

# Giant $K$ factors

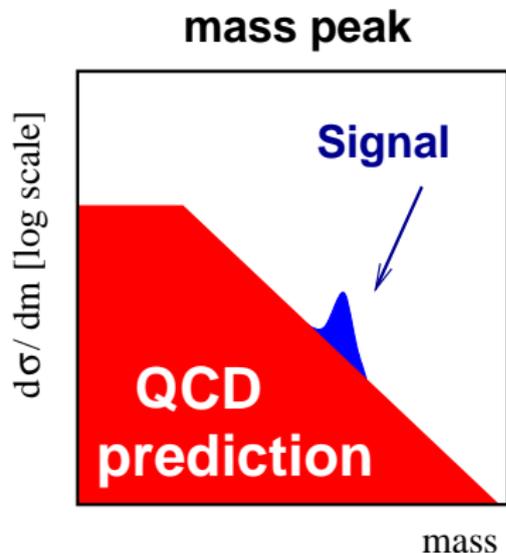
Gavin Salam

CERN, Princeton & LPTHE/CNRS (Paris)

Work performed with Mathieu Rubin and Sebastian Sapeta, [arXiv:1006.2144](https://arxiv.org/abs/1006.2144)

Fermilab Theory Seminar

19 May 2011

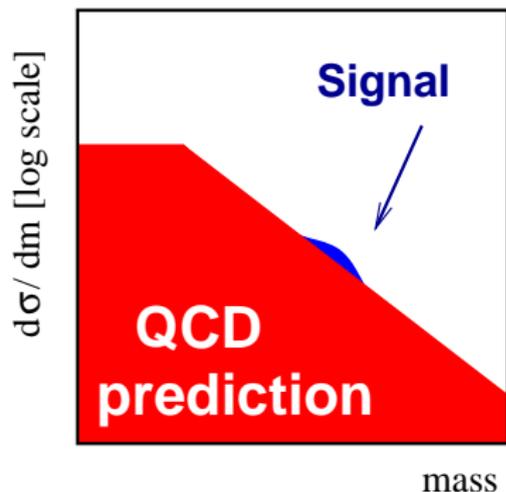


New resonance (e.g.  $Z'$ ) where you see all decay products and reconstruct an invariant mass

QCD may:

- ▶ swamp signal
- ▶ smear signal

leptonic case easy; hadronic case harder

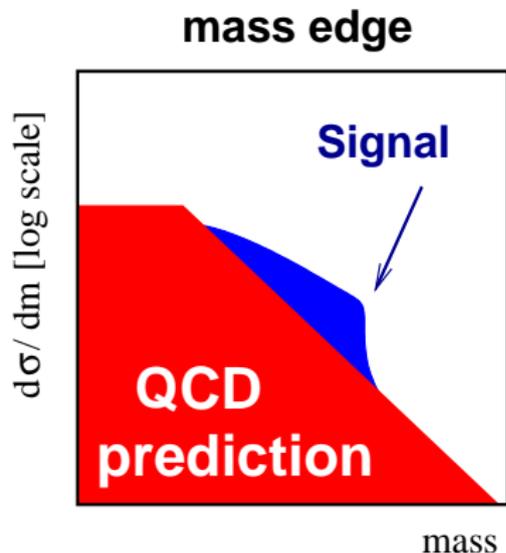


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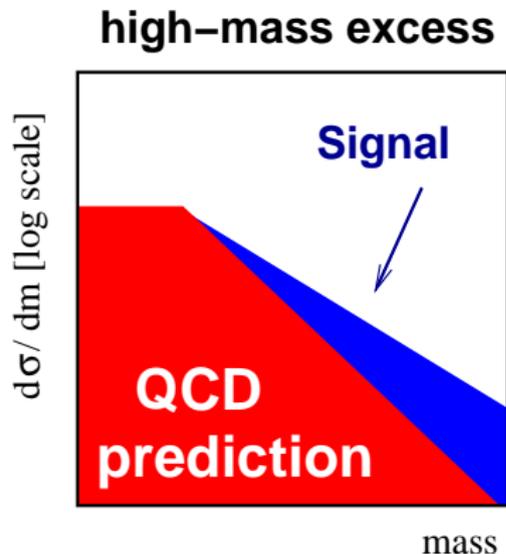


New resonance (e.g. R-parity conserving SUSY), where undetected new stable particle escapes detection.

Reconstruct only *part* of an invariant mass  
→ kinematic edge.

QCD may:

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- ▶ smear signal

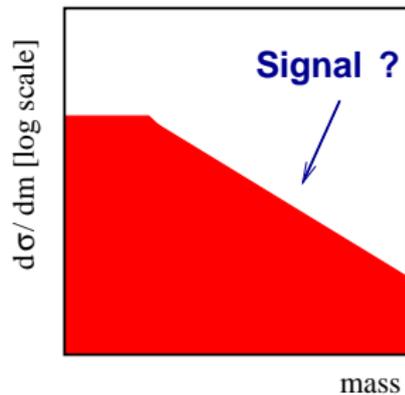
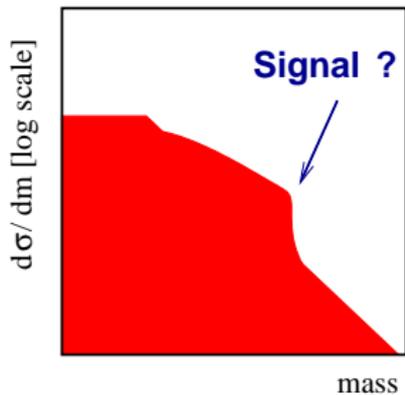
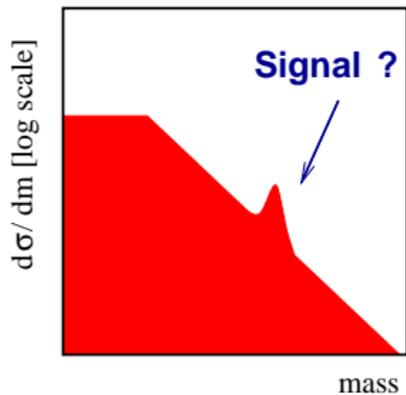


Unreconstructed SUSY cascade. Study *effective* mass (sum of all transverse momenta).

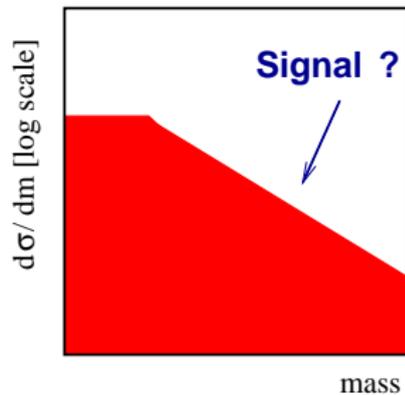
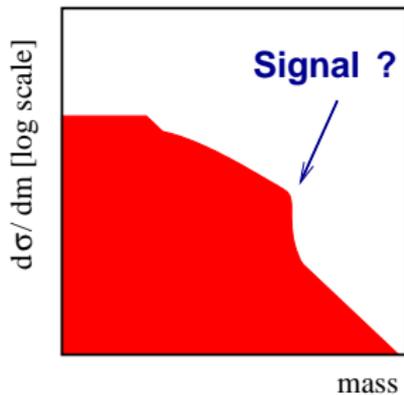
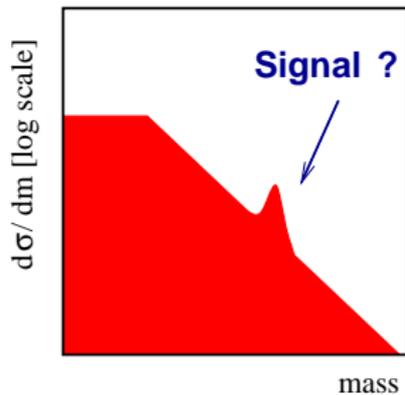
Broad excess at high mass scales.

Knowledge of backgrounds is crucial in declaring discovery.

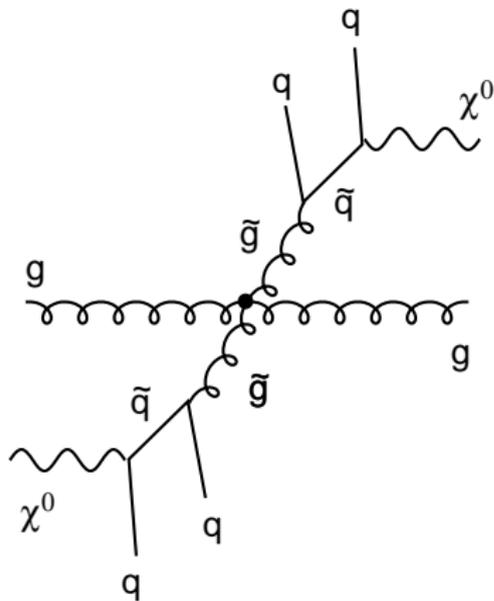
QCD is *one way* of getting handle on background.

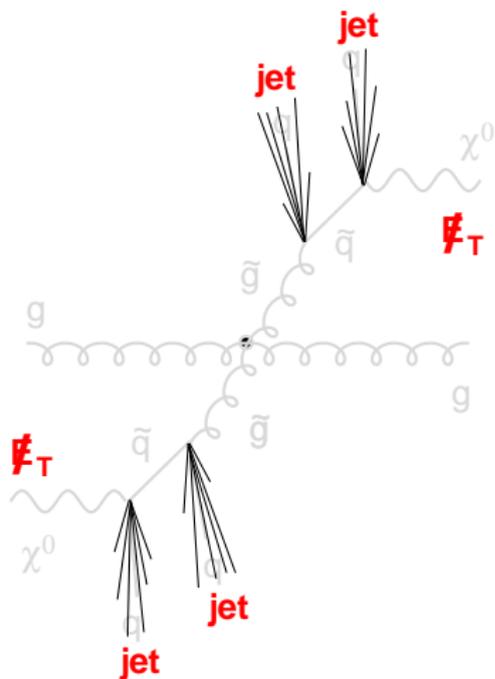


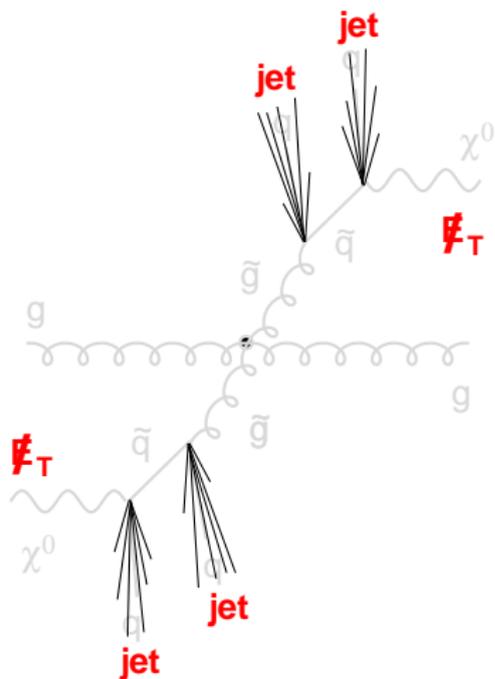
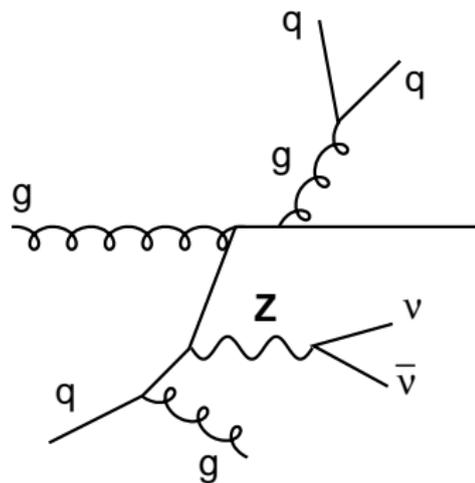
THIS  
TALK



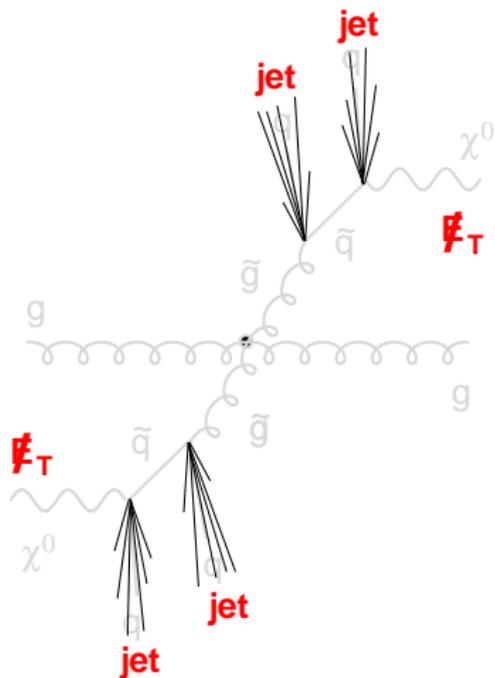
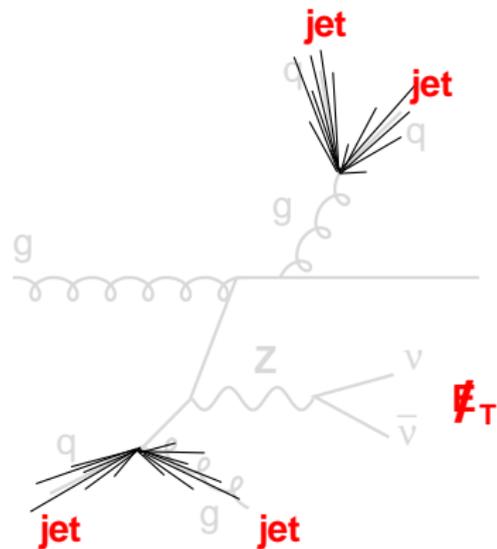
**THIS  
TALK**

Signal

Signal

SignalBackground

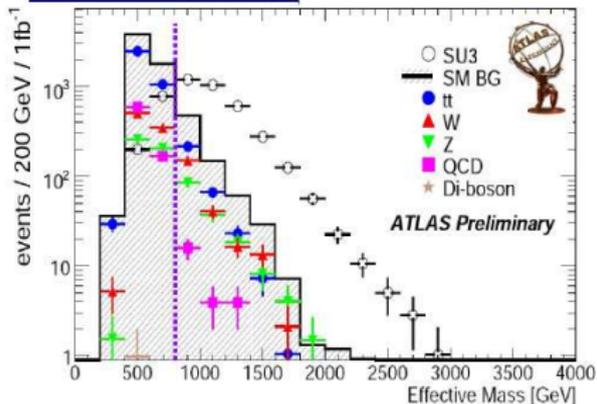
## SUSY example: gluino pair production

SignalBackground

## Atlas selection [all hadronic]

- no lepton
- MET > 100 GeV
- 1<sup>st</sup>, 2<sup>nd</sup> jet > 100 GeV
- 3<sup>rd</sup>, 4<sup>th</sup> jet > 50 GeV
- MET / m<sub>eff</sub> > 20%

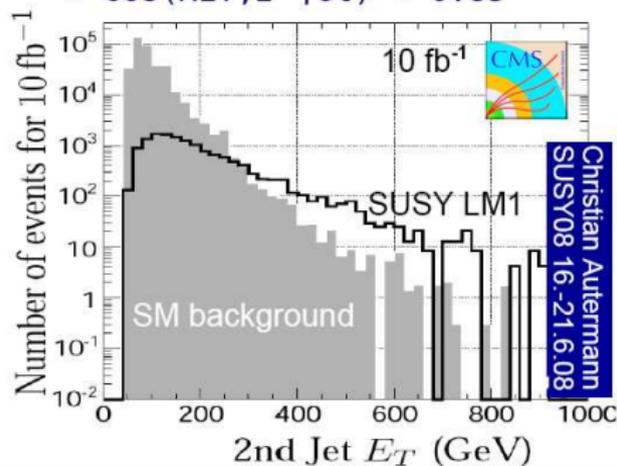
Christian Autermann  
SUSY08 16.-21.6.08  
4



