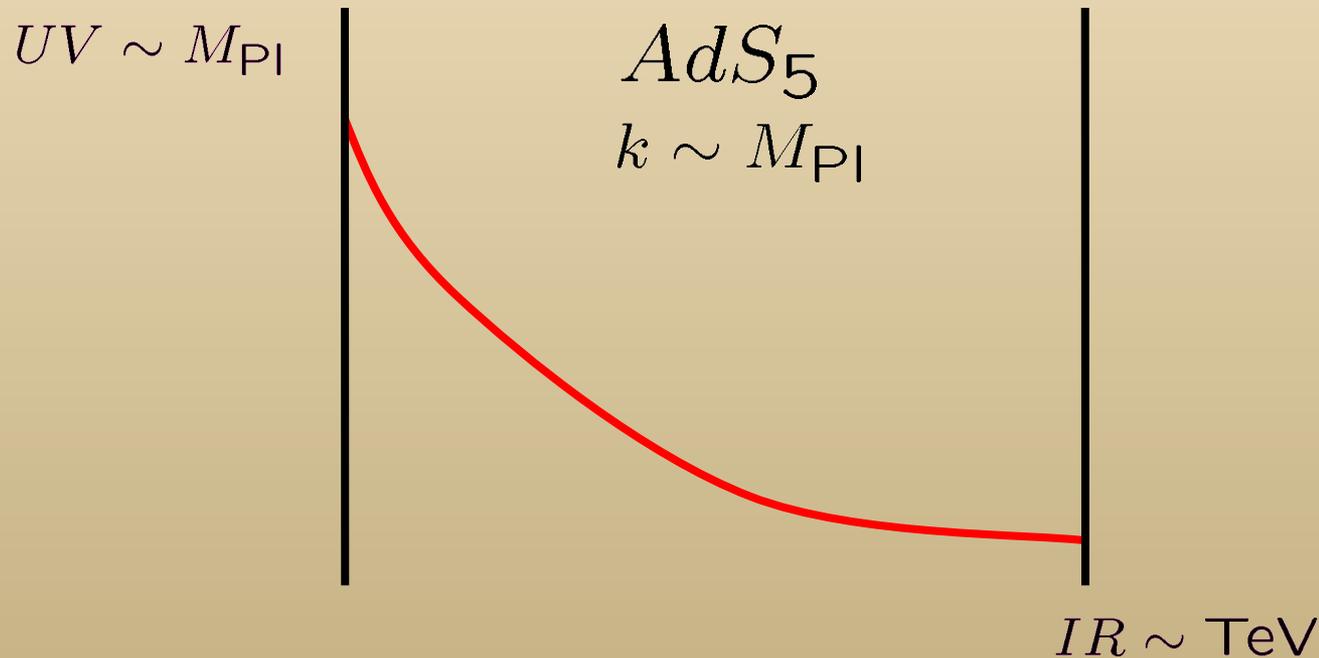
José Santiago
Fermilab

work with M. Carena, E. Pontón, C. Wagner

NPB (06) [ph/0607106]
and hep-ph/0701055
+ work in progress

Randall-Sundrum 101

- RS offers a nice solution to the hierarchy problem

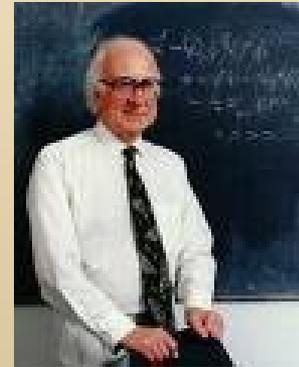


Randall-Sundrum 101

- RS offers a nice solution to the hierarchy problem

$UV \sim M_{\text{Pl}}$

AdS_5
 $k \sim M_{\text{Pl}}$



$\langle \phi \rangle \sim \tilde{k} \equiv ke^{-kL} \sim \text{TeV}$

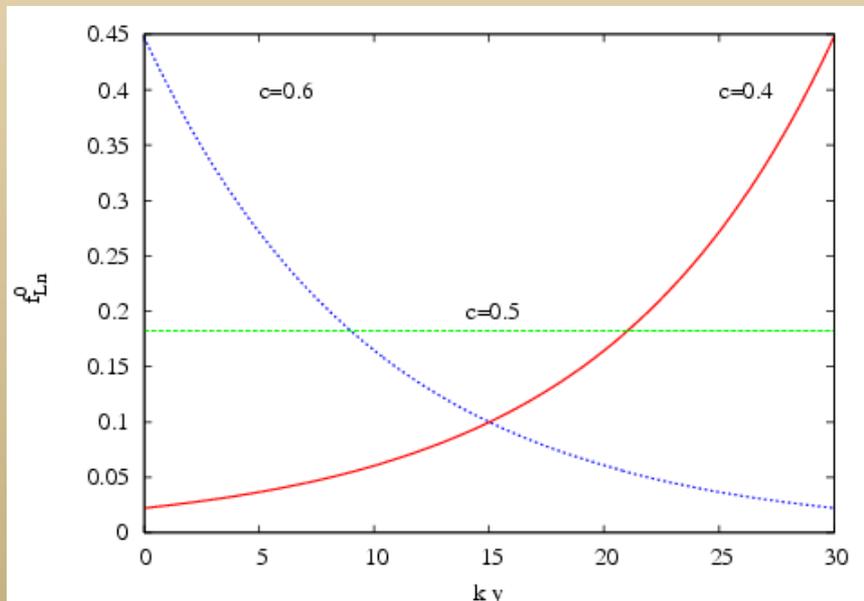
$IR \sim \text{TeV}$

Randall-Sundrum 101

- Fermions (and gauge bosons) in the bulk
 - nice flavor theory: hierarchical masses + ~~FCNC~~
 - Higher-dimensional operators suppressed

Randall-Sundrum 101

- Fermions (and gauge bosons) in the bulk
 - nice flavor theory: hierarchical masses + ~~FCNC~~
 - Higher-dimensional operators suppressed
- Fermion localization from bulk mass $\mathcal{L}_m = c_\psi k \bar{\Psi} \Psi$



+ + b.c. zero mode

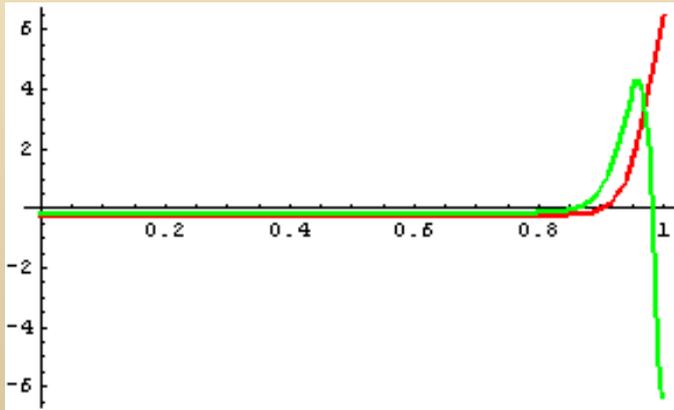
It is also possible to
use different b.c.

~~- + b.c. zero mode~~

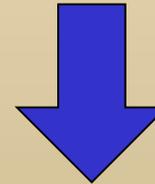
KK modes might be light

Randall-Sundrum faces reality

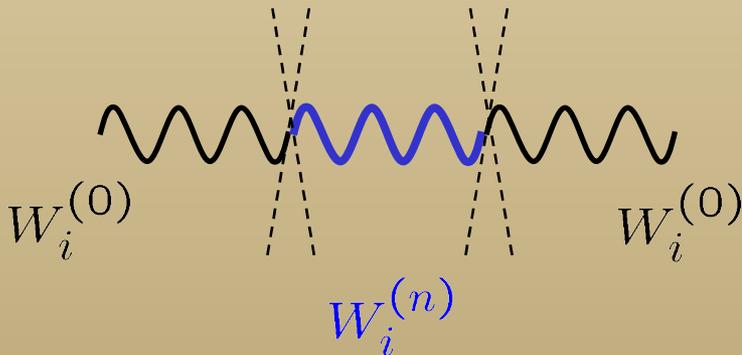
- Gauge boson KK modes localized towards IR brane



