

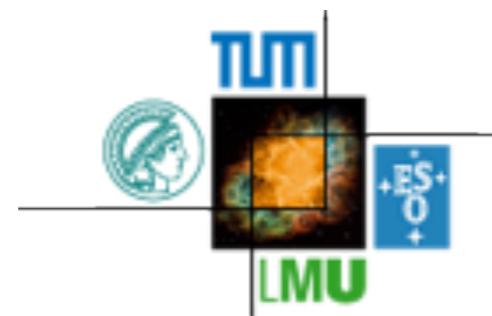
Standard-Model Prediction of ϵ_K at NNLO

Fermilab Theory Seminar, November 3, 2011

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In collaboration with Martin Gorbahn

Phys. Rev. D 82 (2010) 094026 [arXiv:1007.0684];
arXiv:1108.2036



Tension in the Data?

Experiment:

ε_K is measured precisely:

$|\varepsilon_K| = 2.228(11) \times 10^{-3}$.

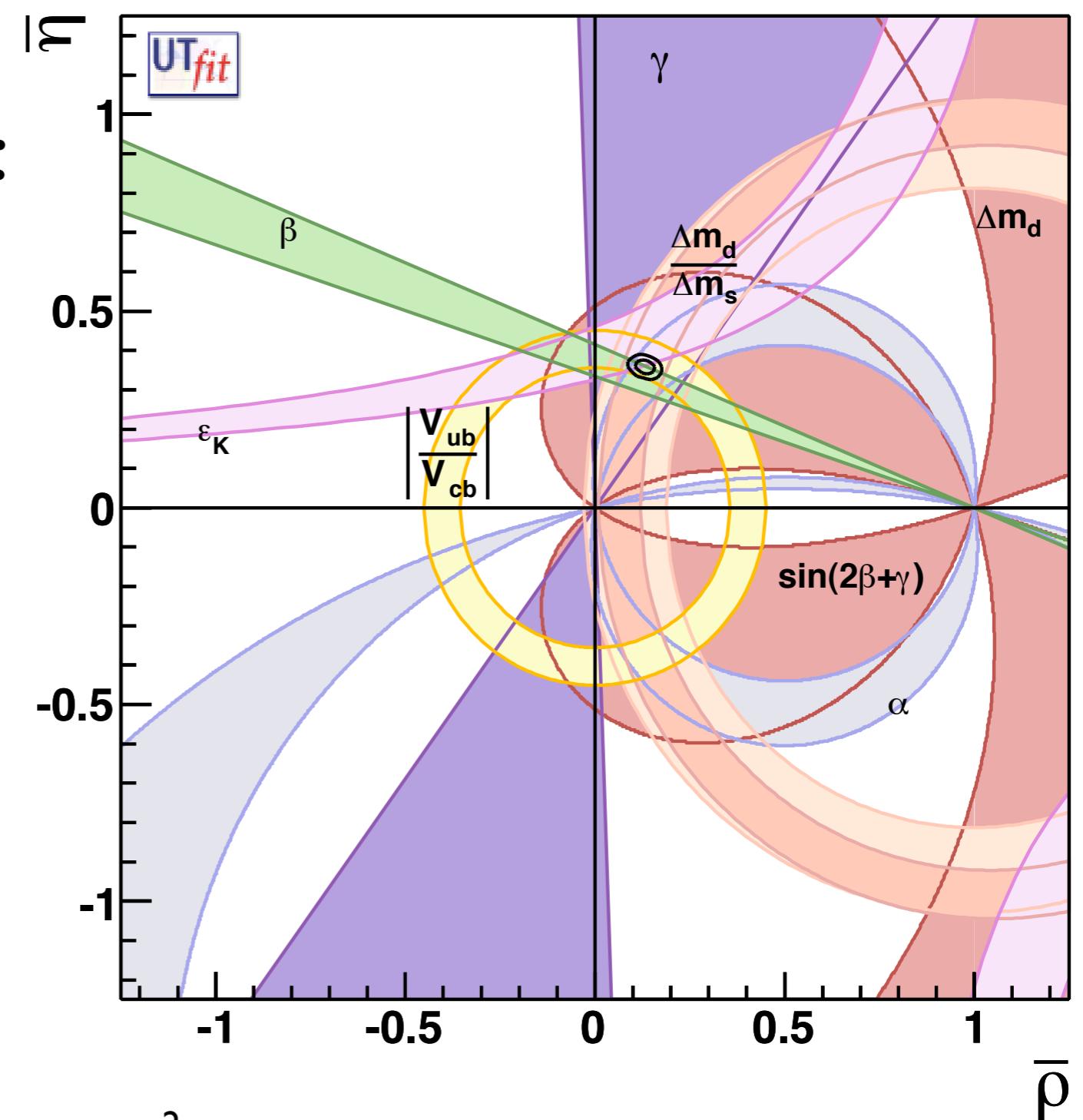
[PDG2010]