## CMS selection [leptonic incl.]

(optimized for 10fb<sup>-1</sup>, using genetic algorithm)

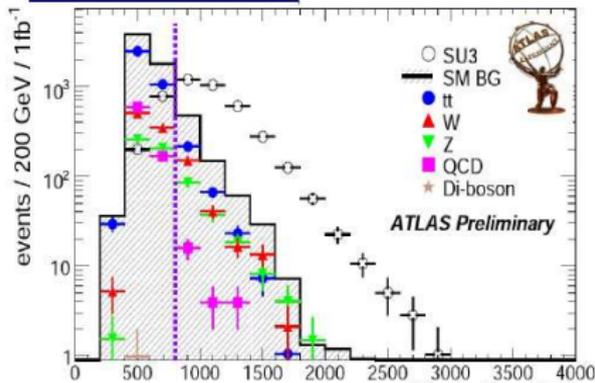
- 1 muon pT > 30 GeV
- MET > 130 GeV
- 1<sup>st</sup>, 2<sup>nd</sup> jet > 440 GeV
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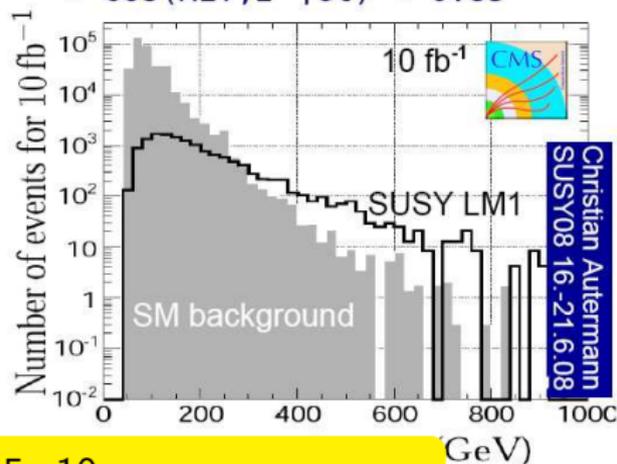
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SUSY ≈ factor 5–10 excess

# How accurate is perturbative QCD?

$$\sigma = c_0 + c_1\alpha_s + c_2\alpha_s^2 + \dots$$

$$\alpha_s \simeq 0.1$$

That implies LO QCD (just  $c_0$ )  
should be accurate to within 10%

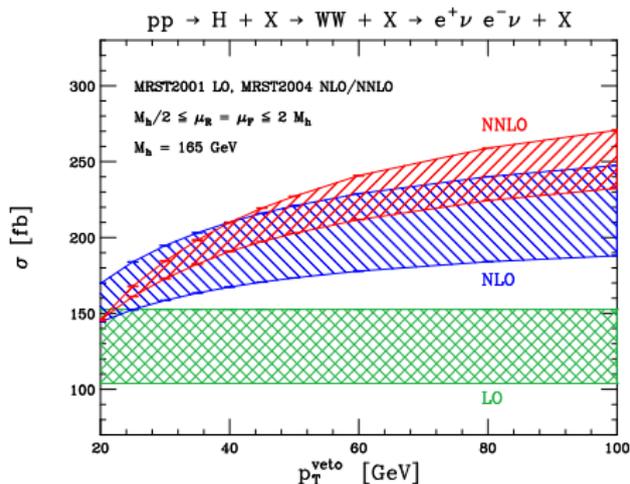
It isn't

Rules of thumb:

LO good to within factor of 2

NLO good to within scale  
uncertainty

This talk is about an example where these rules fail spectacularly,  
the lessons we learn, and the solutions we can apply.



Anastasiou, Melnikov & Petriello '04  
 Anastasiou, Dissertori & Stöckli '07

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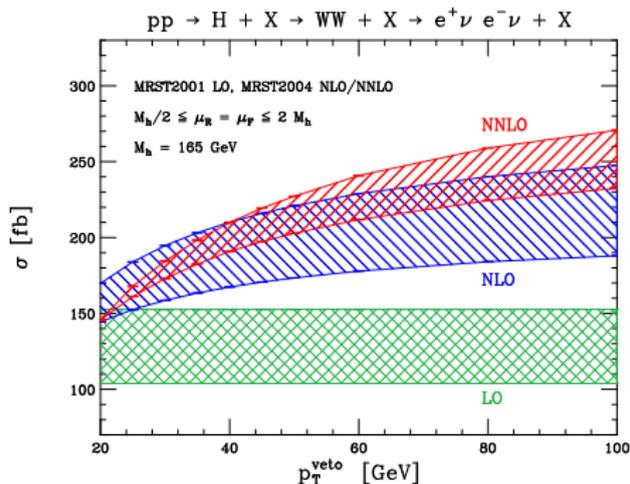
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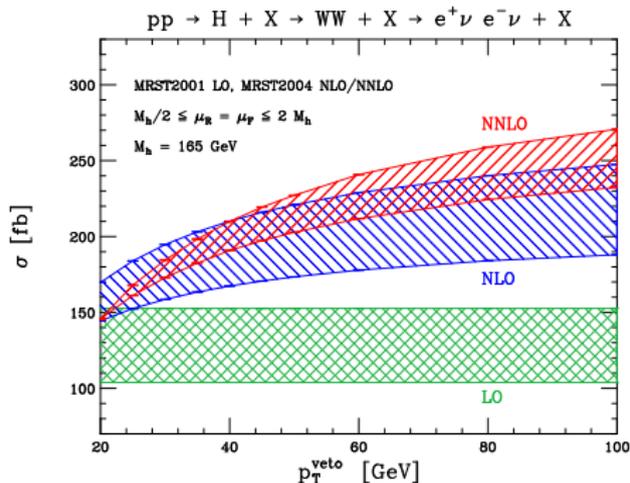
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We don't always have NLO for the background (e.g.  $Z+4$  jets, a  $2 \rightarrow 5$  process).

Though amazing recent progress

$2 \rightarrow 4$ : Blackhat, Rocket, Helac-NLO, BDDP

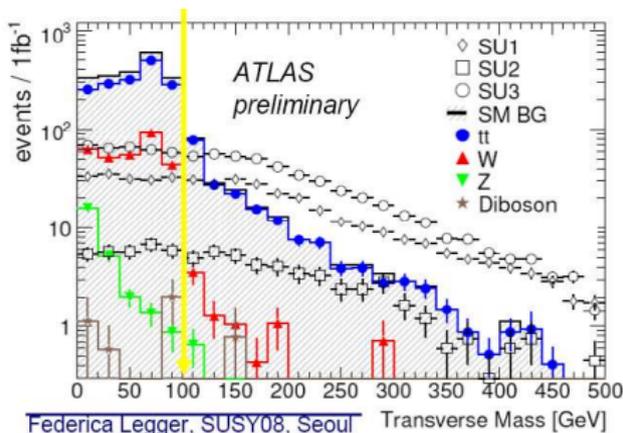
$2 \rightarrow 5$  ( $W+4j$ ): Blackhat

Must then rely on LO (matched with parton showers). How does one verify it?

Common “data-driven” procedure:

[roughly]

- ▶ Get control sample at low  $p_t$
- ▶ SUSY should be small(er) contamination there
- ▶ Once validated, trust LO prediction at high- $p_t$



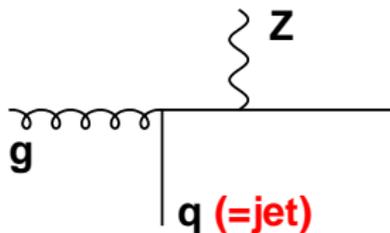
A conservative QCD theory point of view:

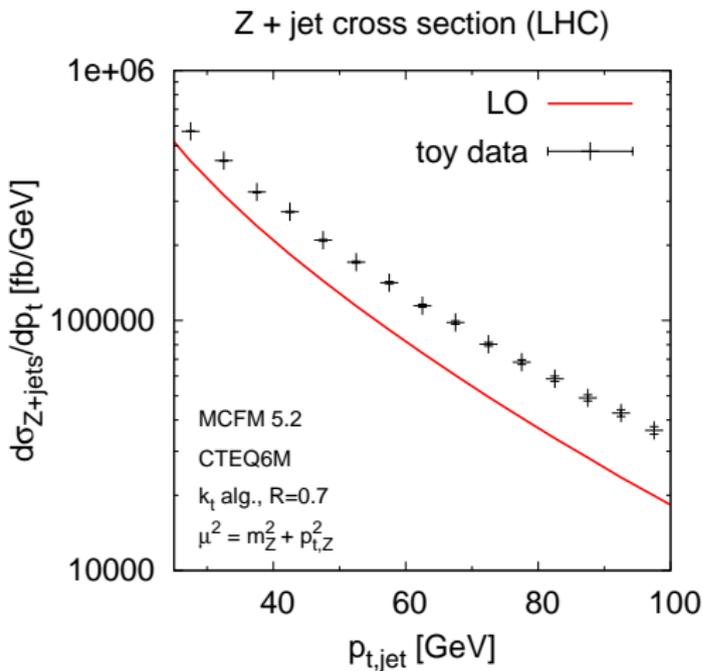
It's hard to be sure: since we can't (yet) calculate  $Z+4$  jets beyond LO.

But we would tend to think it is safe, as long as control data are within usual factor of two of LO prediction

Illustrate issues with toy example:  $Z$ +jet production

- ▶ It's known to NLO and a candidate for "first"  $2 \rightarrow 2$  NNLO  
 $\sim e^+e^- \rightarrow \gamma^*/Z \rightarrow 3$  jets, NNLO: Gehrman et al '08, Weinzierl '08
- ▶ But let's pretend we only know it to LO, and look at the  $p_t$  distribution of the hardest jet (no other cuts — keep it simple)



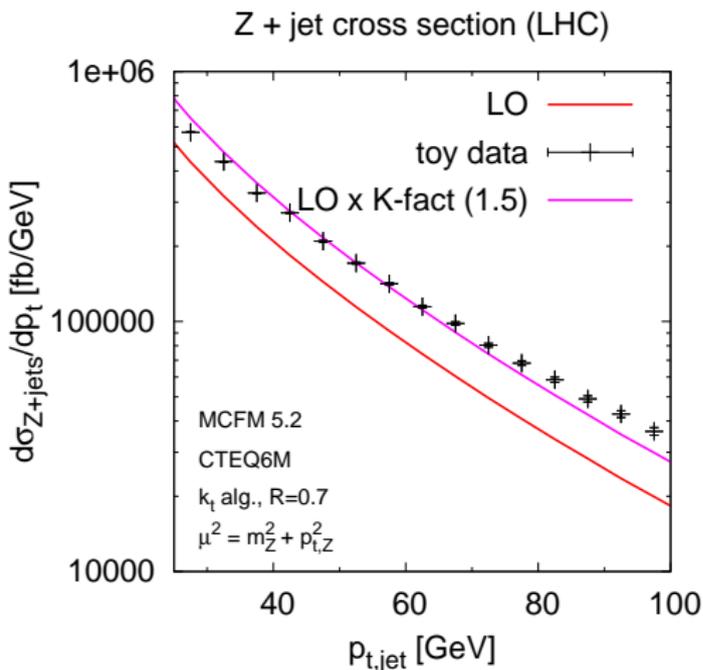
stage 1: get control sampleCheck LO v. data at low  $p_t$ 

- ▶ normalisation off by factor 1.5  
(consistent with expectations)

So renormalise LO by K-factor

- ▶ shape OKish

Don't be too fussy: SUSY  
could bias higher  $p_t$



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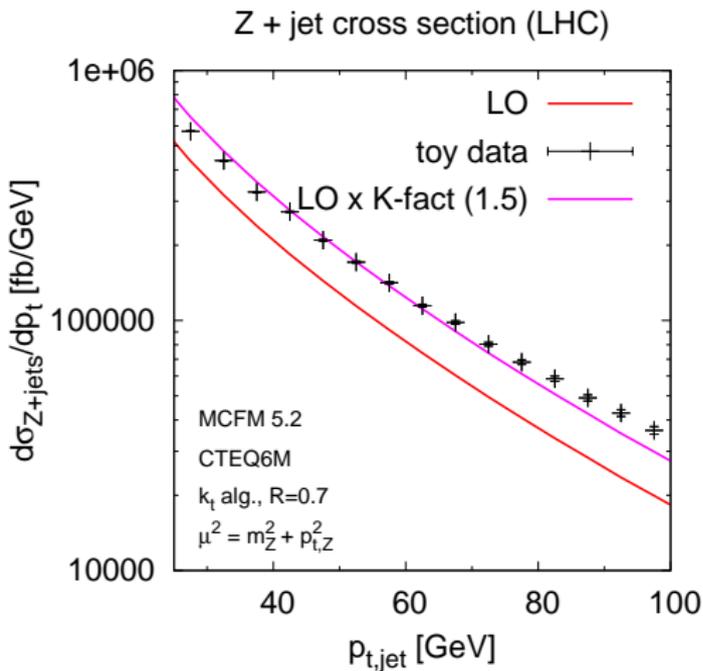
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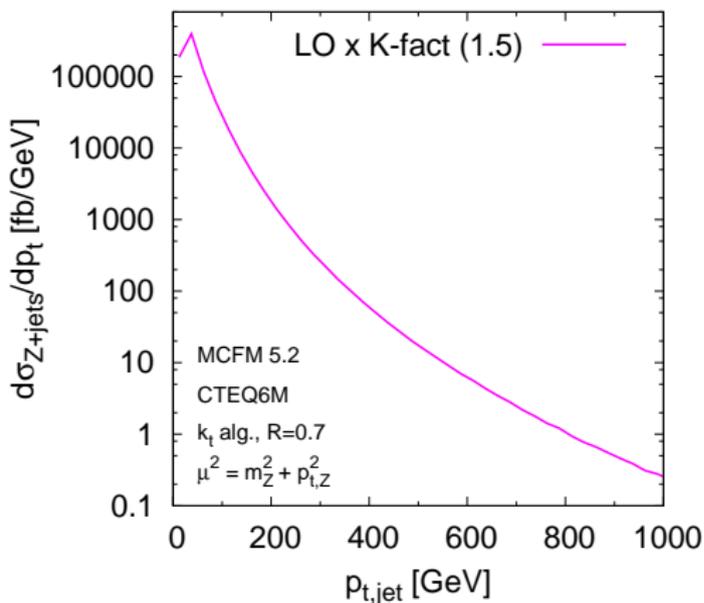
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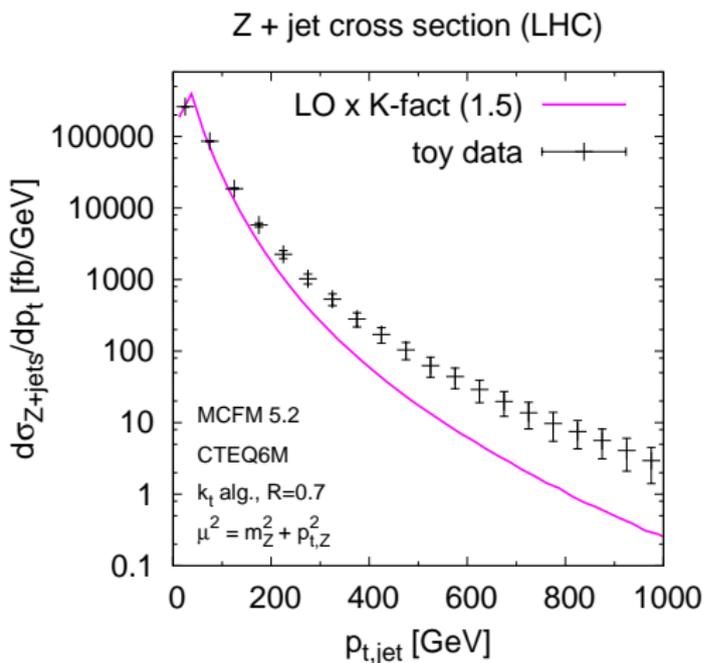
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Z + jet cross section (LHC)