Large mixing with the
Z and W zero modes
through the Higgs



Large T parameter

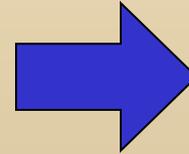


$$M_{KK} \gtrsim 5 - 10 \text{ TeV}$$

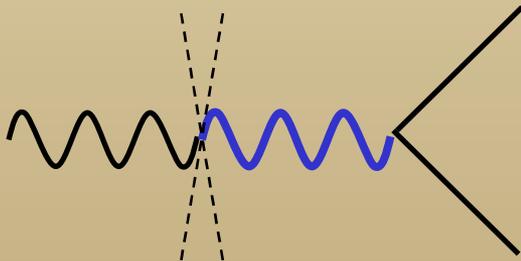
Randall-Sundrum faces reality

- Top (bottom) zero modes localized towards IR brane

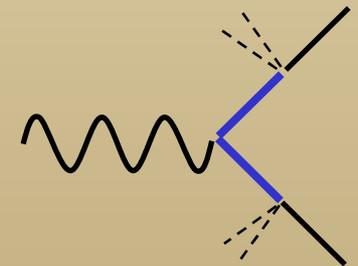
Large gauge and Yukawa couplings to GB and fermion KK modes



Large anomalous Zbb coupling



$$M_{KK} \gtrsim 7 - 8 \text{ TeV}$$



Randall-Sundrum faces reality

Large corrections to the
T parameter and the
Zbb coupling force \tilde{k} to be
beyond the LHC reach

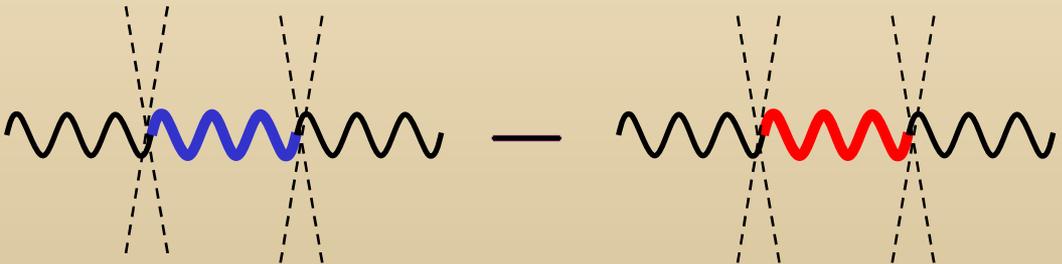
- No Randall-Sundrum physics at the LHC
- The model becomes very fine-tuned

RS needs protection: call custodial symmetry

Agashe, Delgado, May,
Sundrum JHEP (03)

■ Custodial protection of Randall-Sundrum

- Bulk gauge symmetry: $SU(2)_L \times SU(2)_R \times U(1)_X \times P_{LR}$

$$T \propto \text{[Diagram 1]} - \text{[Diagram 2]} \sim 0$$


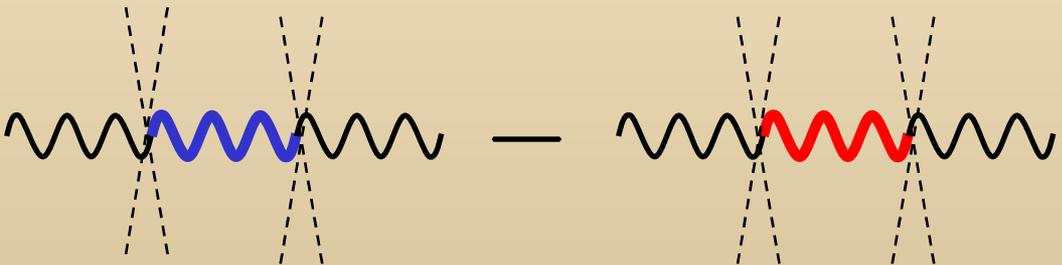
The diagram illustrates the cancellation of terms in the trace T . It shows two Feynman diagrams representing the trace of the product of two operators. The first diagram features a blue wavy line in the middle, and the second diagram features a red wavy line in the middle. Both diagrams consist of external wavy lines and internal dashed lines forming a loop structure. The two diagrams are subtracted, resulting in a trace T that is approximately zero.

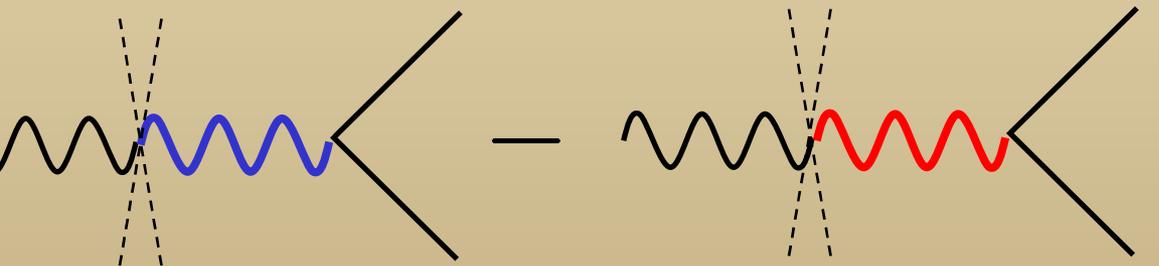
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Agashe, Contino, Da
Rold, Pomarol PLB (06)

■ Custodial protection of Randall-Sundrum

- Bulk gauge symmetry: $SU(2)_L \times SU(2)_R \times U(1)_X \times P_{LR}$

$$T \propto \text{Diagram 1} - \text{Diagram 2} \sim 0$$


$$\delta g_{b_L} \propto \text{Diagram 3} - \text{Diagram 4} \sim 0$$


Symmetry broken by b.c. to $SU(2)_L \times U(1)_Y$ on the UV

Fermion Quantum Numbers

$$P_{LR} \Rightarrow T_R^3(b_L) = T_L^3(b_L)$$

The simplest option is bidoublets under $SU(2)_L \times SU(2)_R$

$$\begin{array}{c}
 \xleftrightarrow{SU(2)_R} \\
 \left(\begin{array}{cc}
 \chi_L^u(-+) & t_L(+,+) \\
 \chi_L^d(-+) & b_L(+,+)
 \end{array} \right)_X \sim (2, 2)_{2/3} \sim \begin{array}{c} U(1)_Q \\ \left(\begin{array}{cc}
 5/3 & 2/3 \\
 2/3 & -1/3
 \end{array} \right) \end{array} \\
 \updownarrow SU(2)_L
 \end{array}$$

The Higgs is also a bidoublet with $Q_X = 0$

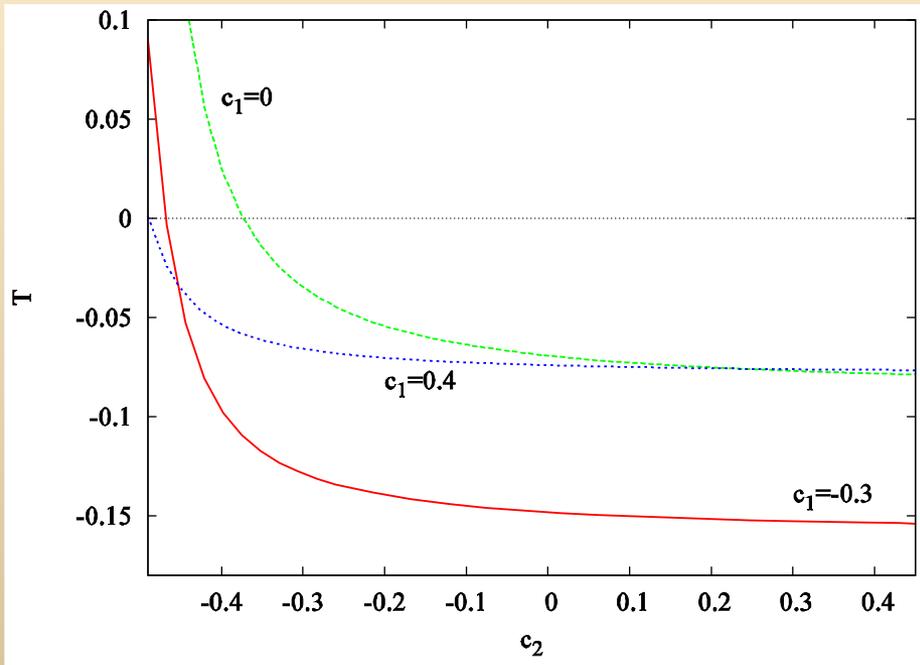
$$t_R(+,+) \sim (1, 1)_{2/3}$$

Calculable corrections to T and Zbb

Carena, Pontón, J.S., Wagner NPB (06)