Theory:

ε_K can be calculated
reliably:

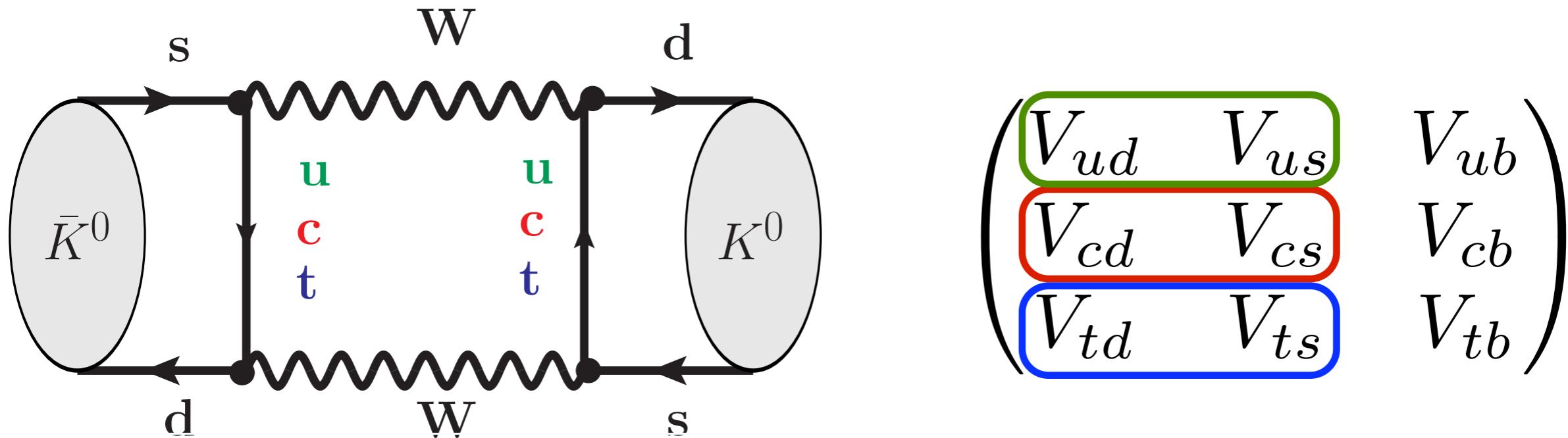
$|\varepsilon_K| = 1.83(27) \times 10^{-3}$.

[NLO SM prediction]



ϵ_K in the Standard Model

- Neutral Kaon mixing is a FCNC transition
- Loop-induced in the Standard Model (SM)