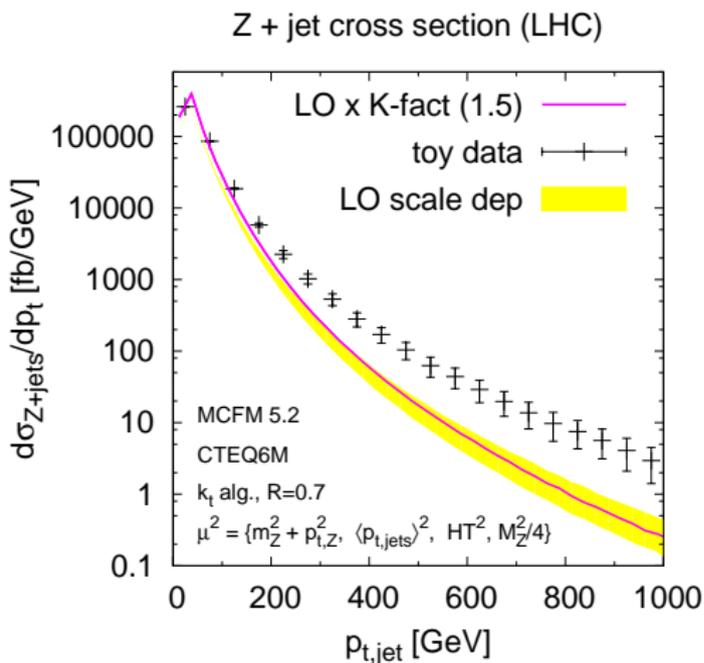
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- ▶ good agreement at low  $p_t$ , by construction
- ▶ excess of factor  $\sim 10$  at high  $p_t$
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[NB: not always done except e.g. Alwall et al. 0706.2569]  
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Is it:

- ▶ QCD + extra signal?
- ▶ just QCD? But then where does a  $K$ -factor of 10 come from?

Here it's just a toy illustration. Later this year it may be for real:

- ▶ Do Nature / Science / PRL accept the paper?

**Discovery of New Physics at the TeV scale**

*We report a  $5.7\sigma$  excess in MET + jets production that is consistent with a signal of new physics ...*

- ▶ Do we proceed immediately with a linear collider?  
It'll take 10–15 years to build; the sooner we start the better
- ▶ At what energy? It would be a shame to be locked in to the wrong energy...

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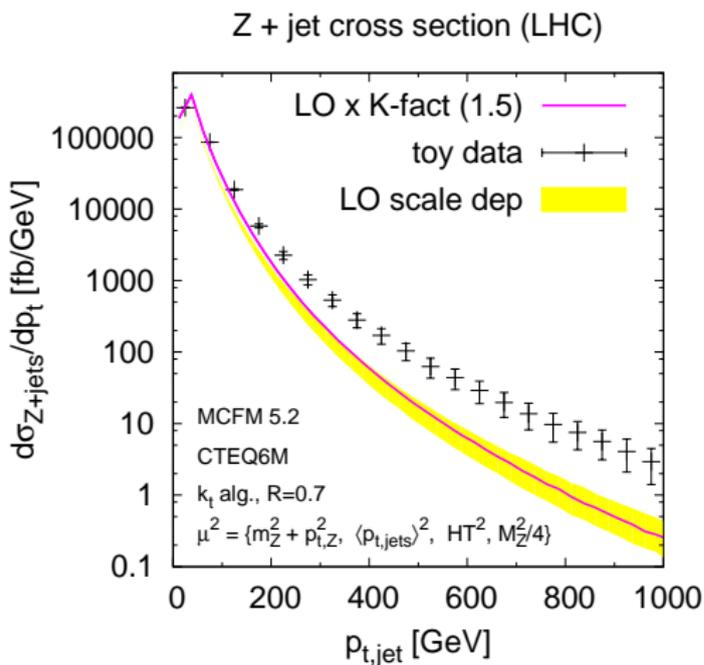
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Unlike for SUSY multi-jet searches, in the Z+jet case we do have NLO.

Once NLO is included the excess disappears

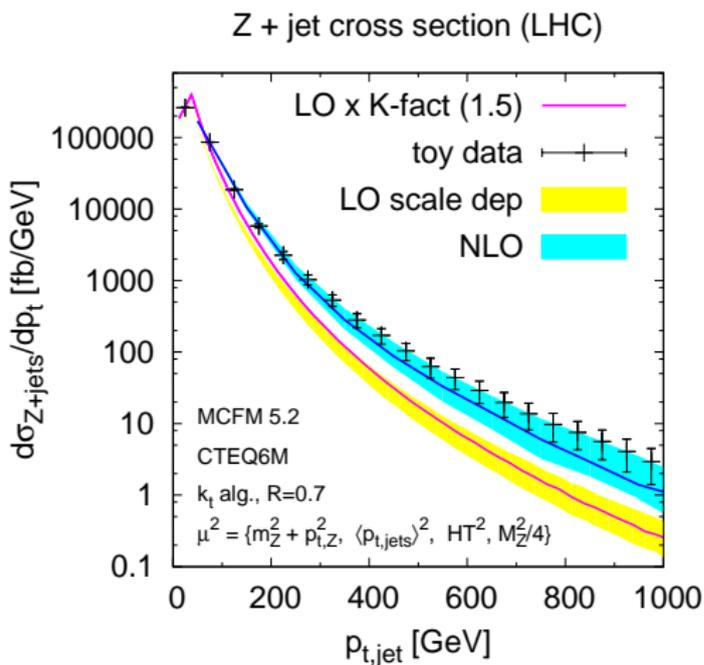
The “toy data” were just the upper edge of the NLO band

Example based on background work for Butterworth, Davison, Rubin & GPS '08

Related observations also by Bauer & Lange '09; Denner, Dittmaier, Kasprzik & Muck '09

*Hold on a second: how does QCD give a K-factor  $\mathcal{O}(5 - 10)$ ?*

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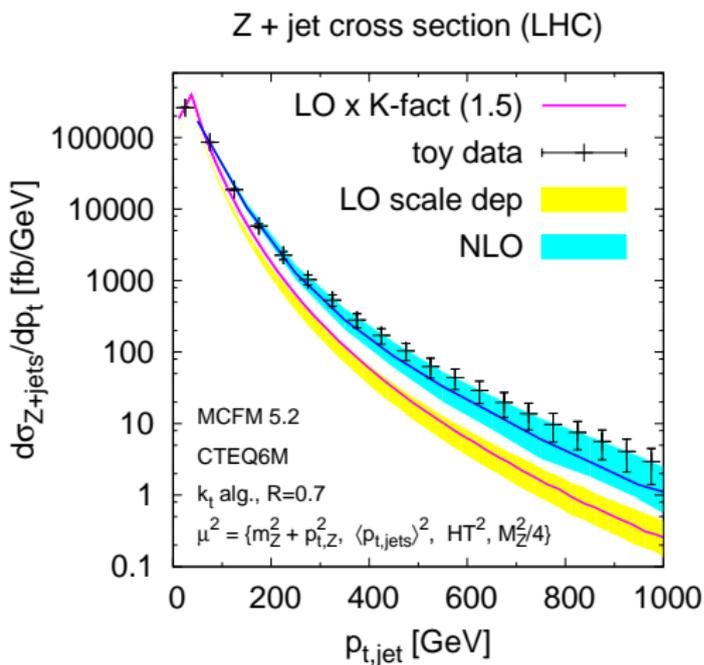
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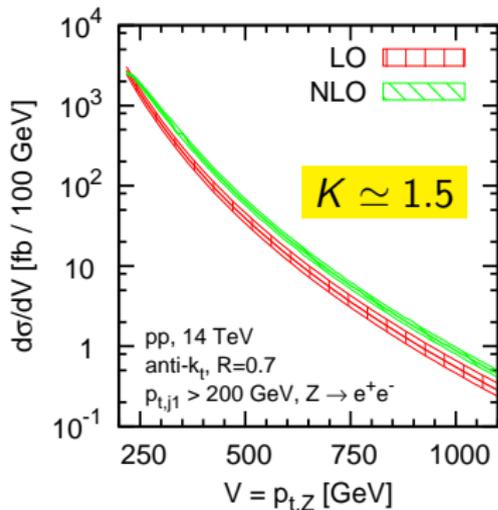
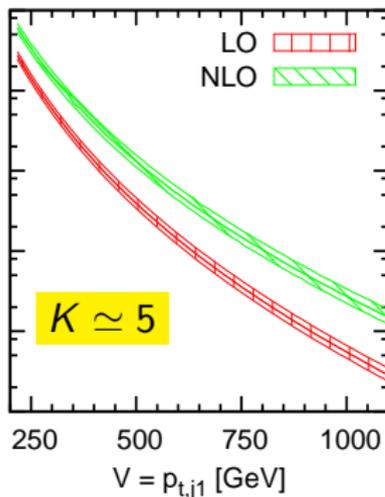
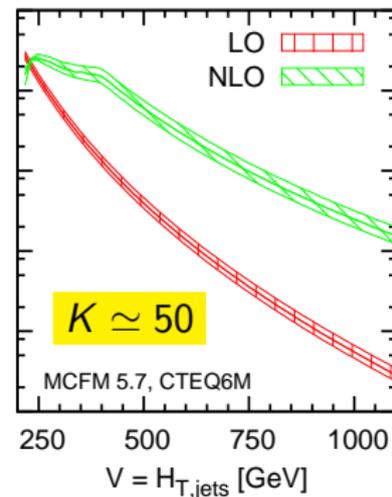
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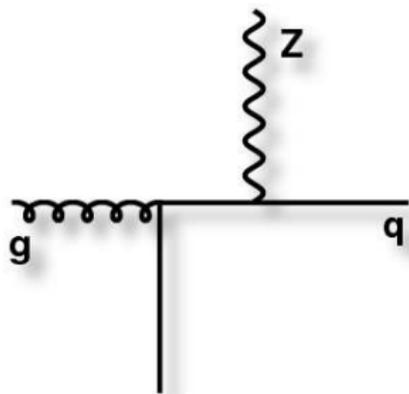
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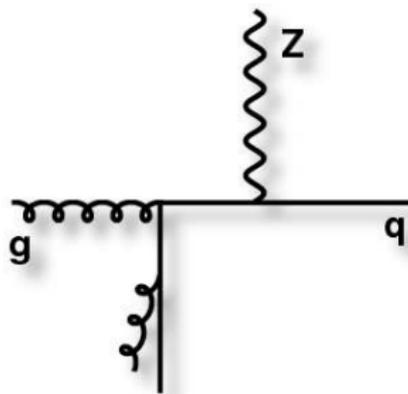
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$p_t$  of Z-boson $p_t$  of jet 1 $H_{T,jets} = \sum_{jets} p_{t,j}$ 

**“Giant K-factors”**

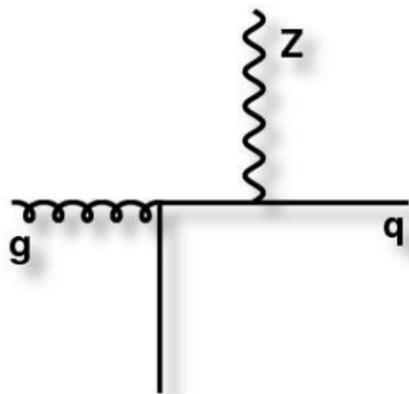
Leading Order

$$\alpha_s \alpha_{EW}$$

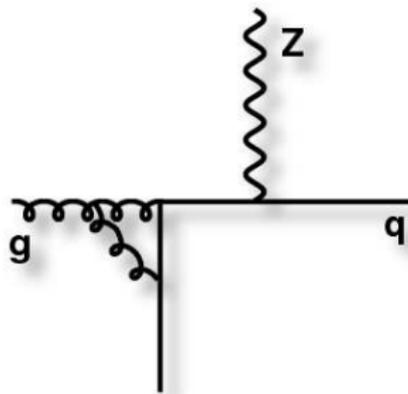
Next-to-Leading Order

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LHC probes scales  $\gg$  EW scale,  $\sqrt{s} \gg M_Z$ . EW bosons are **light**.  
**New logarithmically enhanced topologies appear.**

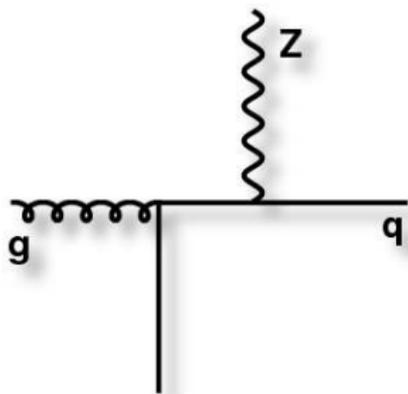
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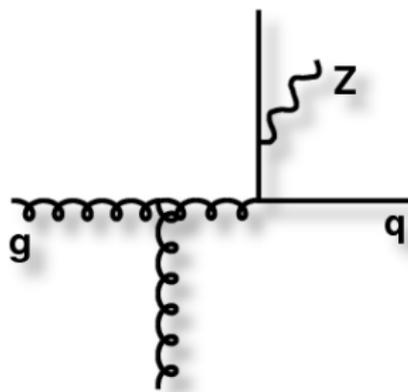
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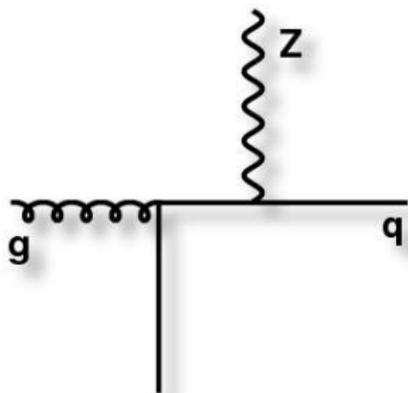
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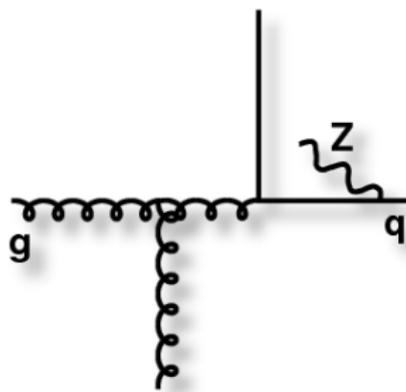
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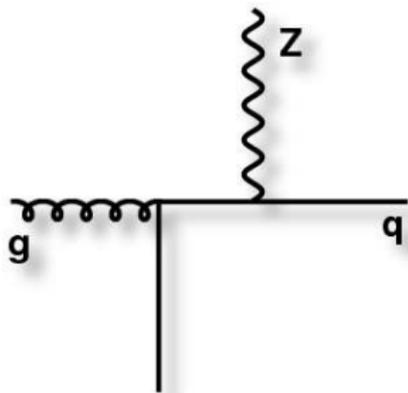
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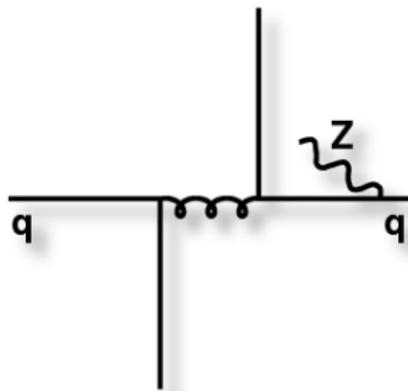
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NLO driven in part by  $qq$  parton luminosity: large at  $pp$  colliders

Is this example not a little contrived?  
After all, experiments would surely notice unexpected event topology such as that here.