- T and Zbb protected (at tree level) by custodial symmetry broken on the UV brane:
 - Quantum corrections are calculable (finite)
 - T and Zbb are insensitive to physics above Λ
- One loop corrections are important
 - Bidoublets contribute **negative** to T
 - Singlets contribute **positive** to T
 - Large positive T leads to large positive δg_{bL}

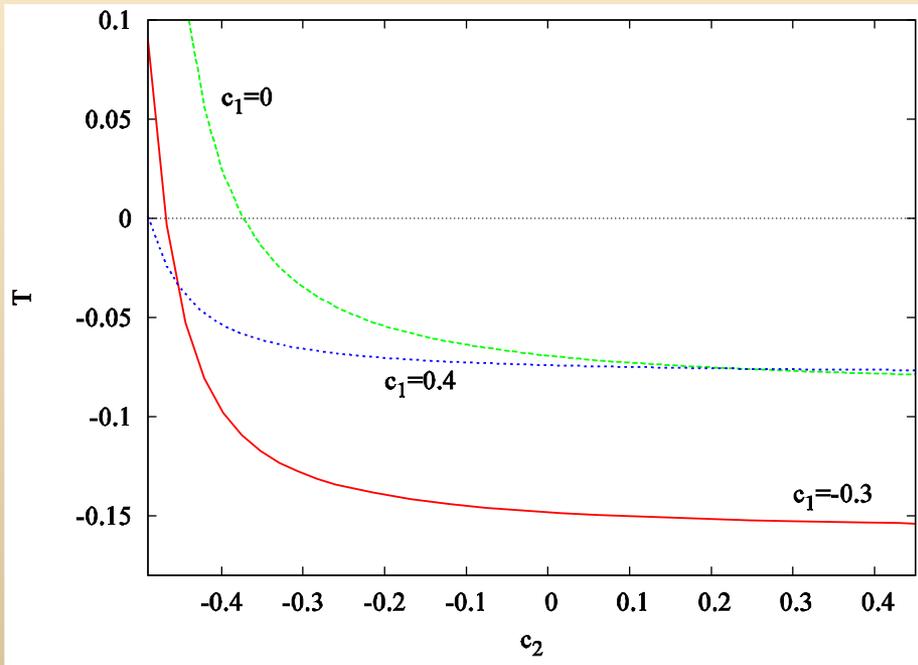
Calculable corrections to T and Zbb



T negative in most of parameter space: the RH top is almost flat (LH top/bottom near IR)

UV ← singlet localization → IR

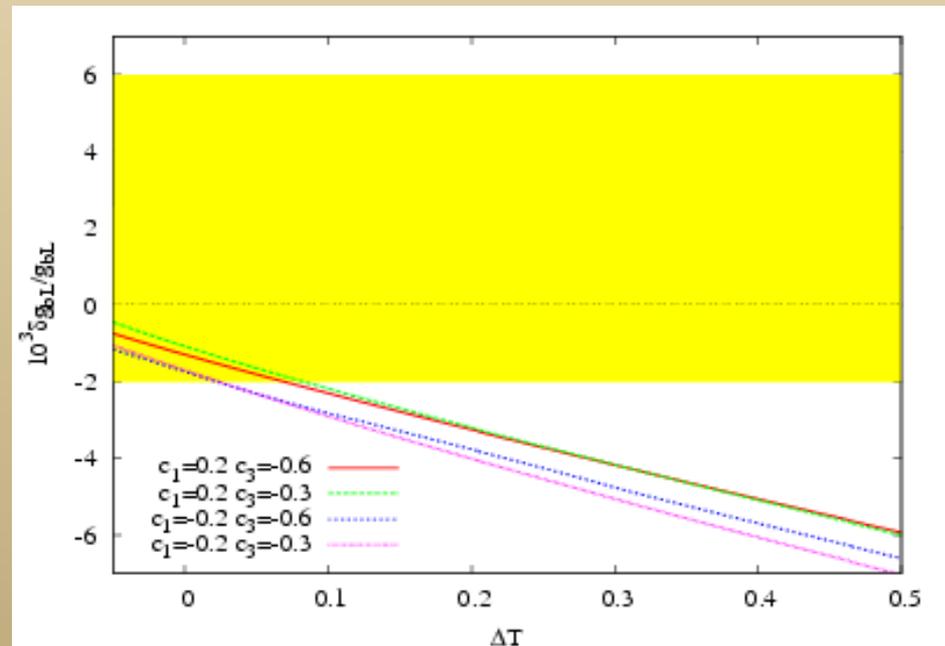
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Positive T leads to large deviations in Zbb



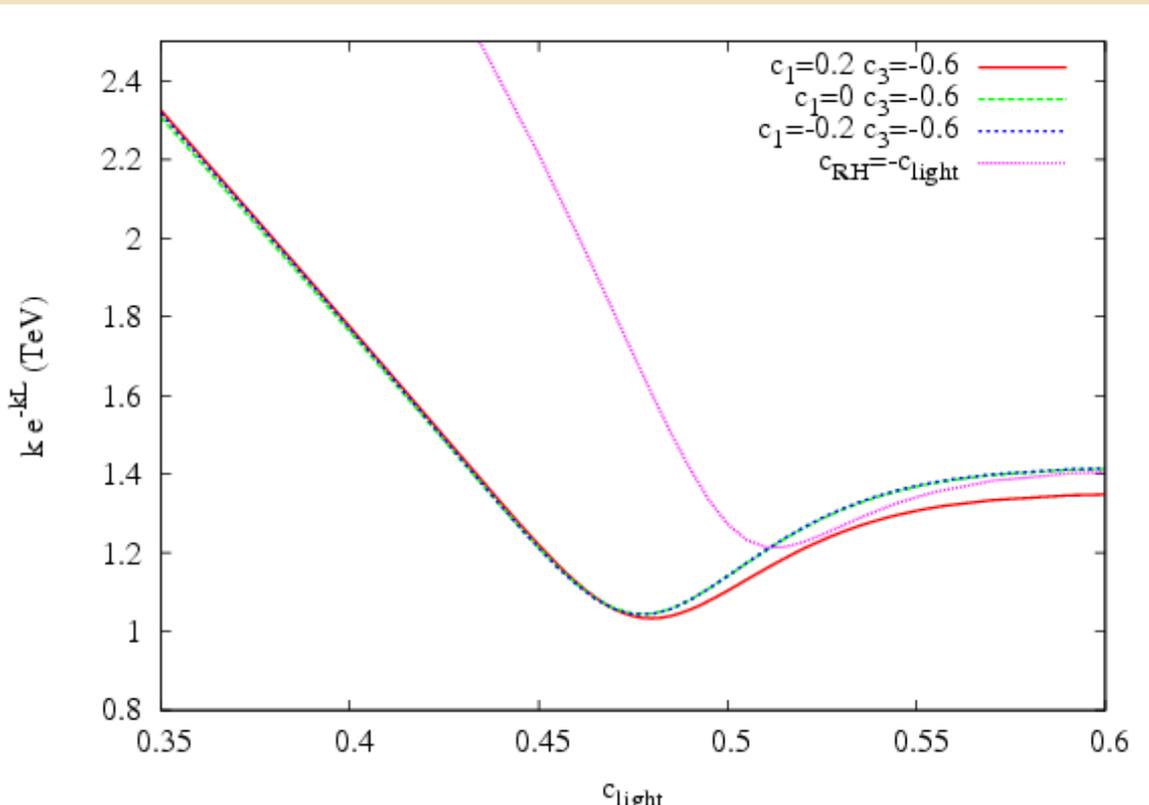
Global Fit to Electroweak Observables

Carena, Pontón, J.S., Wagner hep-ph/0701055

Han, Skiba PRD (05-06)

- We have performed a global fit to all EW observables including:
 - All tree-level effects at leading order in $\mathcal{O}(v^2/\tilde{k}^2)$
 - Leading one loop effects on: **T**, **Zbb** and S
- We compute the χ^2 as a function of the localization parameter of the fermions $t_L, t_R, b_R, f_{\text{light}}^{LH}, f_{\text{light}}^{RH}$ and \tilde{k}
- To reduce number of parameters we optimize with respect to $t_R, f_{\text{light}}^{RH}$