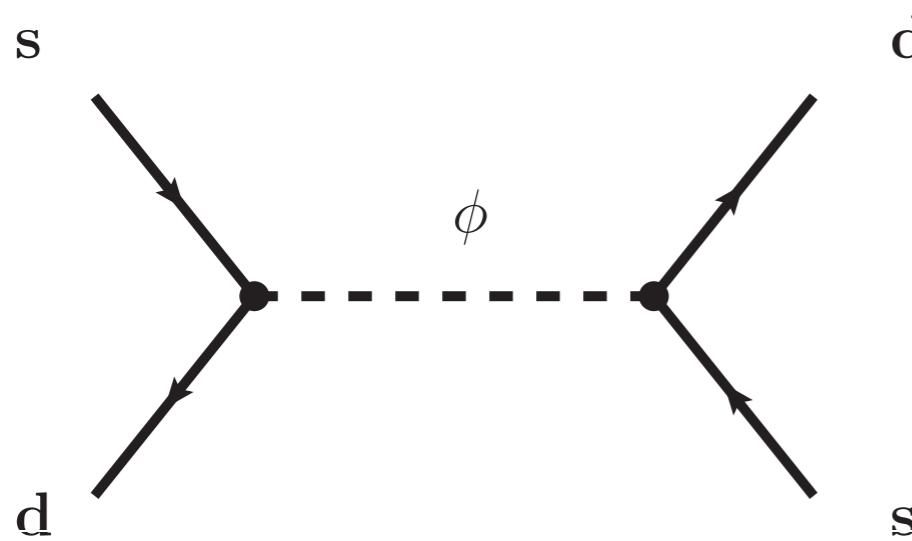


SM \rightarrow CP violation by Complex Phase in the CKM matrix

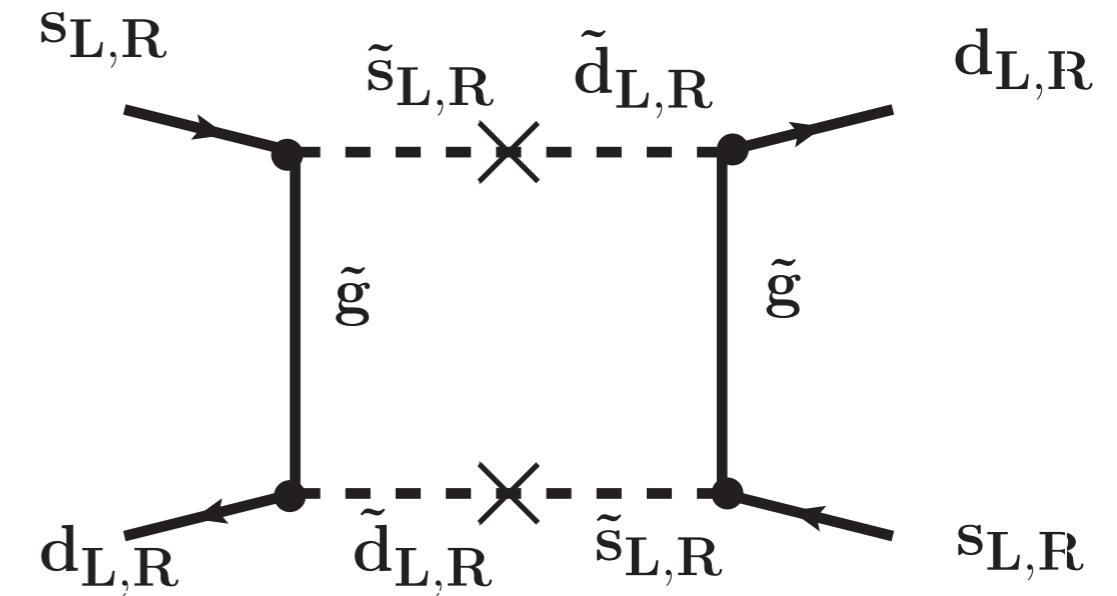
Top-quark contribution CKM suppressed
⇒ sensitivity to deviations from MFV

ϵ_K beyond the Standard Model

Many more sources of CP violation:



2HDM Type III



SUSY models

+ Technicolour, extra dimensions,

Chiral Enhancement of non-SM operators

Strong constraints from ϵ_K !

Outline

- Introduction: Effective Field Theory
- Calculation of ϵ_K
 - Long-distance contributions
 - Short-distance contributions
- Result and error budget
- Summary and outlook

ϵ_K - Definition

ϵ_K describes indirect CP violation in the neutral Kaon system

$$\epsilon_K = \frac{\langle (\pi\pi)_{I=0} | K_L \rangle}{\langle (\pi\pi)_{I=0} | K_S \rangle} = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12})}{\Delta m_K} + \xi \right)$$

$\langle K^0 | H^{| \Delta S | = 2} | \bar{K}^0 \rangle$

$\frac{\text{Im}\langle (\pi\pi)_{I=0} | K^0 \rangle}{\text{Re}\langle (\pi\pi)_{I=0} | K^0 \rangle}$

$$\phi_\epsilon = \arctan \frac{\Delta m_K}{\Delta \Gamma_K / 2}$$

Effective Field Theory

Factorize different energy scales:
electroweak, charm, hadronic

- Wilson coefficients
- Short-distance
- Perturbative



$$\langle H^{|\Delta S|=2} \rangle = C \times \langle \tilde{Q}^{|\Delta S|=2} + Q^{|\Delta S|=1} \bar{Q}^{|\Delta S|=1} \rangle$$

- Hadronic matrix elements
- Long-distance
- Non-perturbative QCD - lattice, ChPT