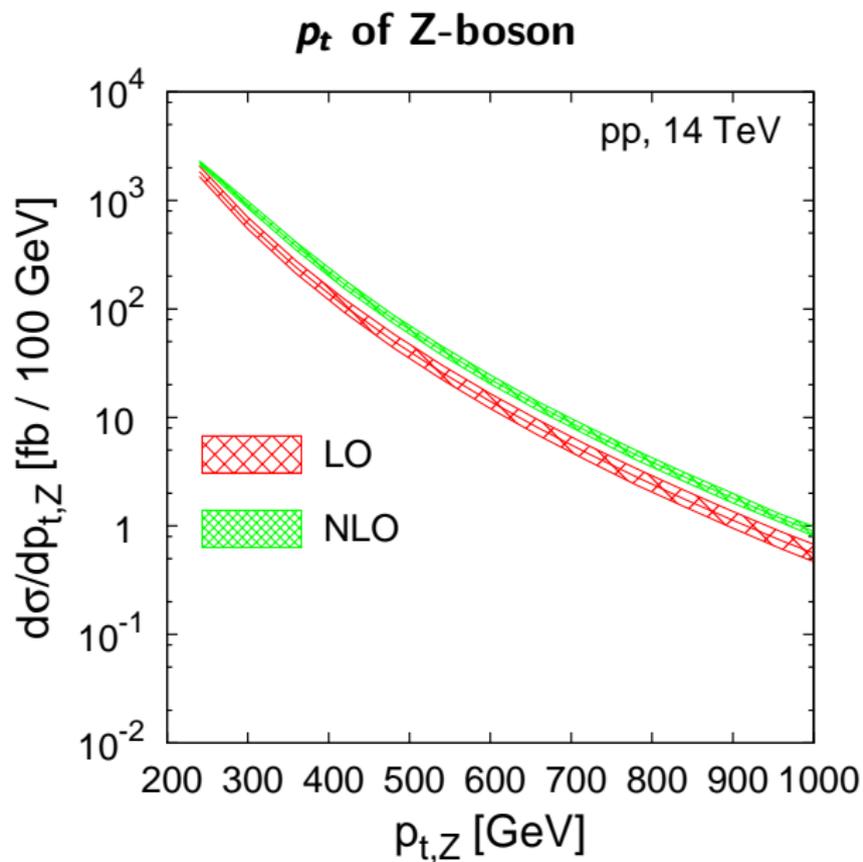
We actually first saw the problem in a more complex process:  $Wb\bar{b}$  as a background to boosted Higgs searches (with “wrong” cuts). *The more complicated the process, the trickier the diagnosis of the problem.*

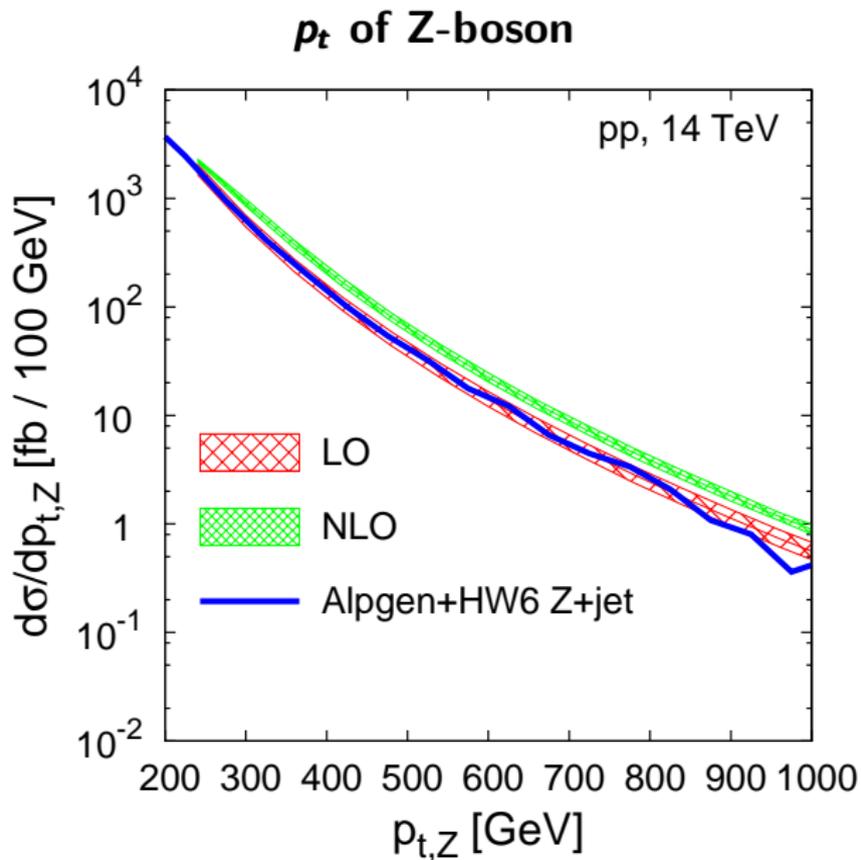
It's enough to get this wrong once, leading to “unwarranted” press-releases and major subsequent embarrassment.

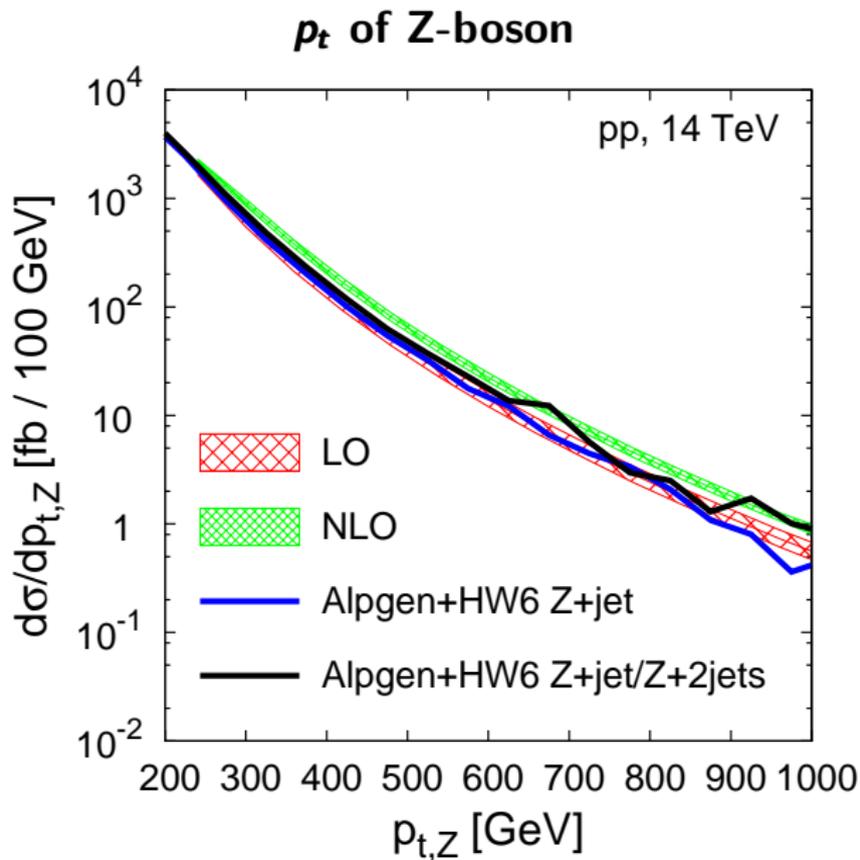
In day-to-day work the experiments don't just use LO.

Instead, they

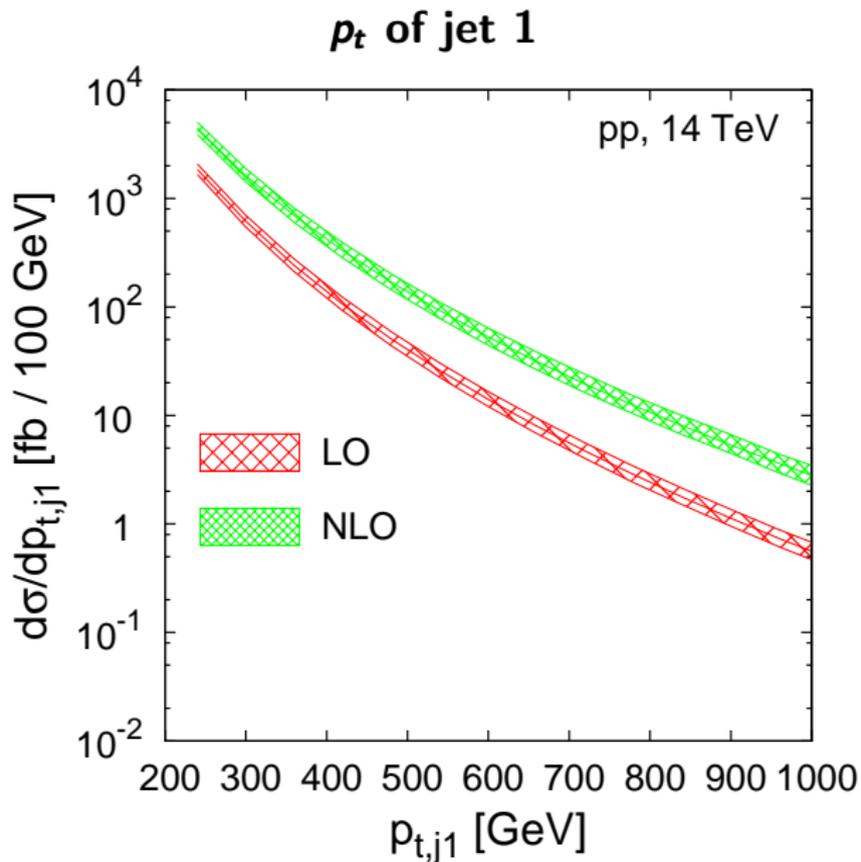
- ▶ Take  $Z + \text{jet}$ ,  $Z + 2 \text{ jet}$  samples, etc.,
- ▶ Attach a parton shower to each of them
- ▶ and combine with procedures such as MLM matching or CKKW matching.