Global Fit to Electroweak Observables



$\tilde{k} \sim 1 - 1.4 \text{ TeV}$
 \Downarrow
 $M_{KK}^{\text{gauge}} \sim 2.5 - 3.5 \text{ TeV}$

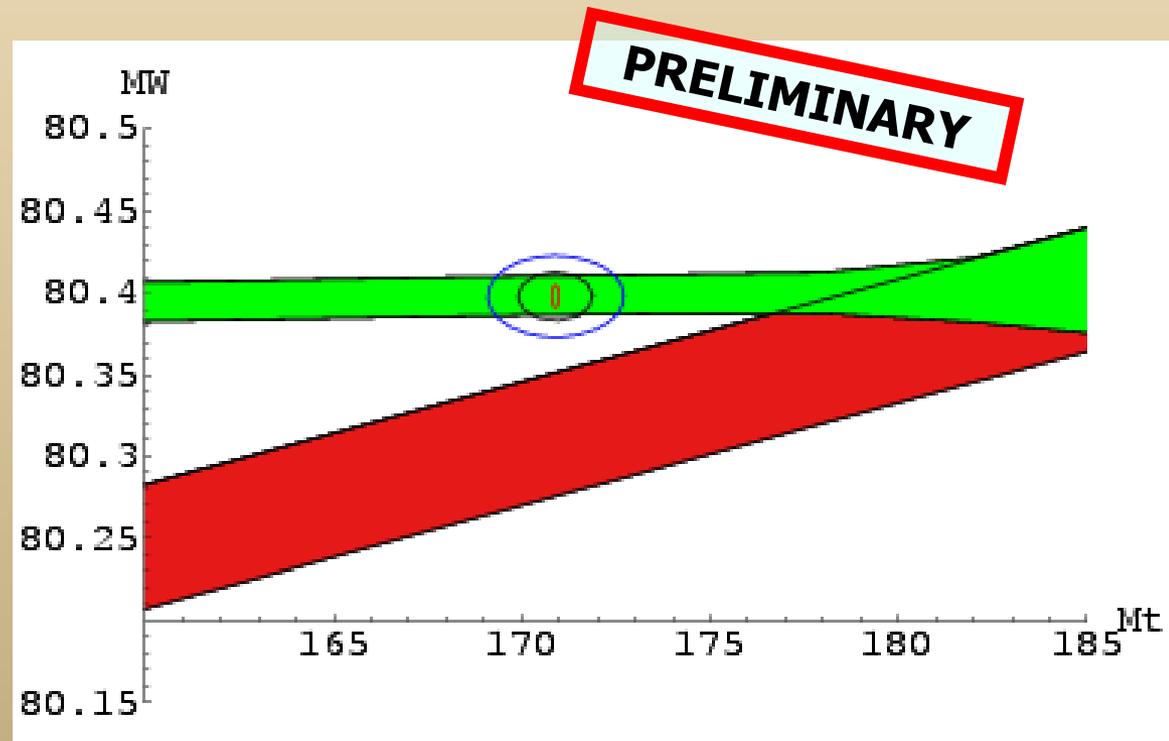
← IR light localization → UV

Global Fit to Electroweak Observables

- New measurements of m_t and m_W generate a tension in the SM
- The contribution in our model goes in the right direction

$$M_w = 80.398 \pm 0.025 \text{ GeV}$$

$$M_{\text{top}} = 170.9 \pm 1.8 \text{ GeV}$$



Fermion Spectrum

First two generation KK modes

q'	Q	$m_{q'}$ (GeV)	decay
q_1^1	2/3	$\sim 200 - 500$	$q_1^1 \rightarrow Zu, (100\%)$
q_1^2	2/3	$\sim 200 - 500$	$q_1^2 \rightarrow Zc, (100\%)$
q_2^1	2/3	$\sim 200 - 500$	$q_2^1 \rightarrow Hu, (100\%)$
q_2^2	2/3	$\sim 200 - 500$	$q_2^2 \rightarrow Hc, (100\%)$
χ_2^{u1}	5/3	$\sim 200 - 500$	$\chi_2^{u1} \rightarrow Wu, (100\%)$
χ_2^{u2}	5/3	$\sim 200 - 500$	$\chi_2^{u2} \rightarrow Wc, (100\%)$
q'^{d1}	-1/3	$\sim 200 - 500$	$q'^{d1} \rightarrow Wu, (100\%)$
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Not a necessary prediction

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Fermion Spectrum

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q_1	2/3	369	$q_1 \rightarrow Zt$, (20%) $q_1 \rightarrow Ht$, (60%) $q_1 \rightarrow Wb$, (20%)
q_2	2/3	373	$q_2 \rightarrow Zt$, (9%) $q_2 \rightarrow Ht$, (70%) $q_2 \rightarrow Wb$, (21%)
u_2	2/3	504	$u_2 \rightarrow Zt$, (13%) $u_2 \rightarrow Ht$, (40%) $u_2 \rightarrow Wb$, (41%) $u_2 \rightarrow Zq_1$, (1.5%) $u_2 \rightarrow Wq'^{d_3}$, (2.5%) $u_2 \rightarrow W\chi_2^{u_3}$, (2.%)
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Solid prediction



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Solid prediction

Collider Phenomenology

■ Tevatron

- Searches for new quarks
- Interesting possibilities for Higgs physics

$$\sigma(pp\bar{p} \rightarrow q'\bar{q}' \rightarrow 2H2j) \sim \sigma(gg \rightarrow H)$$

■ LHC

- GB KK modes discovery up to ~ 5 TeV
- Exotic decay channels

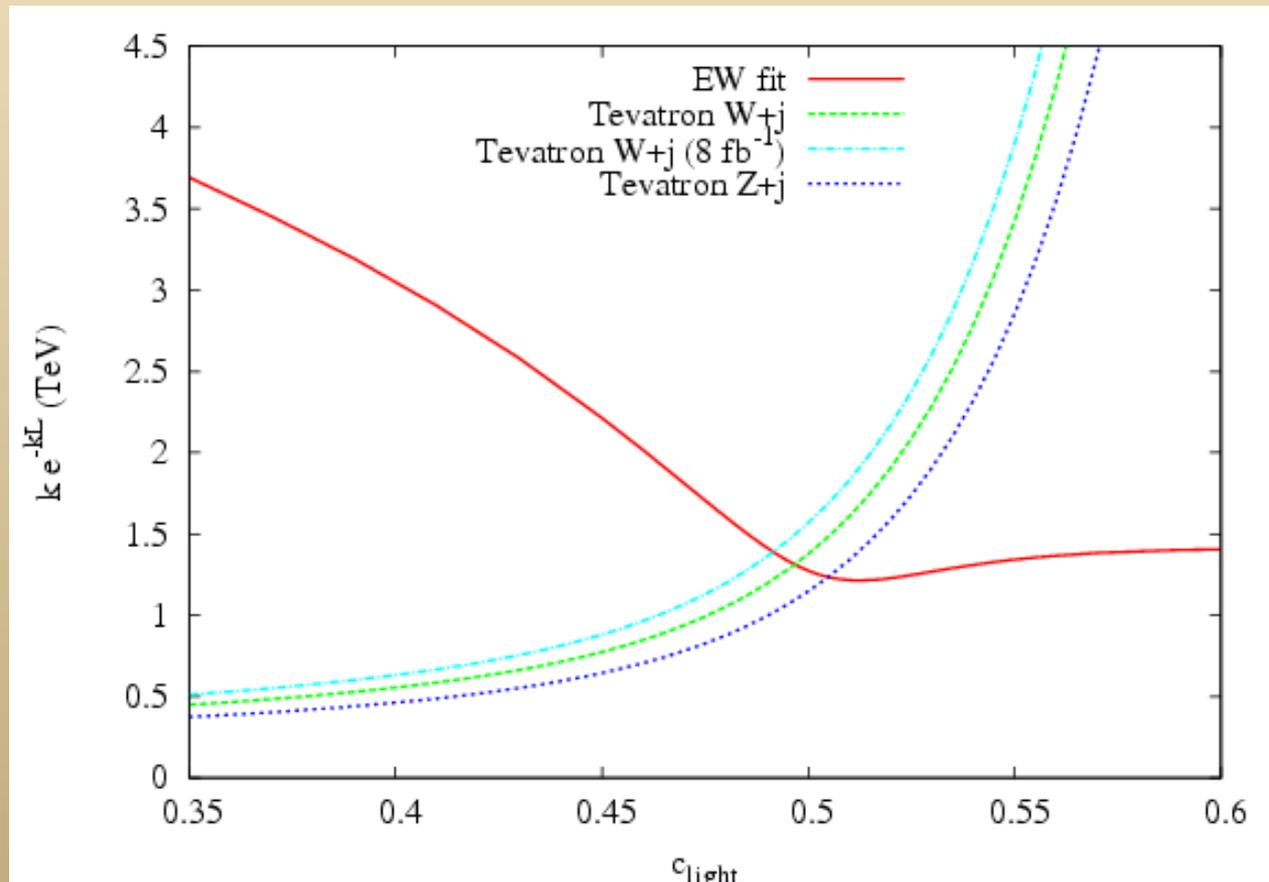
$$pp \rightarrow q'\bar{q}' \rightarrow W^+W^-t\bar{t} \rightarrow W^+W^+W^-W^-b\bar{b}$$