Effective Field Theory: Logs

Generic structure of perturbation series:

|

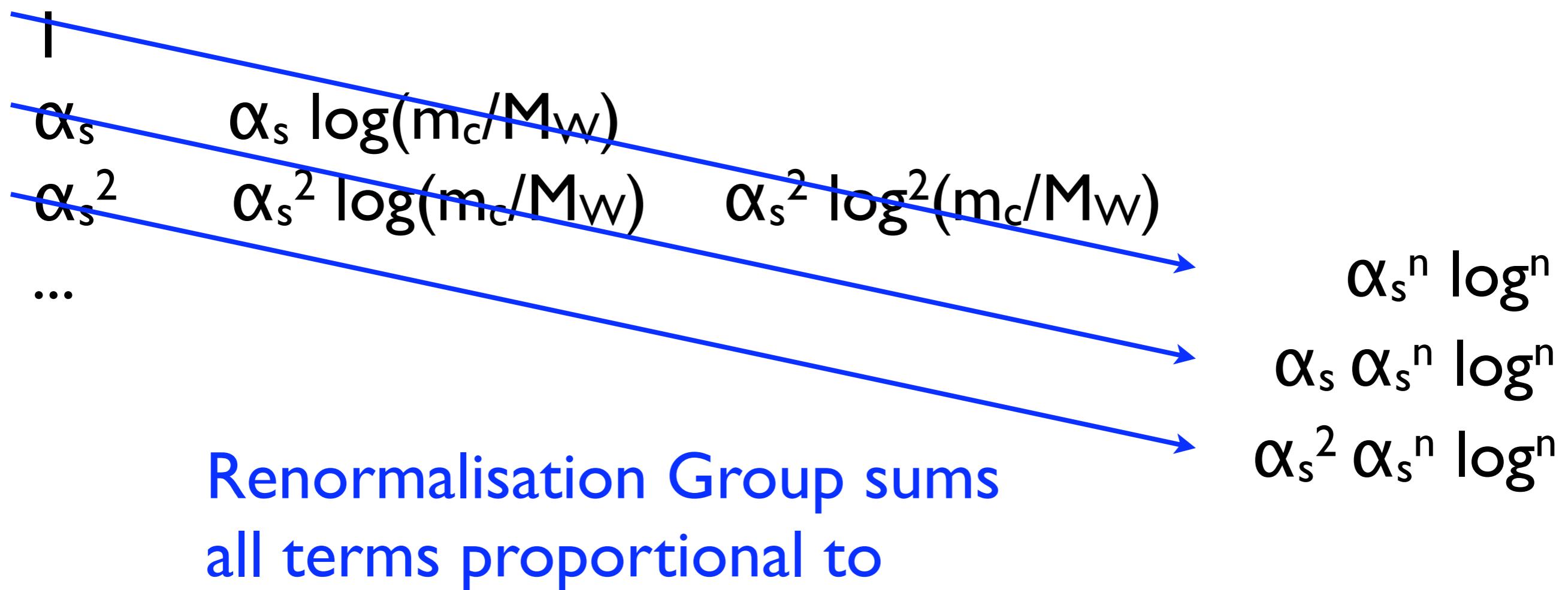
$$\alpha_s \quad \alpha_s \log(m_c/M_w)$$

$$\alpha_s^2 \quad \alpha_s^2 \log(m_c/M_w) \quad \alpha_s^2 \log^2(m_c/M_w)$$

...

Effective Field Theory: Logs

Generic structure of perturbation series:



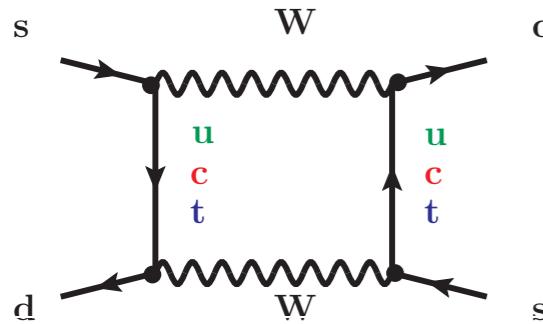
Effective Field Theory: Running

Solution to the Renormalisation Group Equation:
Running Coupling

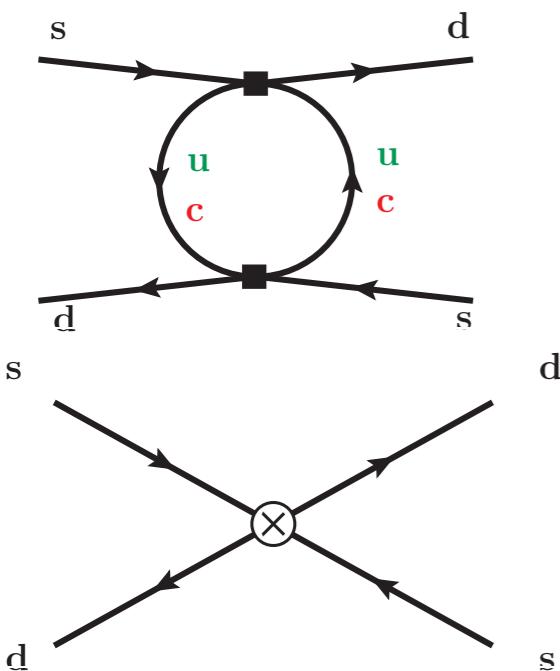
$$C(\mu) = U(\mu, \mu_0)C(\mu_0)$$

$$U(\mu, \mu_0) = T_g \exp \int_{g(\mu_0)}^{g(\mu)} dg' \frac{\gamma^T(g')}{\beta(g')}$$

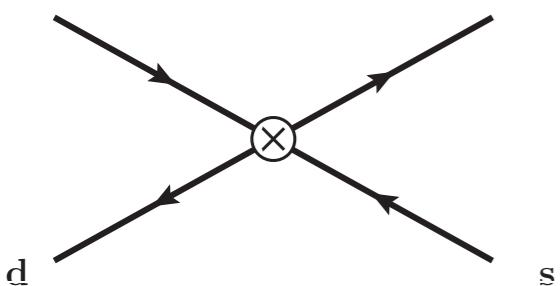
Effective Field Theory



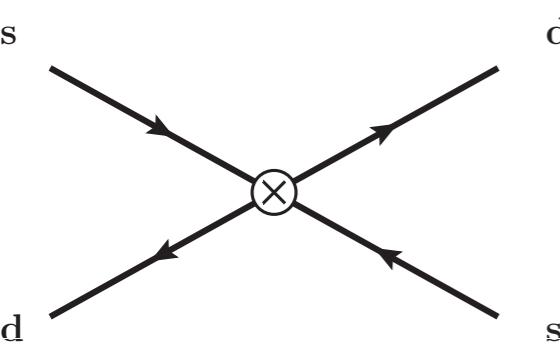
Standard Model



$$\mathcal{H}_{f=5}^{\text{eff}} = \frac{4G_F}{\sqrt{2}} \sum_{i=+,-,3}^6 C_i \left[\sum_{j=+,-} Z_{ij} \sum_{k,l=u,c} V_{ks}^* V_{ld} Q_j^{kl} - \lambda_t \sum_{j=3}^6 Z_{ij} Q_j \right] + \frac{G_F^2}{4\pi^2} \lambda_t^2 \tilde{C}_{S2}^t \tilde{Z}_{S2} \tilde{Q}_{S2}$$



$$+ 8G_F^2 \lambda_c \lambda_t \left[\sum_{k=+,-} \sum_{l=+,-,3}^6 C_k C_l \hat{Z}_{kl,7} + \tilde{C}_7 \tilde{Z}_{77} \right] \tilde{Q}_7 + \text{h.c.} .$$



$$\mathcal{H}_{f=3}^{|\Delta S|=2} = \frac{G_F^2}{4\pi^2} \left[\lambda_c^2 \tilde{C}_{S2}^c(\mu) + \lambda_t^2 \tilde{C}_{S2}^t(\mu) + \lambda_c \lambda_t \tilde{C}_{S2}^{ct}(\mu) \right] \tilde{Z}_{S2} \tilde{Q}_{S2}$$

Definition of B_K

$$B_K = \frac{3}{2} \frac{\langle \bar{K}^0 | Q_{S2} | K^0 \rangle}{f_K^2 M_K^2}$$