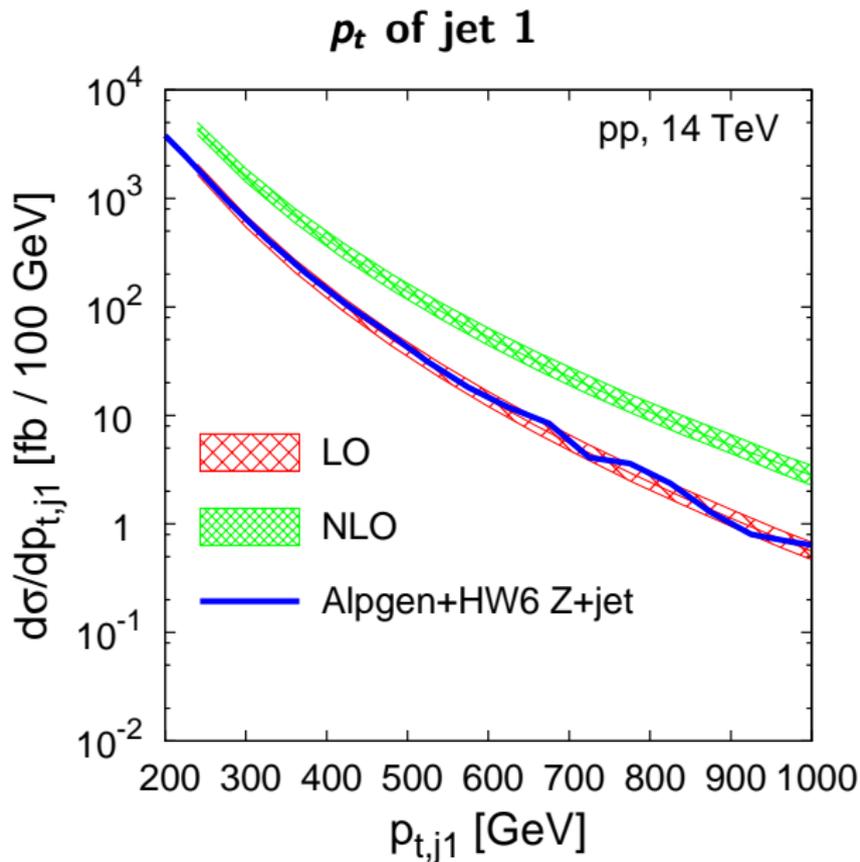


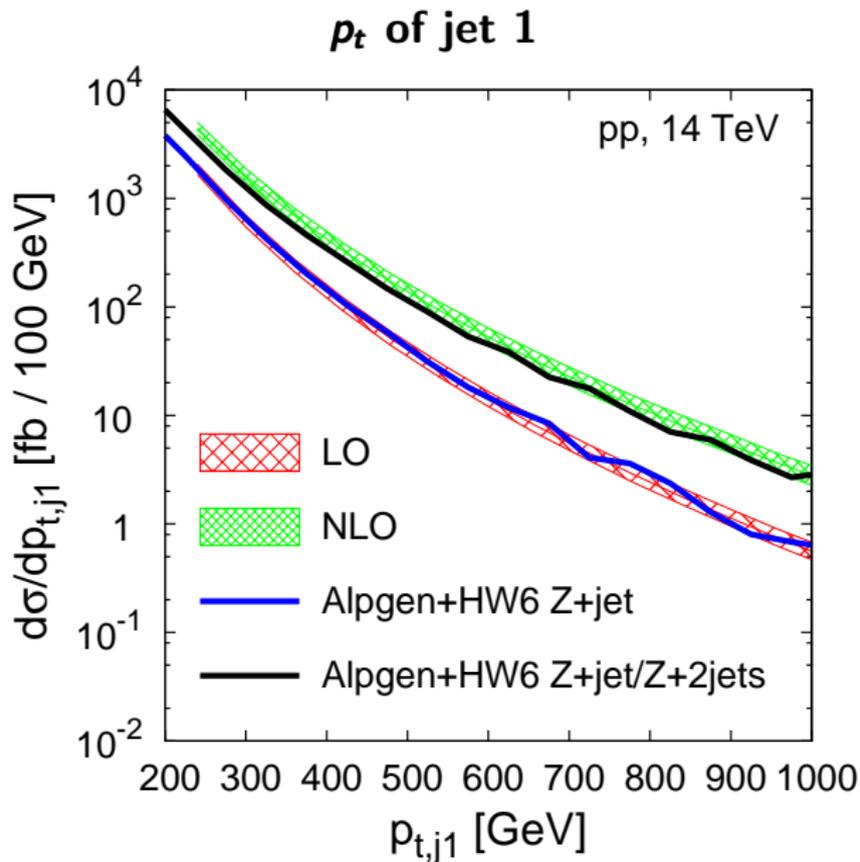




All predictions similar  
and stable

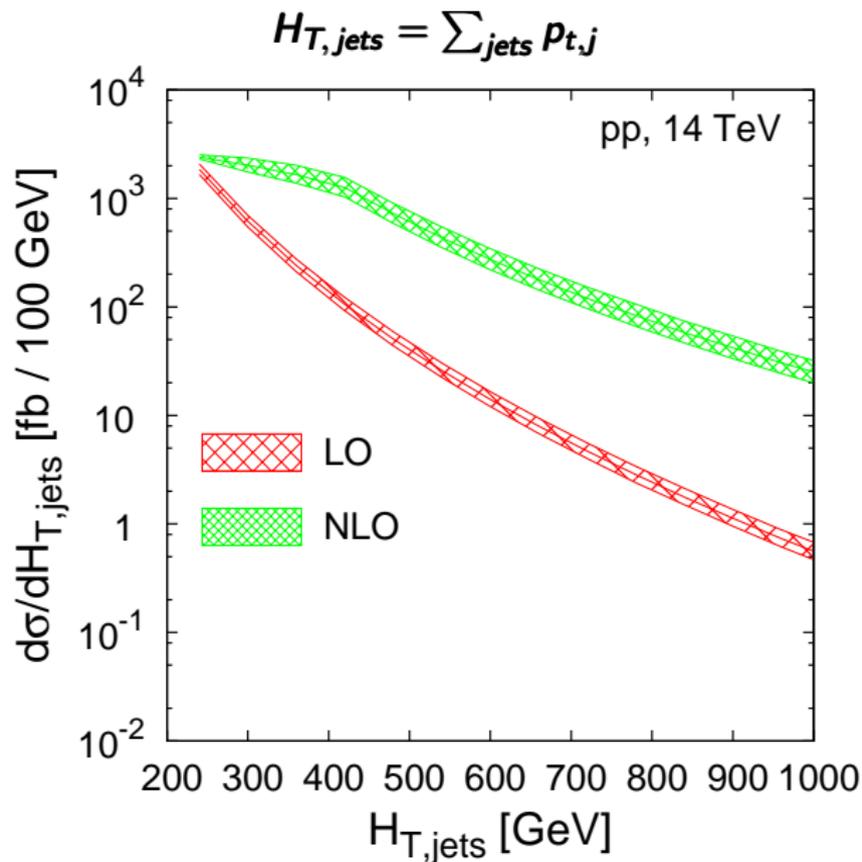


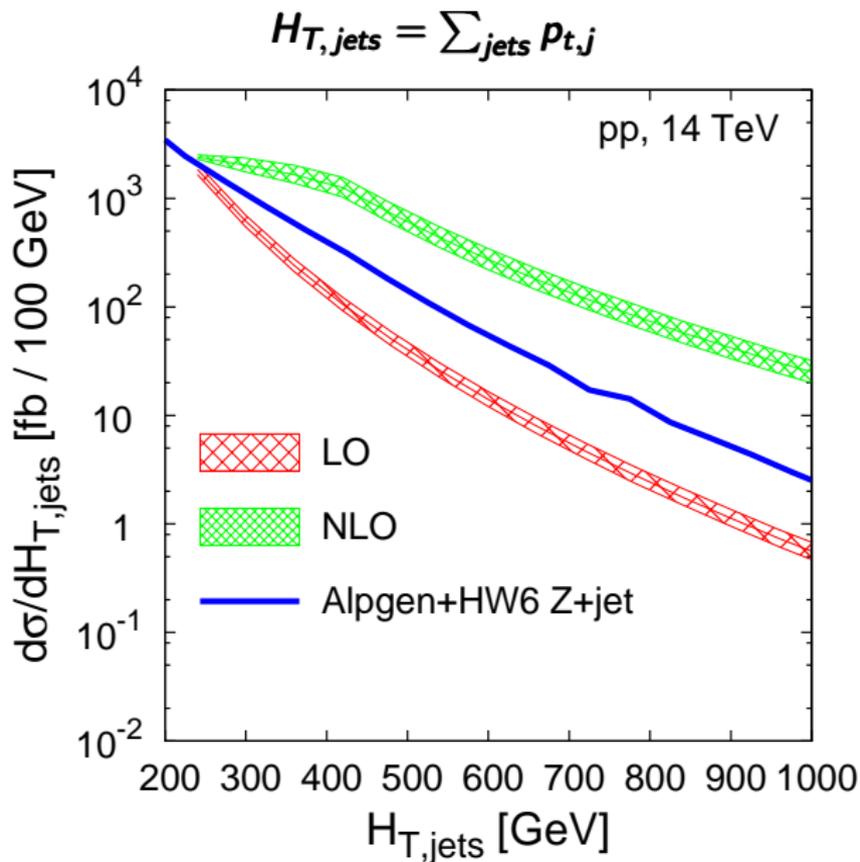
Showered Z+j  $\simeq$  LO



Showered Z+j  $\simeq$  LO

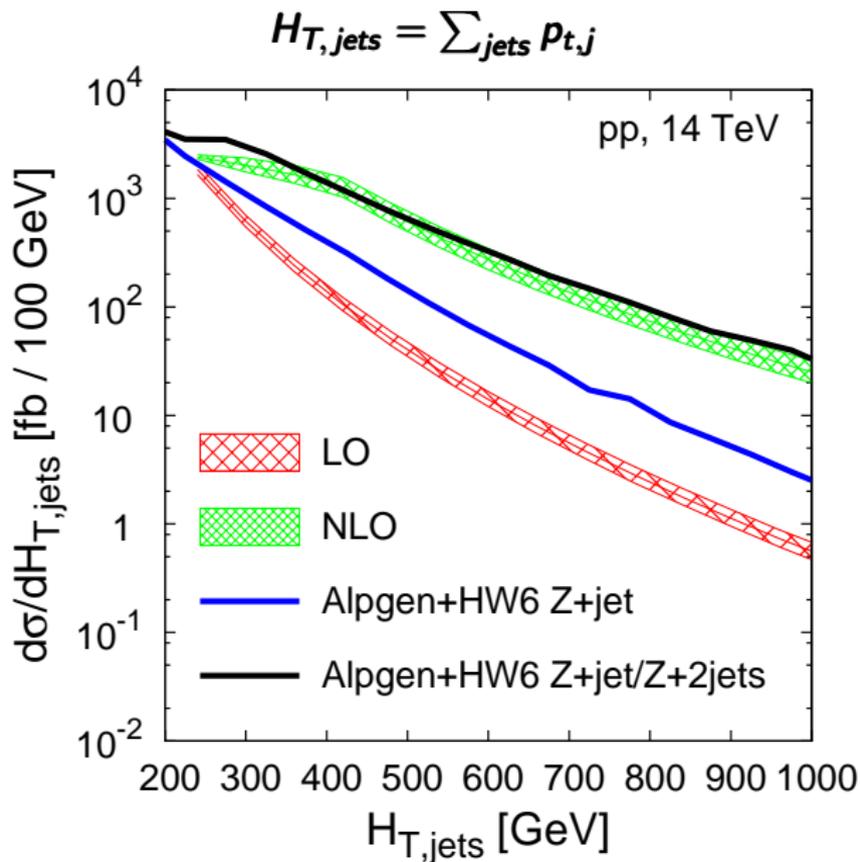
Showered Z+j/Z+2j  
 $\simeq$  NLO





Showered Z+j  $\neq$  LO

cf. also de Aquino,  
Hagiwara, Li & Maltoni '11



Showered Z+j  $\neq$  LO

Showered Z+j/Z+2j  
 $\simeq$  NLO

cf. also de Aquino,  
 Hagiwara, Li & Maltoni '11

It's great that (the widely-used) MLM/CKKW matching procedure correctly approximates the NLO giant  $K$ -factor.

But what happens at NNLO?

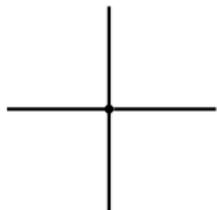
a natural question when LO  $\rightarrow$  NLO convergence is poor

Despite 10 years' calculation, the answer is not yet known

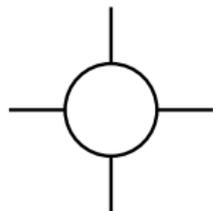
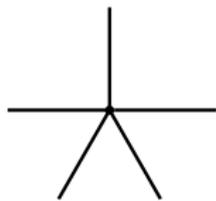
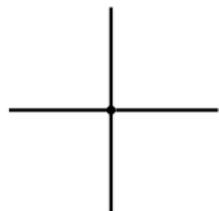
**Our strategy: get an *approximation to NNLO***

specifically, we will approximate a subset of the loop contributions,  
with a method dubbed "LoopSim"

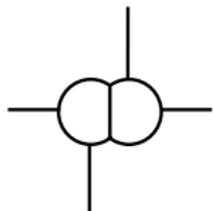
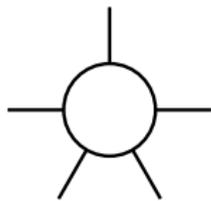
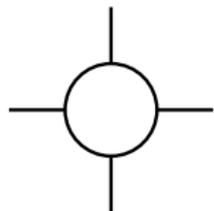
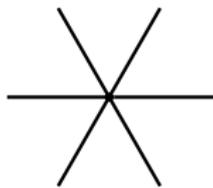
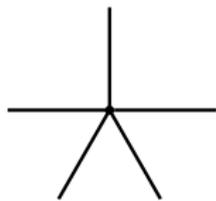
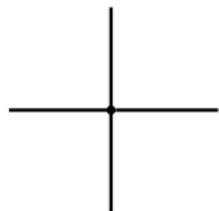
LO

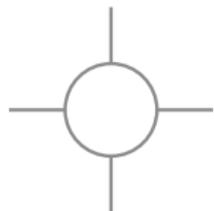
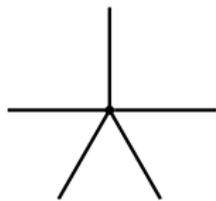
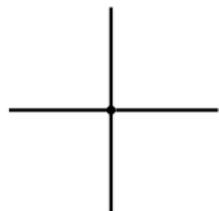


NLO



NNLO



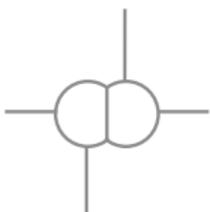
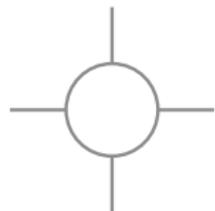
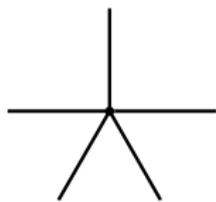
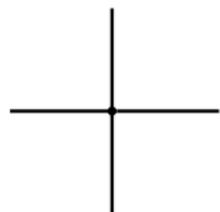
$\bar{n}$ LO

 **Exact**  
 **Approximate**

### Our naming scheme:

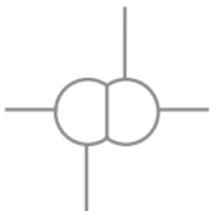
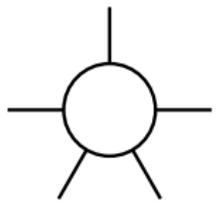
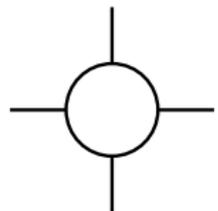
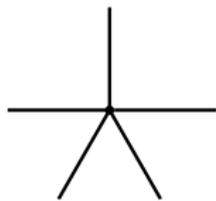
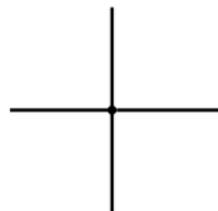
For each loop that we approximate, replace  $N \rightarrow \bar{n}$

- ▶  $\bar{n}$ LO: approx 1-loop diagrams

$\bar{n}\bar{n}\text{LO}$ **Exact****Approximate****Our naming scheme:**

For each loop that we approximate, replace  $N \rightarrow \bar{n}$

- ▶  $\bar{n}\text{LO}$ : approx 1-loop diagrams
- ▶  $\bar{n}\bar{n}\text{LO}$ : approx 1- and 2-loops

$\bar{n}$ NLO**Exact****Approximate****Our naming scheme:**

For each loop that we approximate, replace  $N \rightarrow \bar{n}$

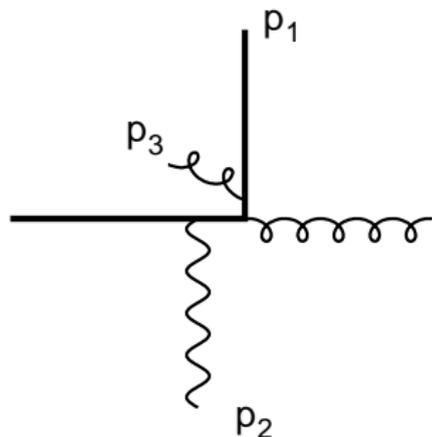
- ▶  $\bar{n}$ LO: approx 1-loop diagrams
- ▶  $\bar{n}\bar{n}$ LO: approx 1- and 2-loops
- ▶  $\bar{n}$ NLO: approx 2-loop only

## First try $Z + \text{jet}$ @ $\bar{n}\text{LO}$ :

Take the “leading” process  
[ $Z + \text{jet}$  @ LO]

and add in process with one extra jet.  
[i.e. include  $Z + 2 \text{ jets}$  @ LO]

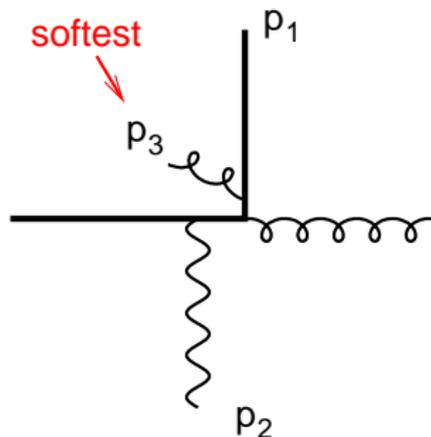
**approximate** the 1-loop  $Z + \text{jet}$  term, by requiring  
cancellation of all divergences  
[those from singly unresolved limit of  $Z + 2 \text{ jets}$ ]



**Z + 2 partons**

$$|M^2(p_1, p_2, p_3)|$$

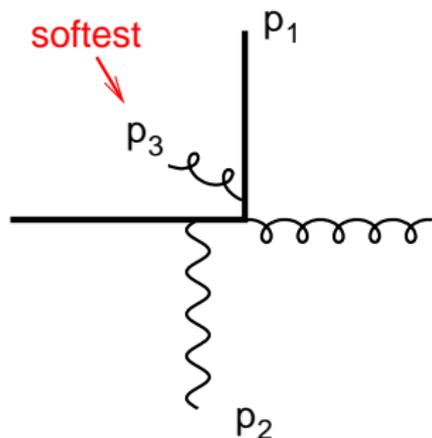
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it  $\equiv$  remove it from event, reshuffle other momenta;  
weight of looped event is  $(-1) \times$  weight of tree-level event



**Z + 2 partons**

$$|M^2(p_1, p_2, p_3)|$$

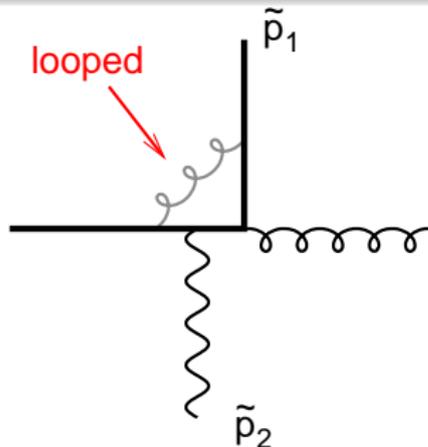
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it  $\equiv$  remove it from event, reshuffle other momenta;  
 weight of looped event is  $(-1) \times$  weight of tree-level event