- t' discovery up to ~ 1 TeV
- Exciting Higgs physics

Tevatron searches

New quark searches

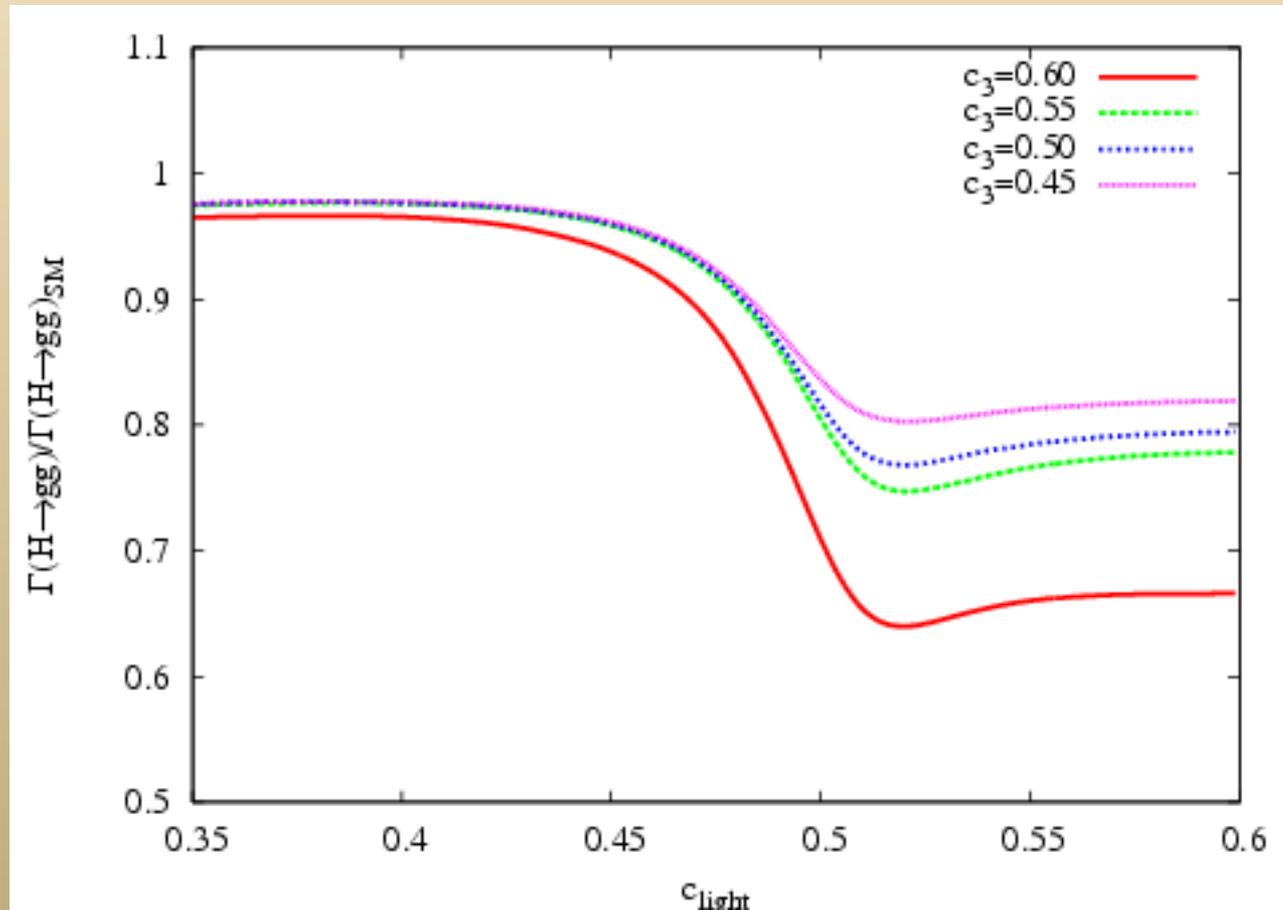
$$m_q \geq \begin{cases} 325 \text{ (410) GeV,} & W + j \text{ with } 0.76 \text{ (projected 8) fb}^{-1} \\ 300 \text{ GeV,} & Z + j \end{cases}$$



Tevatron searches

Higgs physics

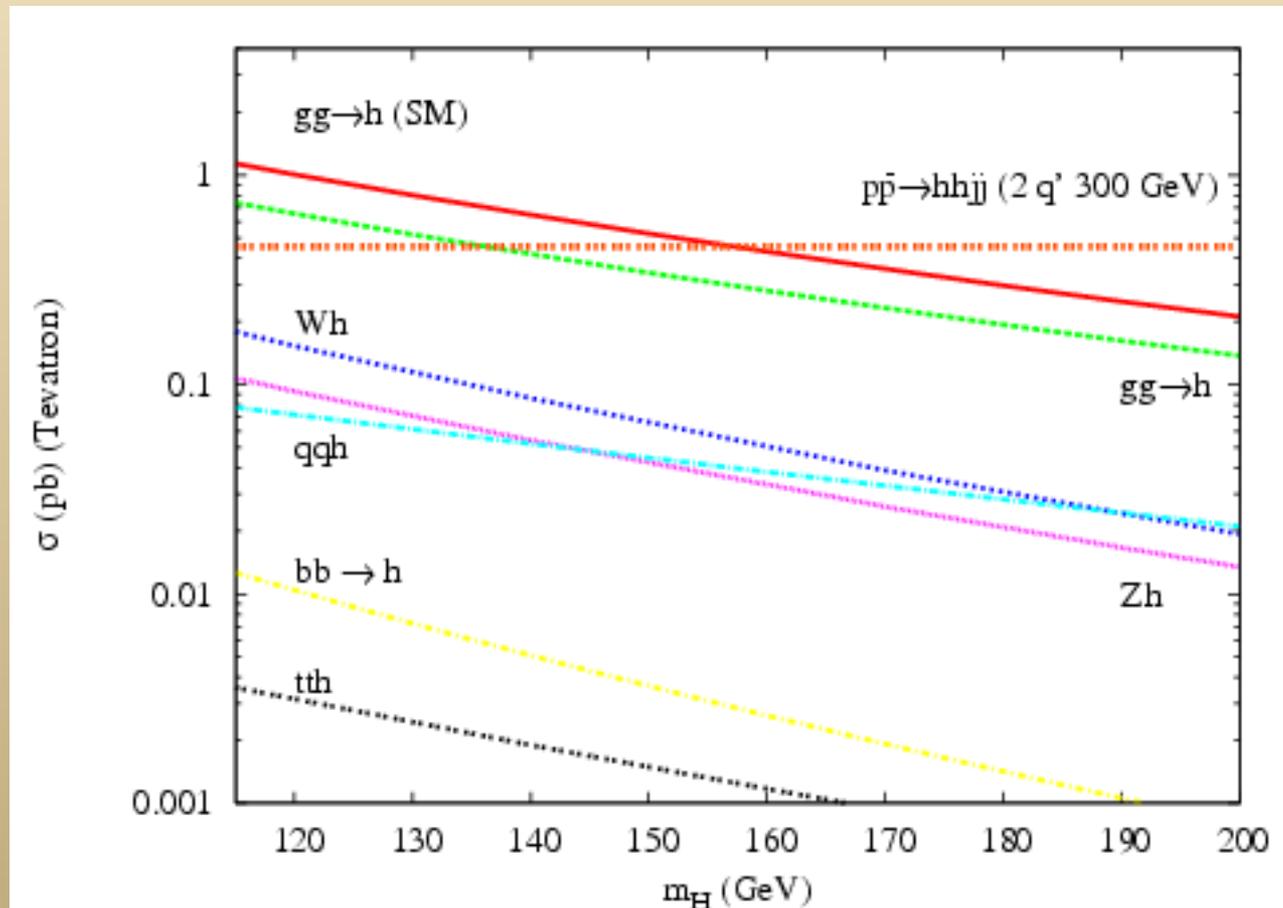
Top mixing with KK modes \Rightarrow reduced top Yukawa