Definition of B_K

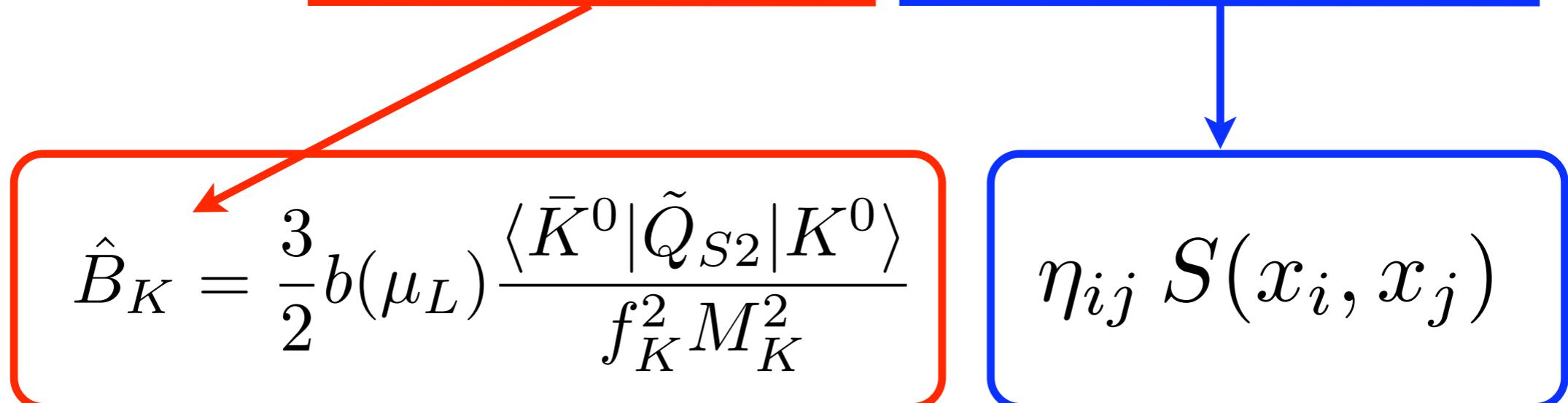
$$B_K = \frac{3}{2} \frac{\langle \bar{K}^0 | Q_{S2} | K^0 \rangle}{f_K^2 M_K^2}$$

$$\langle H_{\text{eff}} \rangle = \langle Q_{S2} (\mu_L) \rangle \quad U(\mu_L, \mu_c) \quad U(\mu_c, \mu_w) \quad C(\mu_w)$$

Definition of B_K

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$$\hat{B}_K = \frac{3}{2} b(\mu_L) \frac{\langle \bar{K}^0 | \tilde{Q}_{S2} | K^0 \rangle}{f_K^2 M_K^2}$$

$$\eta_{ij} S(x_i, x_j)$$

Bare lattice \rightarrow RI-(S)MOM \rightarrow MS-bar [Aoki et. al '08,'10]