**Z + 2 partons**

$$|M^2(p_1, p_2, p_3)|$$

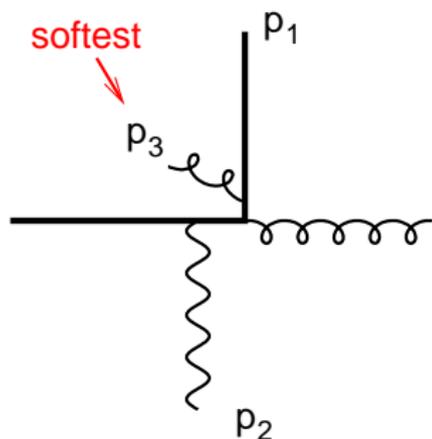
+



**Z + 1 parton + 1 sim. loop**

$$-|M^2(p_1, p_2, p_3)|$$

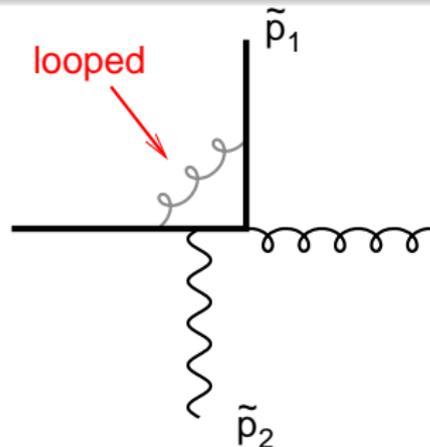
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it  $\equiv$  remove it from event, reshuffle other momenta;  
 weight of looped event is  $(-1) \times$  weight of tree-level event



**Z + 2 partons**

$$|M^2(p_1, p_2, p_3)|$$

+



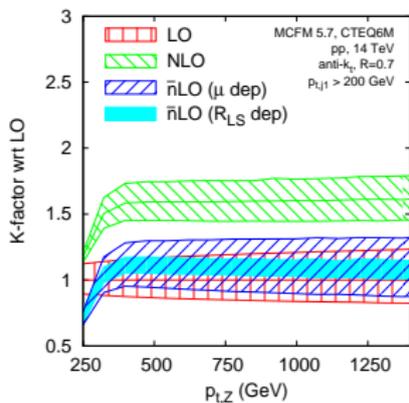
**Z + 1 parton + 1 sim. loop**

$$-|M^2(p_1, p_2, p_3)|$$

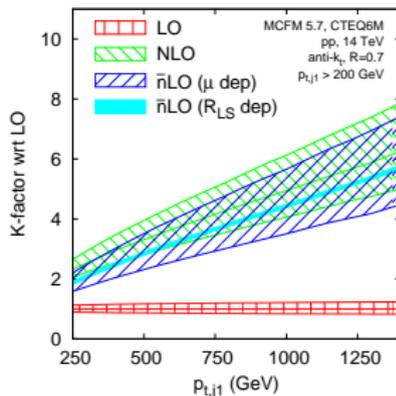
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it  $\equiv$  remove it from event, reshuffle other momenta; weight of looped event is  $(-1) \times$  weight of tree-level event

This cancels all the “single-unresolved” divergences in the Z+2 events

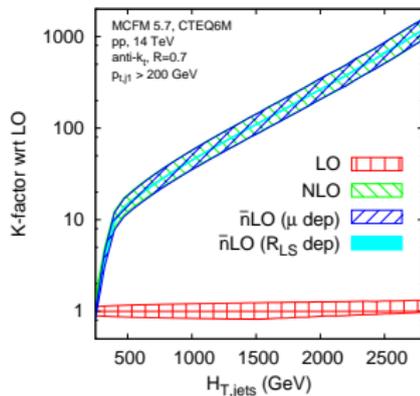
### $p_t$ of Z-boson



### $p_t$ of jet 1



### $H_{T,jets} = \sum_{jets} p_{t,j}$



When the  $K$ -factors are large,  $\bar{n}$ LO agrees well with NLO

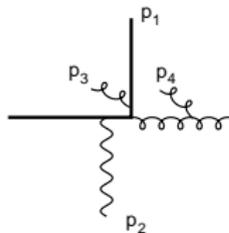
Works with similar “quality” to MLM matching

**MLM/CKKW matching also effectively provide  $\bar{n}$ LO type accuracy**

**How does LoopSim compare to MLM/CKKW?**

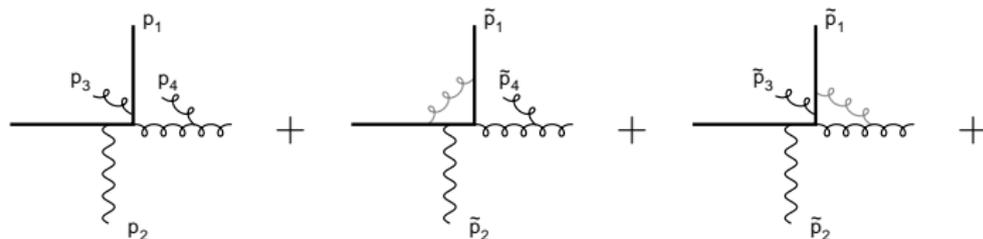
1. Does not rely on shower (✓: simplicity; ✗: not easily integrated with shower MCs)
2. Does not need arbitrary separation of  $Z+1/Z+2$ /etc. samples with (hard-to-choose) momentum cutoff
3. Can easily be extended beyond LO matching

add tree-level Z+3,  
cancel divergences in single + doubly unresolved limits:  $\bar{n}\bar{n}$ LO



$$|M^2(p_1, p_2, p_3, p_4)|$$

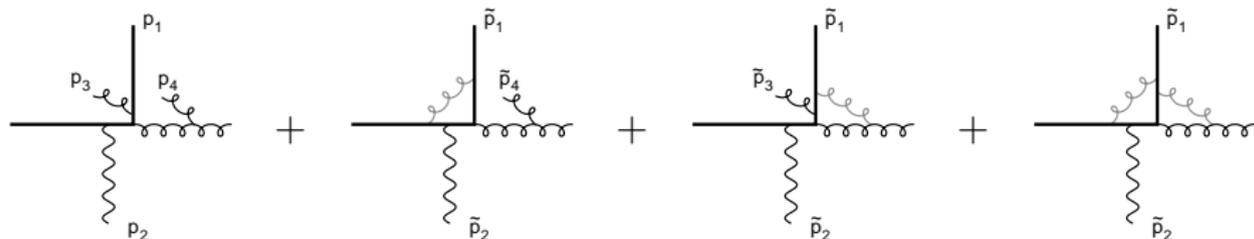
add tree-level Z+3,  
cancel divergences in single + doubly unresolved limits: **n̄nLO**



$$|M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)|$$

Separately loop either of the 2 softest emissions: provides approx of 1-loop

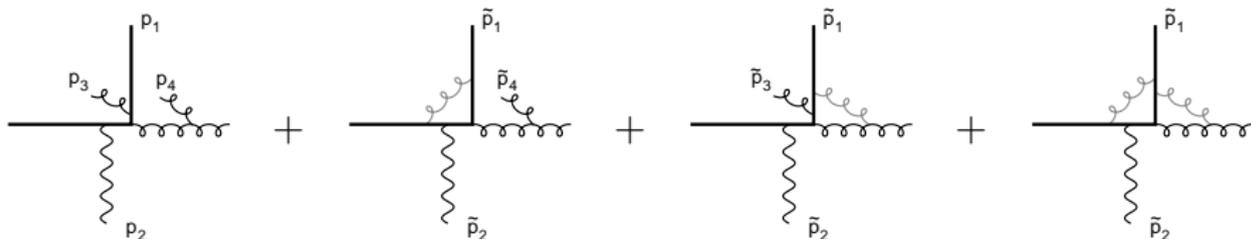
add tree-level Z+3,  
cancel divergences in single + doubly unresolved limits:  $\bar{n}\bar{n}$ LO



$$|M^2(p_1, p_2, p_3, p_4)| \quad -|M^2(p_1, p_2, p_3, p_4)| \quad -|M^2(p_1, p_2, p_3, p_4)| \quad +|M^2(p_1, p_2, p_3, p_4)|$$

Simultaneously loop each of the 2 softest emissions: provides approx of 2-loop  
Total of tree plus approx 1- and 2-loop pieces gives zero

add tree-level Z+3,  
cancel divergences in single + doubly unresolved limits: **n̄nLO**



$$|M^2(p_1, p_2, p_3, p_4)| \quad -|M^2(p_1, p_2, p_3, p_4)| \quad -|M^2(p_1, p_2, p_3, p_4)| \quad +|M^2(p_1, p_2, p_3, p_4)|$$

Simultaneously loop each of the 2 softest emissions: provides approx of 2-loop  
Total of tree plus approx 1- and 2-loop pieces gives zero

add in (exact Z+2 @ 1-loop) – (approximate Z+2 @ 1-loop)  
+ extra simulated 2-loop piece to cancel new Z+2@1-loop divergences

**This is n̄NLO**

The 2-loop piece has the topology of the LO diagram.

The “mistake” we make in approximating it should therefore be a “pure”  $\mathcal{O}(\alpha_s^2)$  correction, without any large enhancements from new NLO type topologies.

$$\begin{aligned}\sigma_{\bar{n}\text{NLO}} &= \sigma_{\text{NNLO}} + \mathcal{O}(\alpha_s^2 \sigma_{\text{LO}}) \\ &= \sigma_{\text{NNLO}} \left( 1 + \mathcal{O}\left(\frac{\alpha_s^2}{K_{\text{NNLO}}}\right) \right)\end{aligned}$$

$$K_{\text{NNLO}} = \frac{\sigma_{\text{NNLO}}}{\sigma_{\text{LO}}} \sim K_{\text{NLO}} \gg 1$$

The *relative* contribution of the neglected piece is suppressed by the large  $K$ -factor.

n̄NLO should be a good approximation to NNLO when the  $K$ -factor is large and due to new higher-order topologies.

# Testing $\bar{n}$ NLO, in 3 processes

[making use of existing NLO calculations]

1.  $Z@NLO$  and  $Z+j@NLO \rightarrow Z@\bar{n}NLO$

with MCFM; compare to exact NNLO from DYNNLO

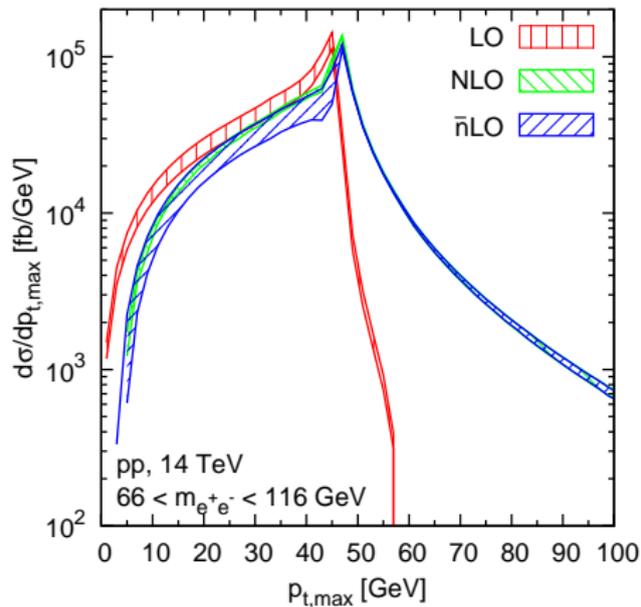
2.  $Z+j@NLO$  and  $Z+2j@NLO \rightarrow Z+j@\bar{n}NLO$

with MCFM

3.  $2j@NLO$  and  $3j@NLO \rightarrow 2j@\bar{n}NLO$

with NLOjet++

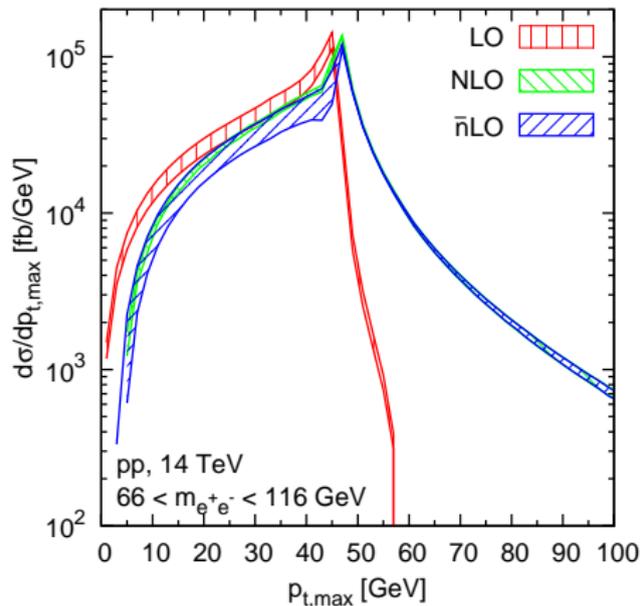
## $\bar{n}$ LO v. NLO



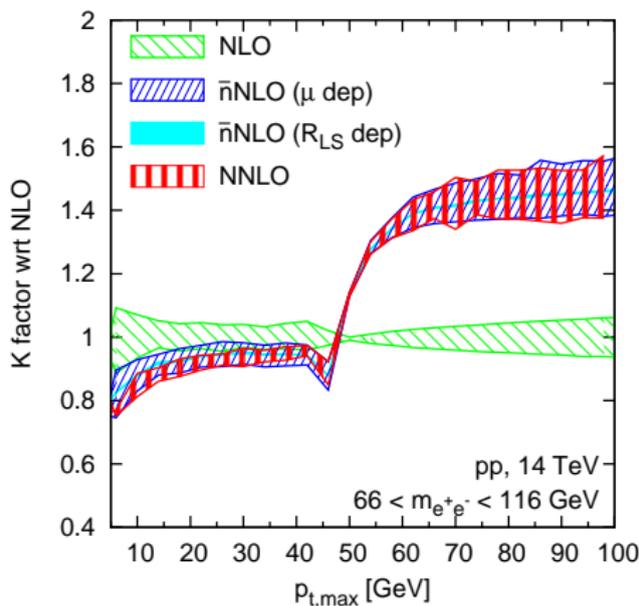
Z (i.e. DY) with Z+j from MCFM & LoopSim

For  $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$  (giant  $K$ -factor!) it had to work  
For  $p_{t,\ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$  it's remarkable that it still works