Tevatron searches

Higgs physics

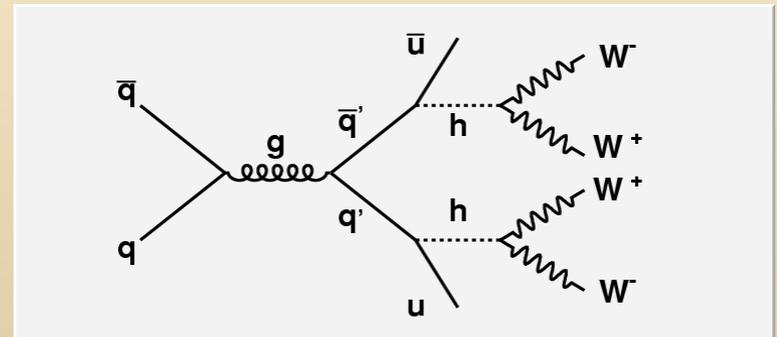
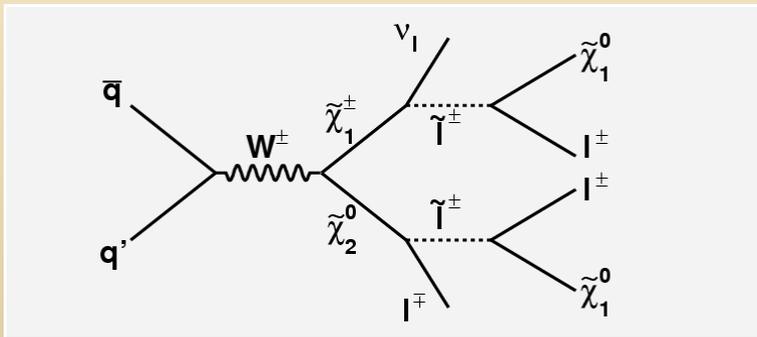
But new production mechanism mediated by q'



Tevatron searches

LS dileptons: SUSY vs q' ($M_H=170$ GeV)

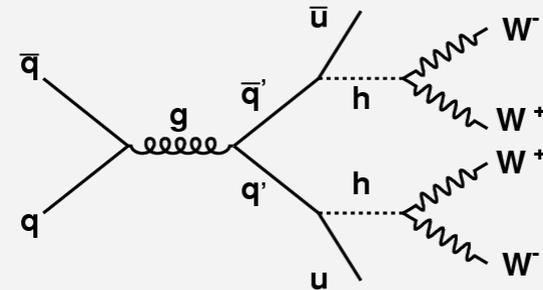
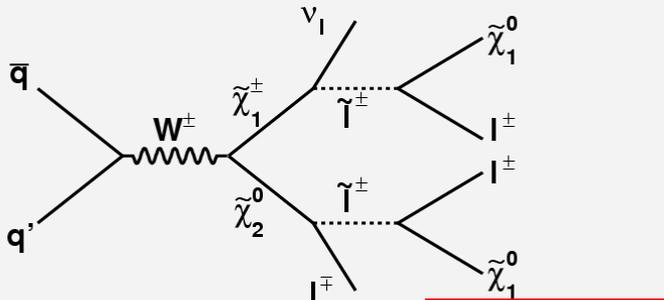
D0 Note 5126-CONF (1fb^{-1})



Tevatron searches

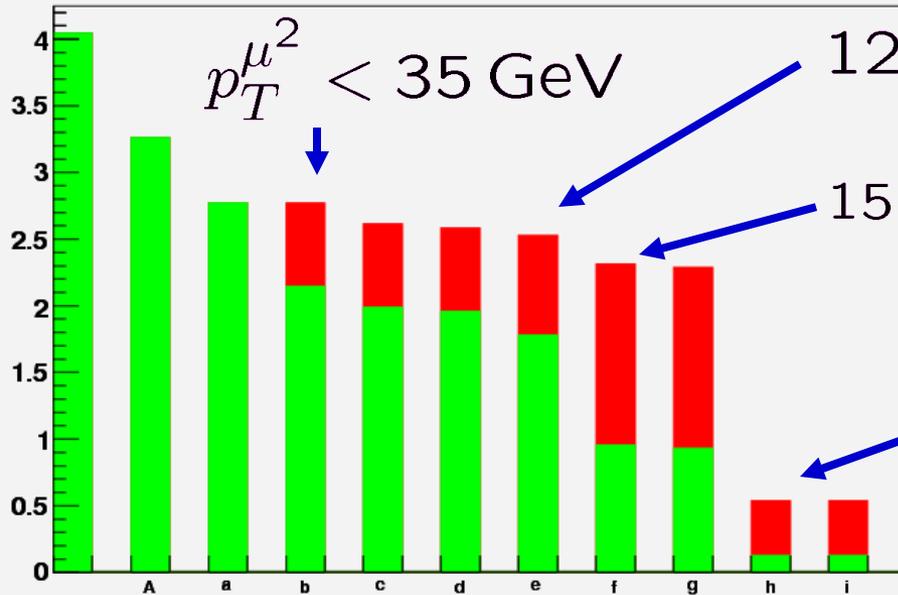
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PRELIMINARY