## $\bar{n}$ LO v. NLO



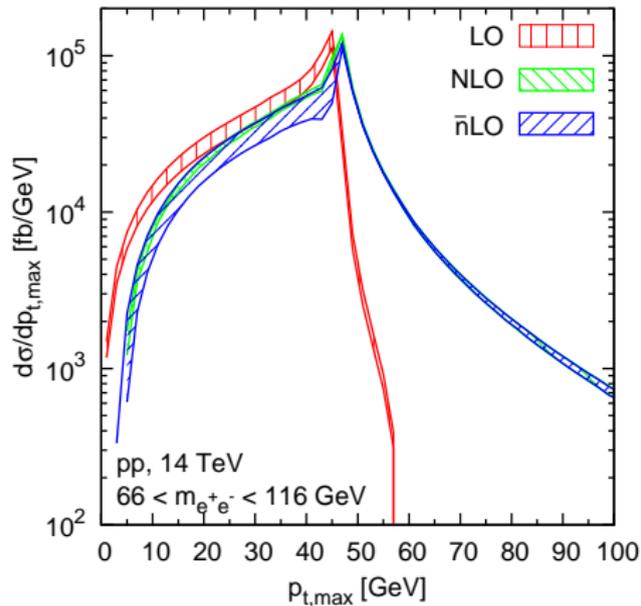
## $\bar{n}$ NLO v. NNLO



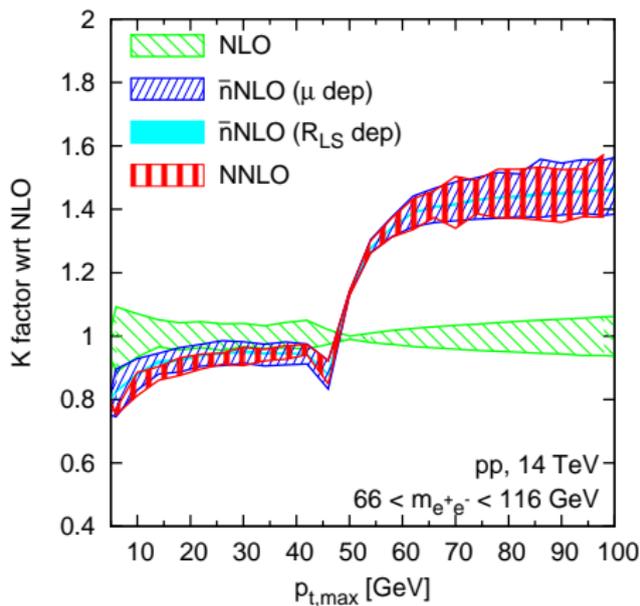
NNLO from DYNNOLO, Z (i.e. DY) with Z+j from MCFM & LoopSim

For  $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$  (giant  $K$ -factor!) it had to work  
 For  $p_{t,\ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$  it's remarkable that it still works

## $\bar{n}$ LO v. NLO



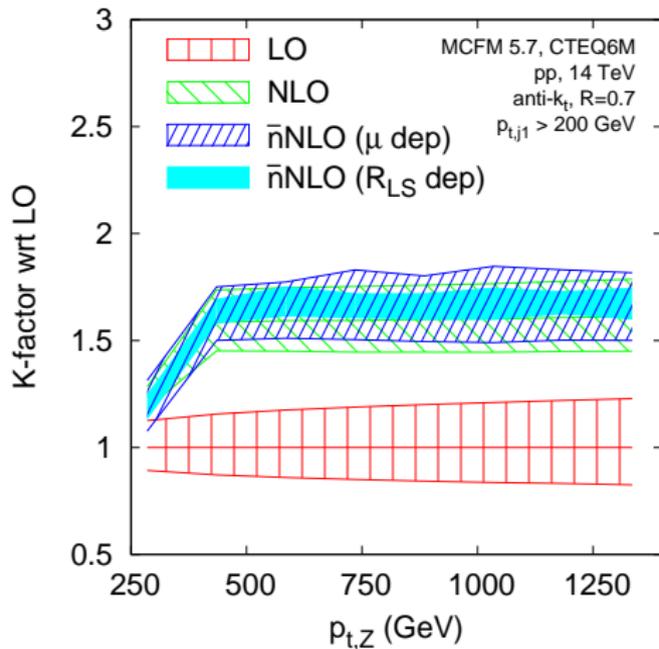
## $\bar{n}$ NLO v. NNLO



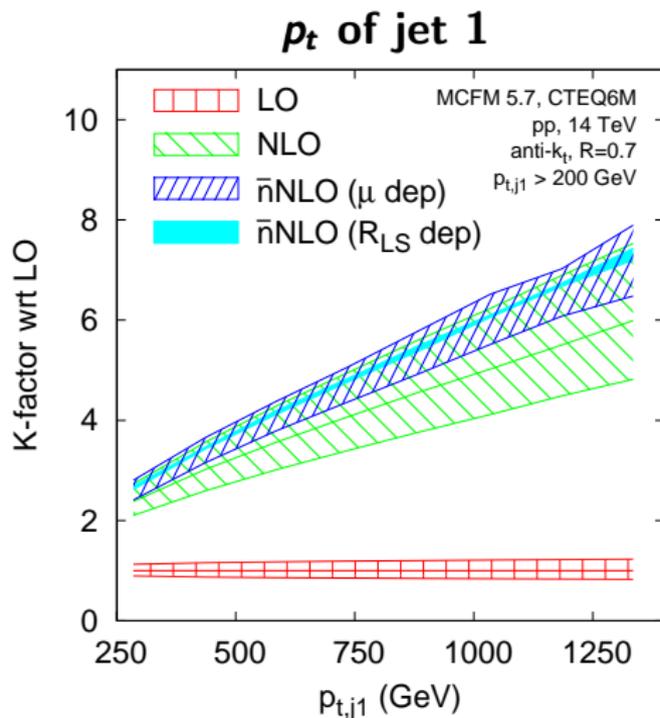
NNLO from DYNNOLO, Z (i.e. DY) with Z+j from MCFM & LoopSim

For  $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$  (giant  $K$ -factor!) it had to work  
 For  $p_{t,\ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$  it's remarkable that it still works

$p_t$  of Z-boson

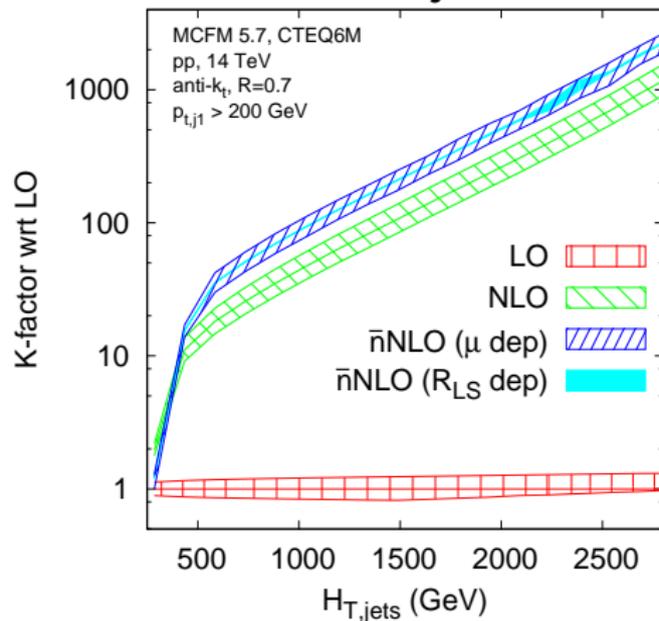


- ▶  $p_{tZ}$  distribution didn't have giant  $K$ -factors.
- ▶  $\bar{n}$ NLO brings no benefit  
 To get improvement you would need exact 2-loop terms



- ▶  $p_{tj}$  distribution seems to converge at  $\bar{n}$ NLO
- ▶ scale uncertainties reduced by  $\sim$  factor 2

$$H_{T,jets} = \sum_{jets} p_{t,j}$$



- ▶ Significant further enhancement for  $H_{T,jets}$
- ▶  $\bar{n}$ NLO brings clear message:

**$H_{T,jets}$  is not a good observable!**

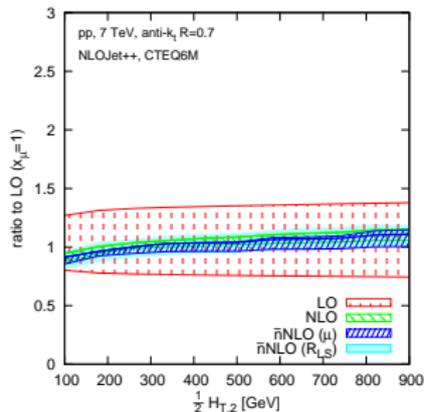
## $H_T$ (effective mass) type observables are widely used in searches

- ▶  $H_T$  has a steeply falling distribution (like  $p_{tj}$ ,  $p_{tZ}$ )
- ▶ At each order (NLO, NNLO), an extra (soft) jet contributes to the  $H_T$  sum e.g. from ISR
- ▶ That shifts  $H_T$  up, which translates to a substantial increase in the cross section

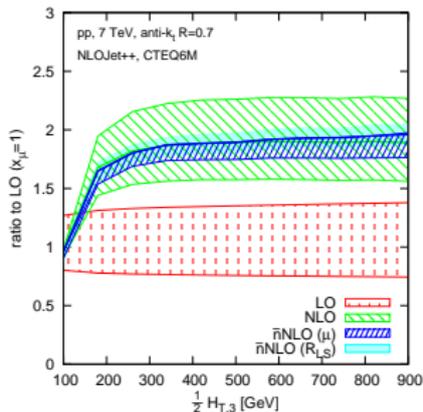
We can test this hypothesis for plain jet events, using a truncated sum,

$$H_{T,n} = \sum_{i=1}^n p_{t,\text{jet } i}$$

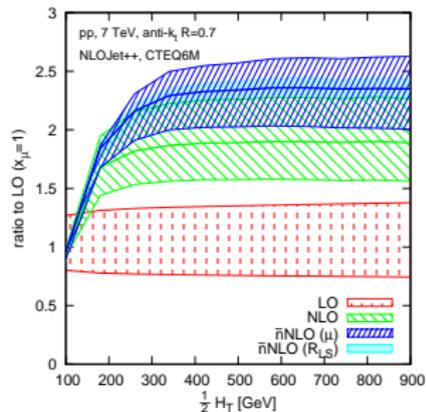
## $H_{T,2}$



## $H_{T,3}$



## $H_{T,\infty}$



**A clear message:**

for a process with  $n$  objects at lowest order, use  $H_{T,n}$

Do you know what gets used in your experiment's searches?

Be aware that giant  $K$ -factors exist

Always look one order beyond the leading order, for example with  
MLM/CKKW matching

New tool to get good predictions in such cases: **LoopSim**

Basically an “operator” to generate approximations to unknown loops

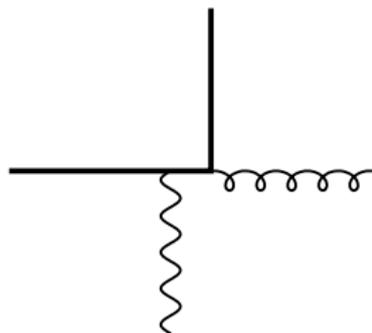
Combine  $Z+j@NLO$ ,  $Z+2j@NLO$  to get “ $\bar{n}NLO$ ”  $Z+jet$

It sometimes works even beyond “giant- $K$ -factor” regions

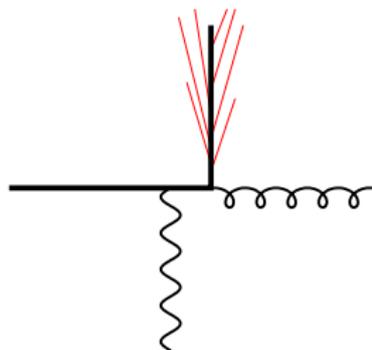
Watch out for  $H_T$

Even for simple processes, it converges very poorly  
unless you define it carefully (limit number of objects in sum)