events/ fb^{-1}



$12 \text{ GeV} < M_{\mu\mu} < 110 \text{ GeV}$

$15 \text{ GeV} < M_T(E_T, p_T^{\mu^2}) < 65 \text{ GeV}$

$\text{Sig}(E_T)$

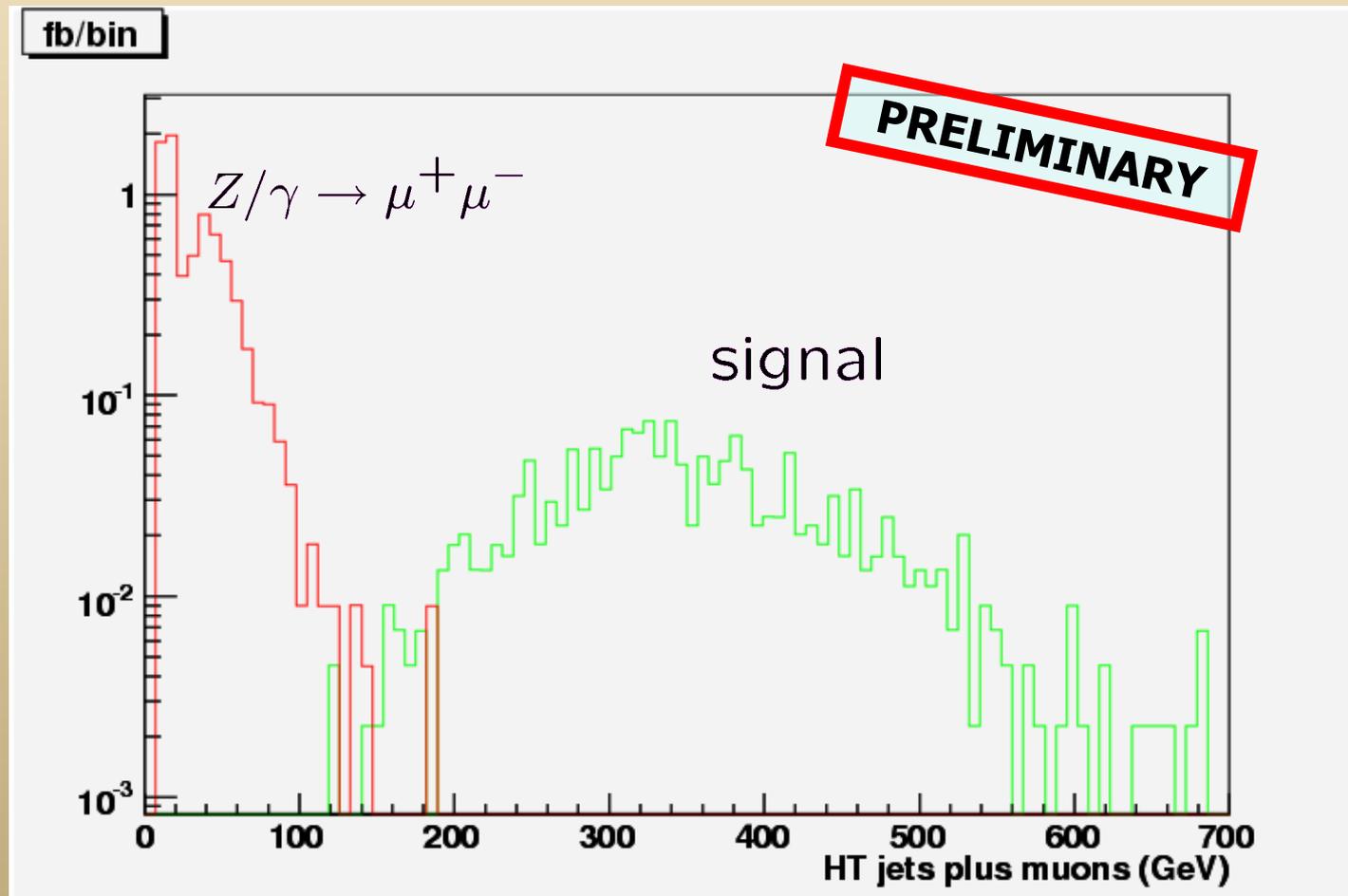
Simulations using
MG/ME+Pythia+PGS4

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Alternative cuts: signal is harder than backgrounds

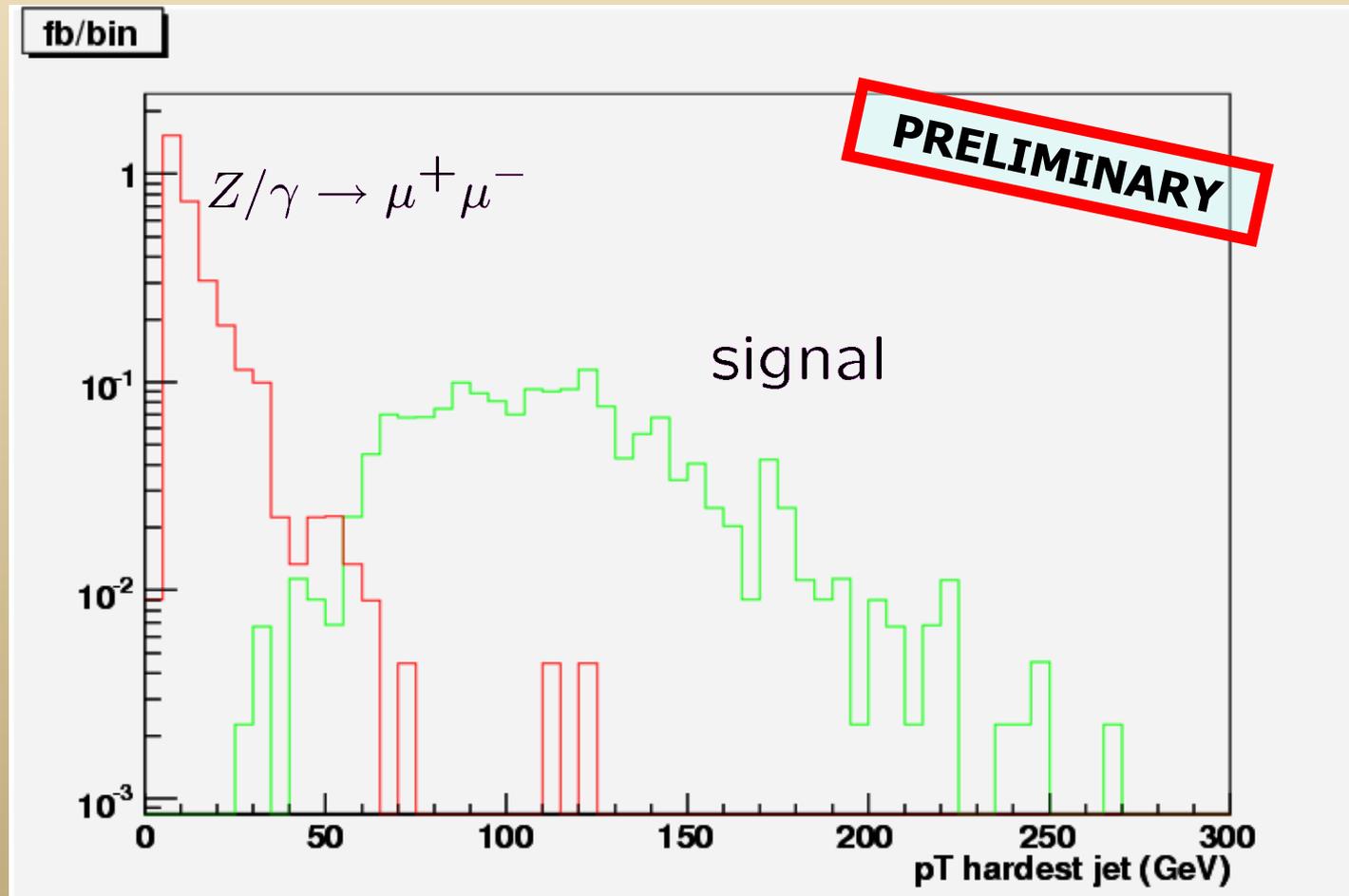


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Tevatron searches

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D0 Note 5126-CONF (1fb⁻¹)

Main backgrounds considered

QCD, WZ , ZZ , $Z/\gamma + \text{jet}(s)$ and $W + \text{jet}(s)$

Result of the analysis for 1 fb⁻¹

	sum	data
selection \mathcal{B}	14922±981	15234
$M_{\mu^\pm\mu^\mp} \in [25 - 65]$ GeV/c ² (a)	3479±232	3569
$p_T^{\mu^2} < 35$ GeV/c (b)	3479±232	3358
$p_T^{\mu^2} > 8$ GeV/c (c)	8.9±1.8	10
$p_T^{\mu^1} > 13$ GeV/c (d)	6.5±1.4	6
$M_{\mu^\pm\mu^\pm} \in [12 - 110]$ GeV/c ² (e)	4.9±1.2	2
$M_T(\cancel{E}_T, p_T^{\mu^2}) \in [15 - 65]$ GeV/c ² (f)	2.9±0.8	2
$\cancel{E}_T > 10$ GeV (g)	2.3±0.7	1
Sig(\cancel{E}_T) > 12 GeV ^{1/2} (h)	1.7±0.6	1
$\cancel{E}_T \times p_T^{\mu^2} > 160$ GeV ² /c (i)	1.1±0.4	1

Our signal, assuming improved analysis is 2.5 events per fb⁻¹

In the absence of signal ~2 fb⁻¹ could set a 95% C.L. bound

LHC searches

- $Z^{(1)}, \gamma^{(1)}$ accessible at the LHC in Drell-Yan up to

$$M_{Z^{(1)}, \gamma^{(1)}} \sim 6 \text{ TeV} \quad \text{Ledroit, Moreau, Morel ph/0703262}$$

- KK gluons require careful treatment of *top jets* and improved fake jet efficiencies but

$$M_{g^{(1)}} \sim 4 - 5 \text{ TeV}, 100 \text{ fb}^{-1} \quad \text{Agashe, Belyaev, Krupovnickas, Perez, Virzi ph/0612015}$$

Lillie, Randall, Wang ph/0701166

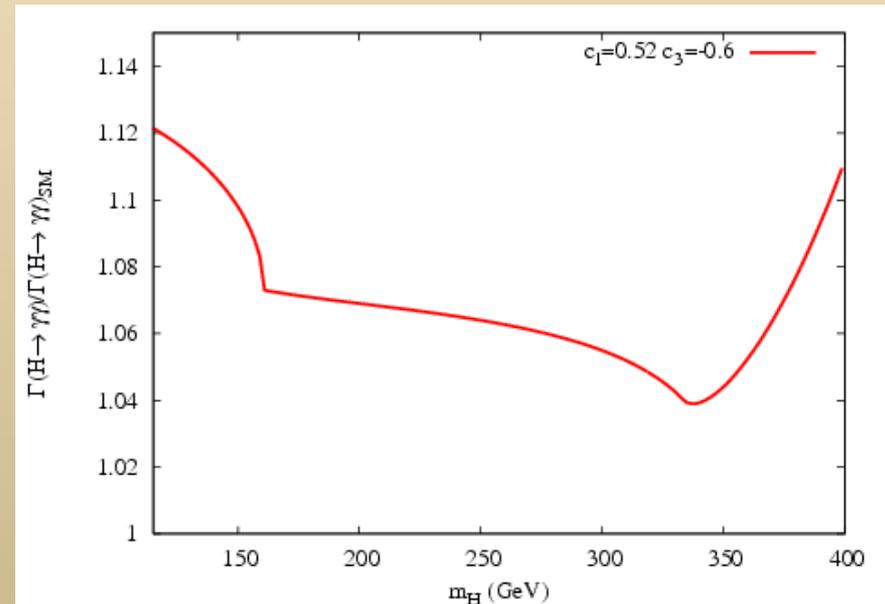
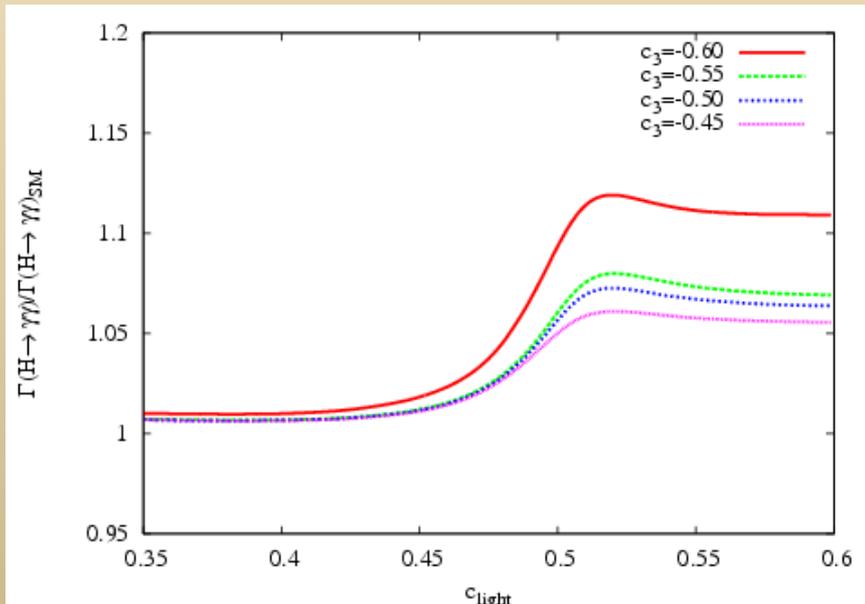
- Exotic decay channels with 4 or more W 's may allow early discovery of q' with $M_{q'}=500 \text{ GeV}$ with 10 fb^{-1}

Dennis, Ünel, Servant, Tseng ph/0701158

- t' with masses $\sim 1 \text{ TeV}$ can be discovered at the LHC

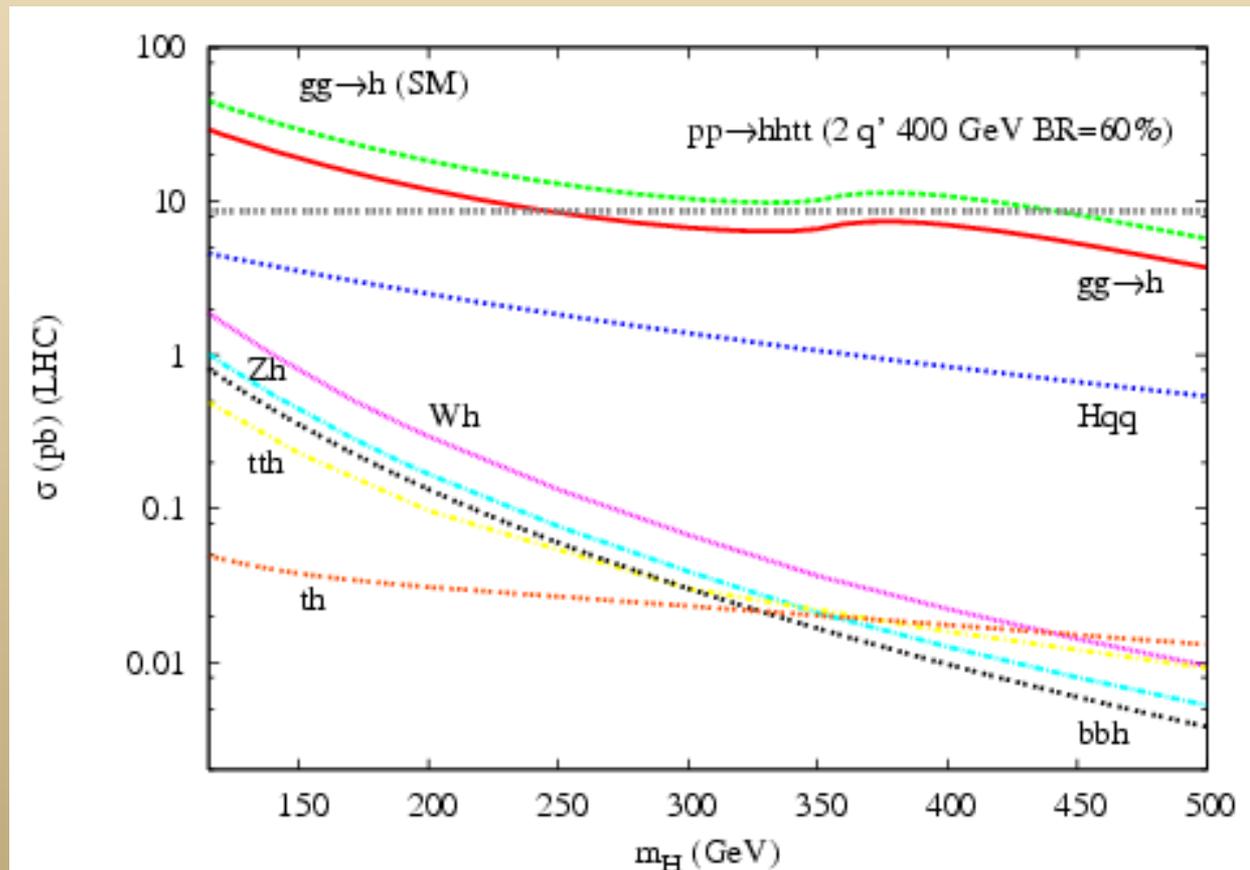
LHC Higgs physics

- Gluon fusion production reduced up to a factor 0.65
- $H \rightarrow \gamma\gamma$ slightly enhanced (up to 1.12 for light Higgs)



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- But new production mechanism mediated by q'
- New possibilities for early discovery

Aguilar-Saavedra JHEP (06)

$$pp \rightarrow T\bar{T} \rightarrow W^+ b H \bar{t} / H t W^- \bar{b} \rightarrow W^+ b W^- \bar{b} H$$

$$pp \rightarrow T\bar{T} \rightarrow H t H \bar{t} \rightarrow W^+ b W^- \bar{b} H H$$

5σ discovery for $m_T=500$ GeV and $m_H=115$ GeV with 8 fb^{-1}

T is a vector-like singlet $\Rightarrow BR(T \rightarrow Ht) \approx 25\%$

But need dedicated analysis!

Conclusions

- Custodial protection of T and Zbb: new opportunities for Randall-Sundrum models
- Quantum corrections calculable and relevant:
 - Fix some features of the model (t_L, b_L near IR brane)
 - Help ease the tension after new M_{top}, M_W measurements
- Very rich, interesting and new phenomenology
 - Light ($\sim 2 - 3$ TeV) gauge boson KK modes
 - Light ($\sim 0.4 - 1$ TeV) quarks, some with exotic charges
 - Spectacular signatures: multigauge bosons, LS leptons, ...
 - Important effects in Higgs physics