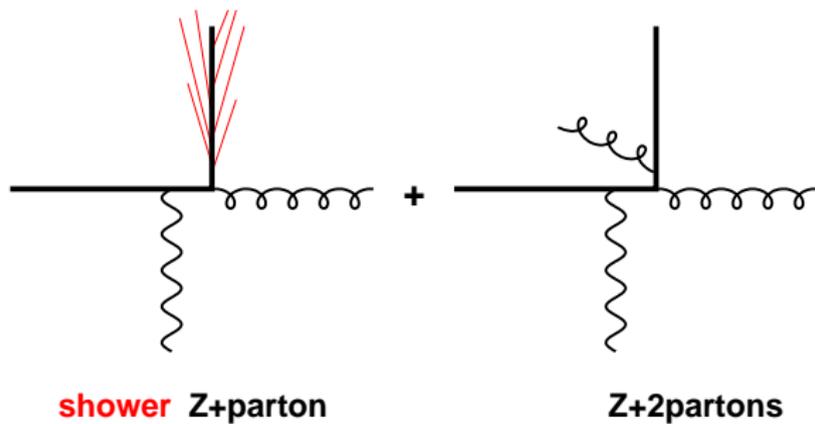
# EXTRAS

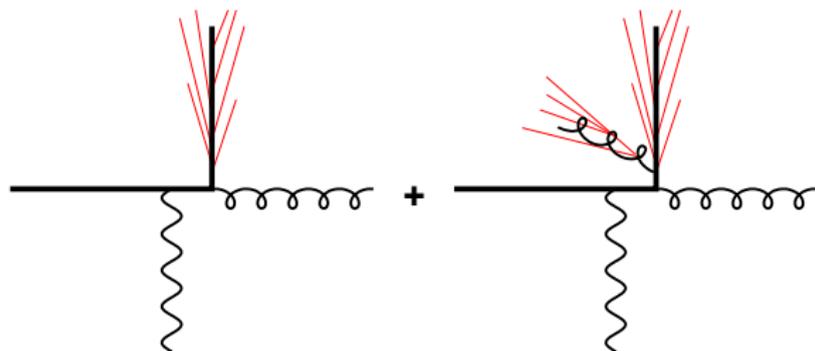


**Z+parton**



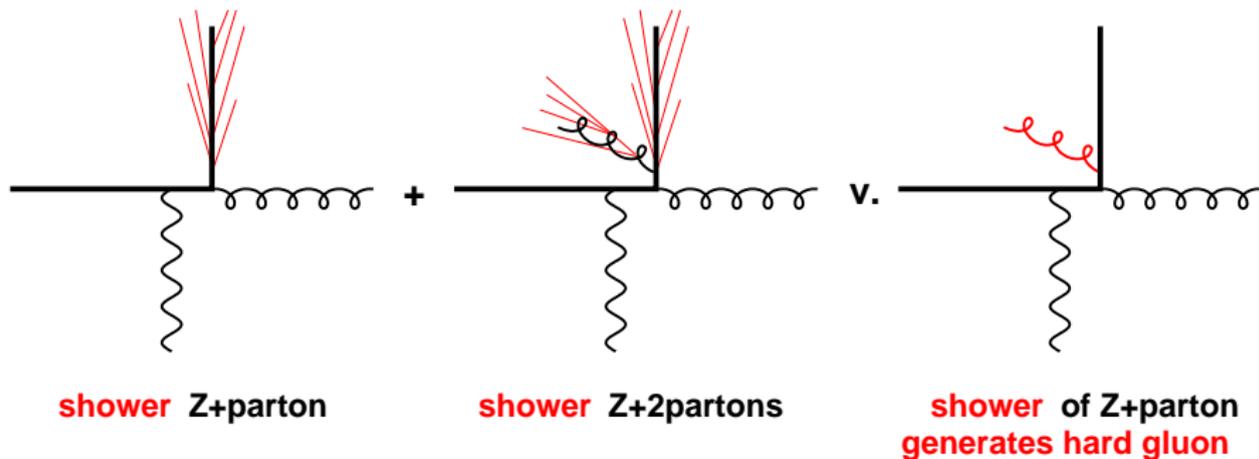
**shower** Z+parton

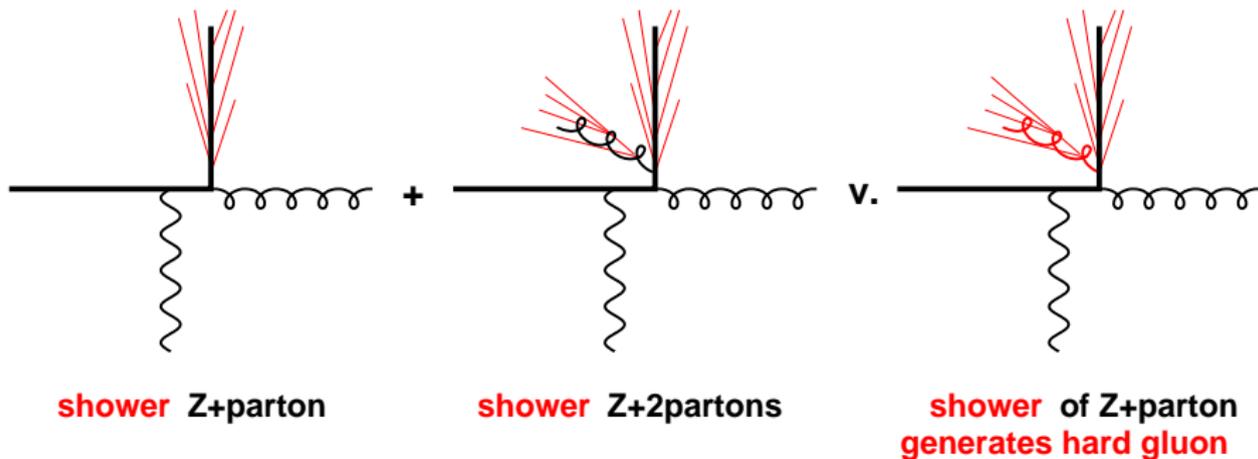


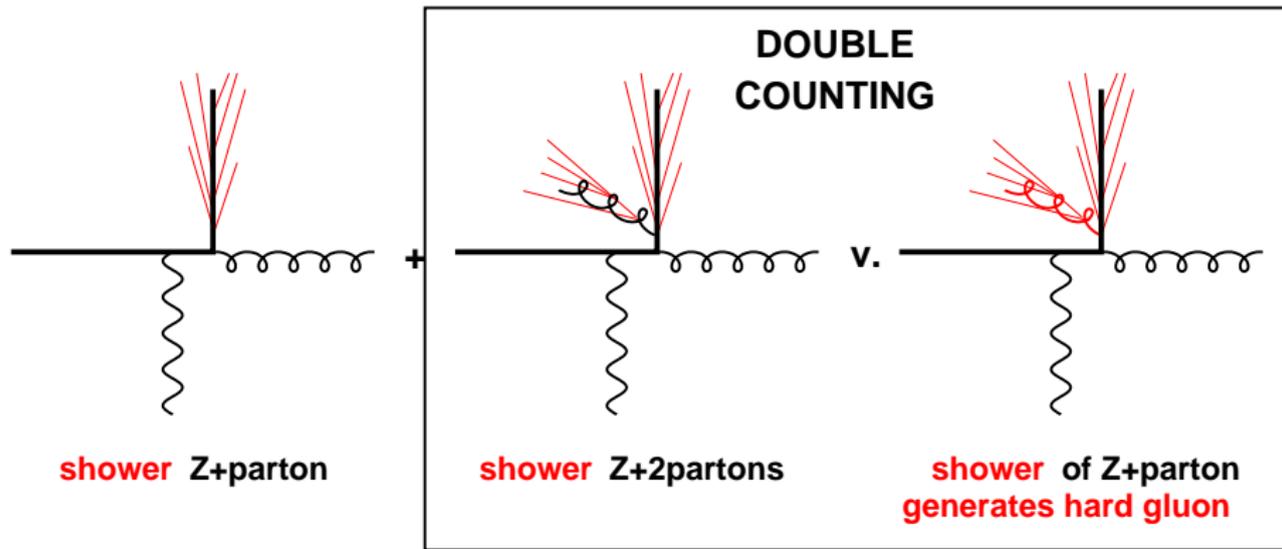


**shower** Z+parton

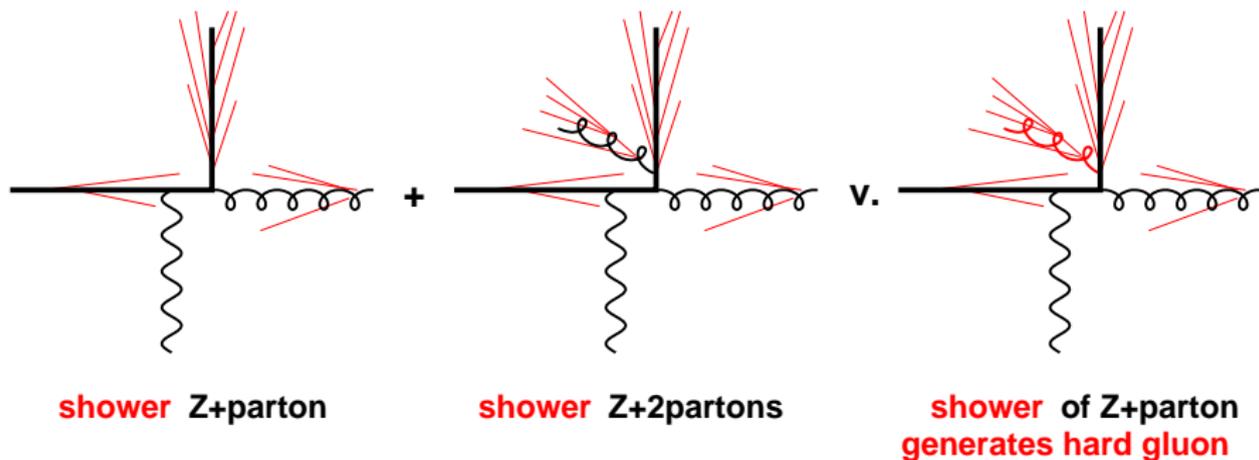
**shower** Z+2partons







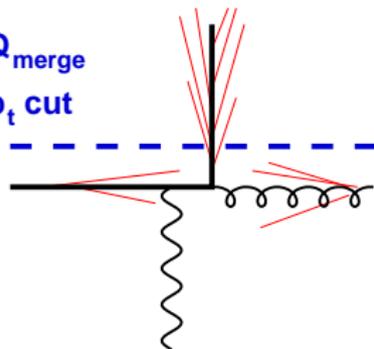
Z + parton implicitly includes part of Z + 2 partons  
It's just that the 2nd parton isn't always explicitly "visible"



- ▶ MLM merging relies on parton shower to help figure out what fraction of  $Z + \text{parton}$  is really  $Z + 2 \text{ partons}$ .
- ▶ In a few slides, we will try to do that without the parton shower

ACCEPT

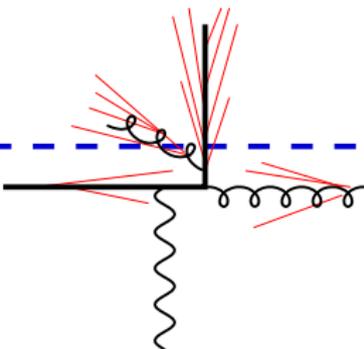
$Q_{\text{merge}}$   
 $p_t \text{ cut}$



shower  $Z+\text{parton}$

ACCEPT

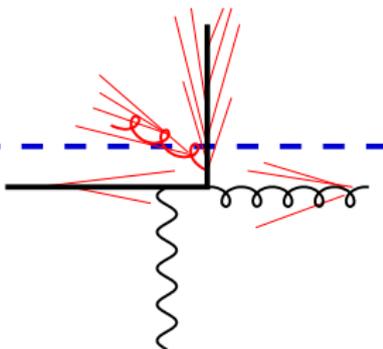
+



shower  $Z+2\text{partons}$

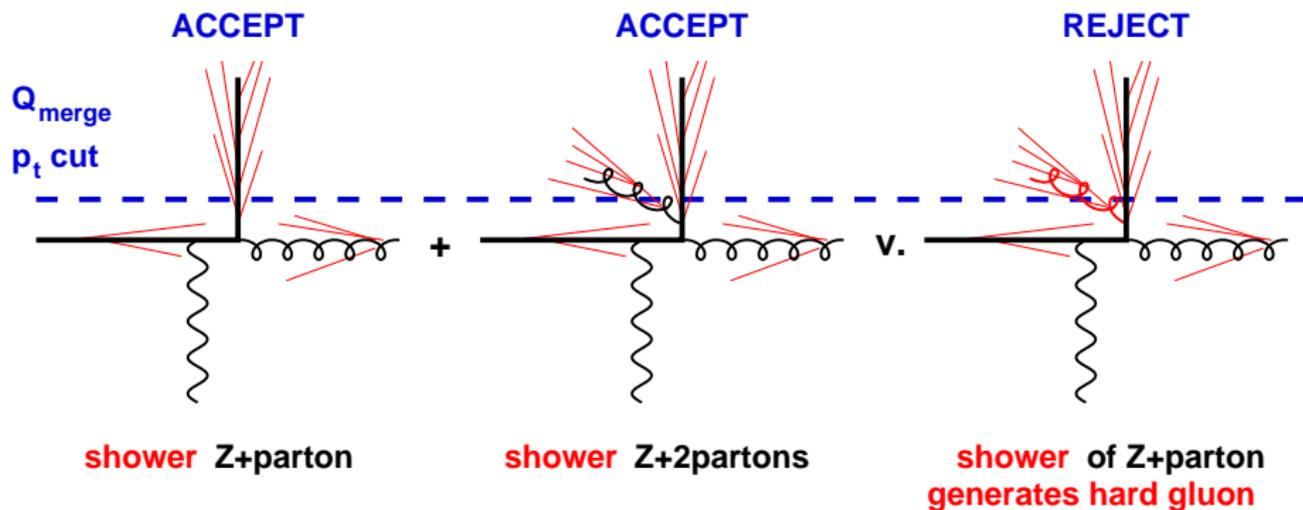
v.

REJECT



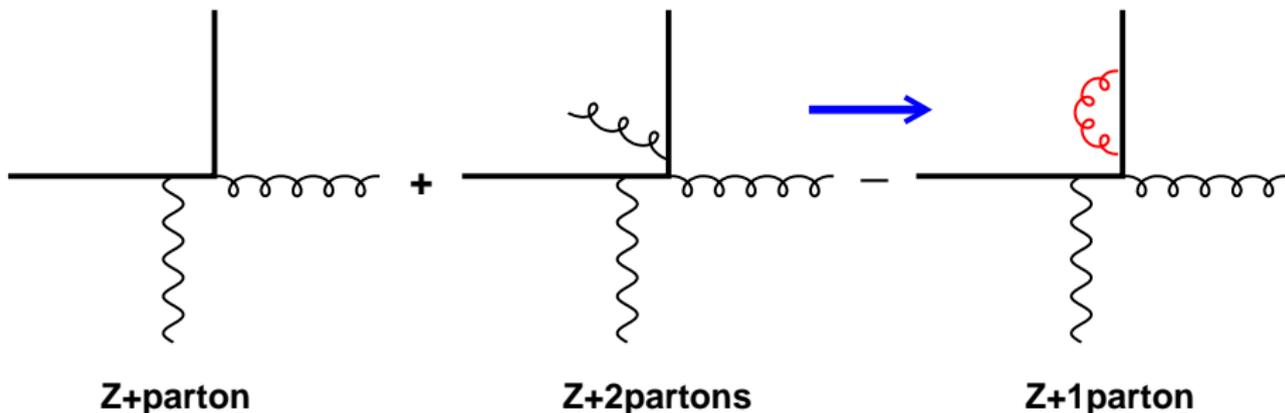
shower of  $Z+\text{parton}$   
generates hard gluon

- ▶ MLM merging relies on parton shower to help figure out what fraction of  $Z + \text{parton}$  is really  $Z + 2 \text{ partons}$ .
- ▶ In a few slides, we will try to do that without the parton shower



- ▶ MLM merging relies on parton shower to help figure out what fraction of  $Z + \text{parton}$  is really  $Z + 2 \text{ partons}$ .
- ▶ In a few slides, we will try to do that without the parton shower

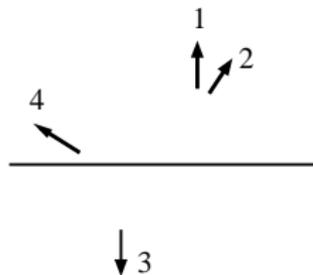
**SUBTRACT**



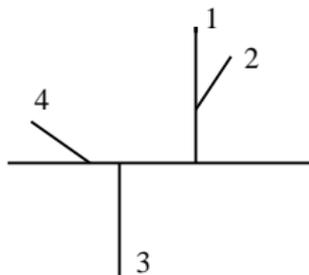
**softest particle of Z+2 is "looped"  
= removed from event (kinematics reshuffled)**

- ▶ For every  $Z + 2$  parton ( $2 \rightarrow 3$ ) event, figure out what what  $2 \rightarrow 2$  event it would really have come from  
"Loop" the softest parton  
[Don't actually explicitly calculate any loop diagrams: simulate the loops]
- ▶ Subtract that  $2 \rightarrow 2$  event  
Unlike MLM, no cutoffs on  $2 \rightarrow 3$  events  
If done properly, divergences will cancel

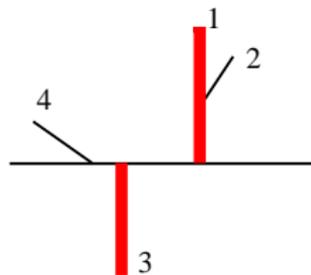
(a) Input event



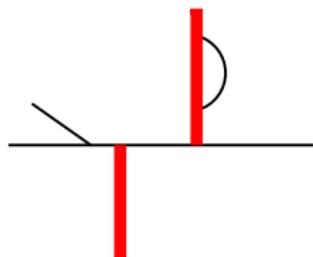
(b) Attributed emission seq.



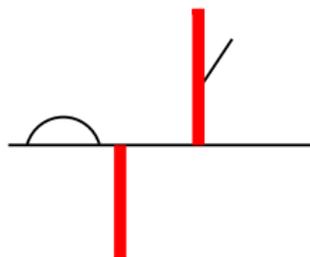
(c) Born particle ID



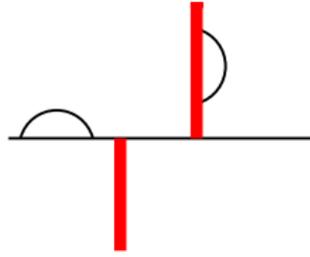
(d) Output 1-loop event



(e) 2nd output 1-loop event



(f) Output 2-loop event



- ▶ Use jet algorithm to assign a branching structure to event à la CKKW
- ▶ The particles that are softest are the ones that will be “looped”

Define operators:

$U_\ell(\text{event } E) \equiv$  all simulated  $\ell$ -loop events from  $E$

$$U_\forall(\text{event}) \equiv \sum_{\ell=0} U_\ell(\text{event})$$

“U” stands for unitarisation (cancellation of all divergences)  
sum of all diagrams (essentially) adds up to zero

To combine  $Z+j$  with  $Z+2j$  take

$$Z+j@{\bar{n}}\text{LO} \equiv Z+j@\text{LO} + U_\forall(Z+2j@\text{LO})$$

we use “ $\bar{n}$ LO” to emphasize that this is a crude approximation  
to an actual NLO calculation — the exact loops are missing  
NB:  $U_\forall$  here includes  $\ell = 0, 1$

Just replace simulated loops with exact loops  
Apply LoopSim to exact 1-loop to get (e.g.) simulated 2-loop terms

---

$E_{n,\ell} \equiv$  event with  $n$  partons and  $\ell$  exact loops  
 $U_{\forall,\ell} \equiv$  operator to apply when  $\ell$  exact loops known

$$U_{\forall,1}(E_{n,0}) = U_{\forall}(E_{n,0}) - U_{\forall}(U_1(E_{n,0}))$$

$$U_{\forall,1}(E_{n,1}) = U_{\forall}(E_{n,1})$$

---

$$Z+j@ \bar{n}\text{NLO} = Z+j@\text{NLO} + U_{\forall,1}(Z+2j@\text{NLO}_{\text{only}})$$

Extension to NLO, NNLO, multi-leg, etc. is almost trivial in LoopSim

Not the case in methods that merge with parton showers too