Outline

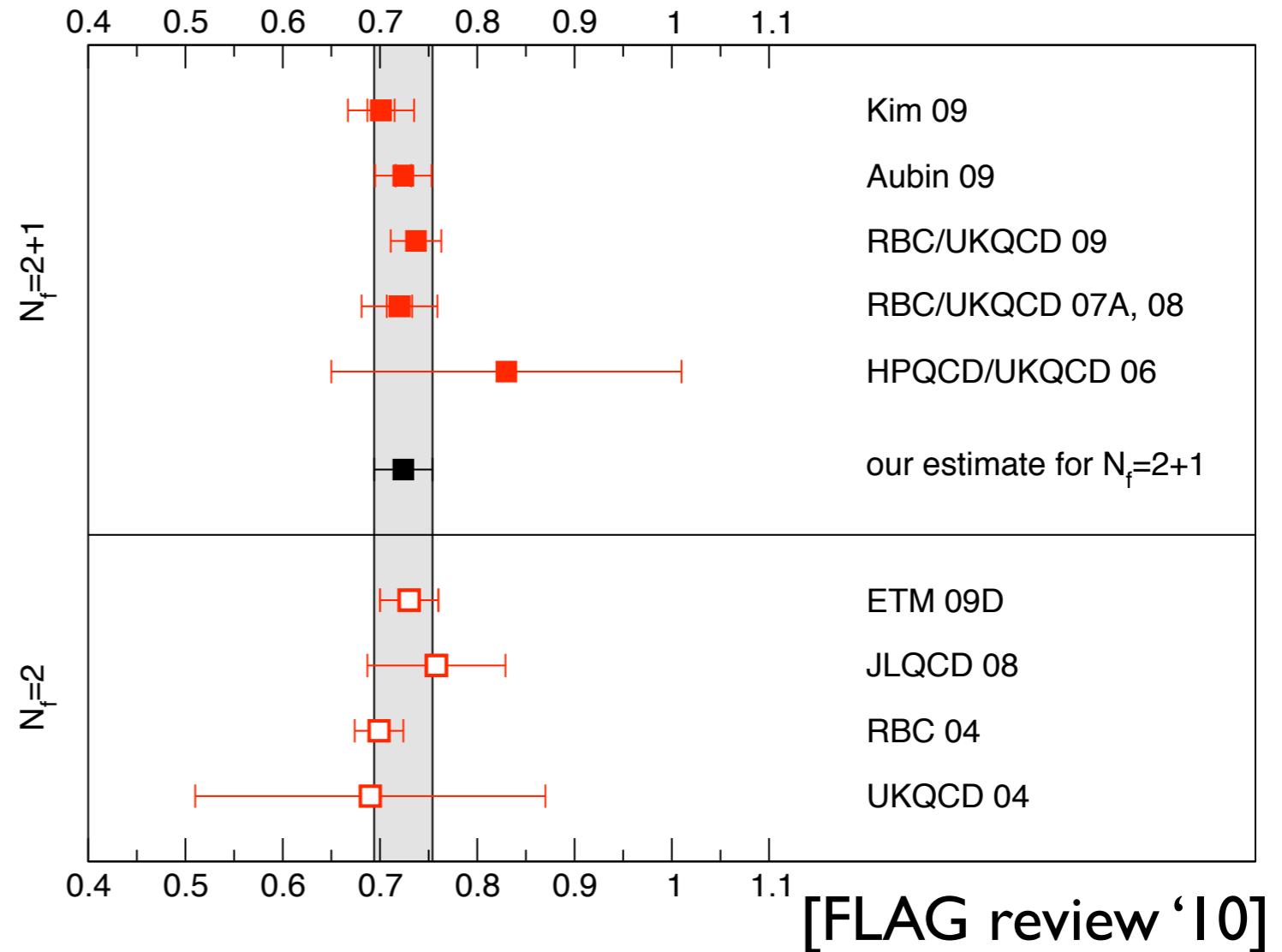
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Determination of B_K

$$\langle K^0 | \tilde{Q}^{|\Delta S|=2} | \bar{K}^0 \rangle \propto \widehat{B}_K = 0.749(27) \text{ [Aoki et al '10]}$$
$$= 0.725(26) \text{ [Aubin et al '09]}$$

Huge progress:

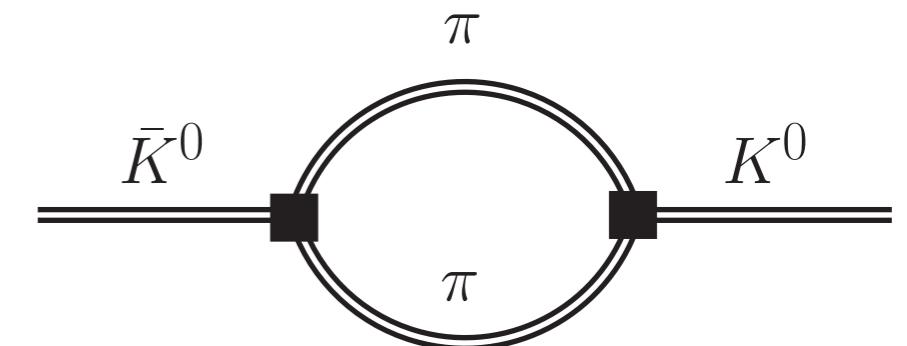
- Unquenched calculations
- Small pion masses
- Different lattice spacings
- Main error: matching to perturbative renormalisation



LD-Contributions from $\Delta S=1$ Hamiltonian

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12})}{\Delta m_K} + \xi \right)$$

Dispersive (real) and
absorptive (imaginary) part of



$$\int d^4x \langle K^0 | H^{|\Delta S|=1}(x) H^{|\Delta S|=1}(0) | \bar{K}^0 \rangle$$

- Estimate ξ from ε'/ε : -6% [Nierste '02, Buras et al. '09]
- Estimate absorptive part in ChPT: +2.4% [Buras et al. '10]

Combine with prefactor to $K_\varepsilon = 0.94(2)$

LD-Contributions from $\Delta S=2$ Hamiltonian

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12})}{\Delta m_K} + \xi \right)$$

$$\langle K^0 | H^{|\Delta S|=2} | \bar{K}^0 \rangle$$

Higher-dimensional
operators contribute
less than 1%

[Cata et al. '03,'04]

Outline

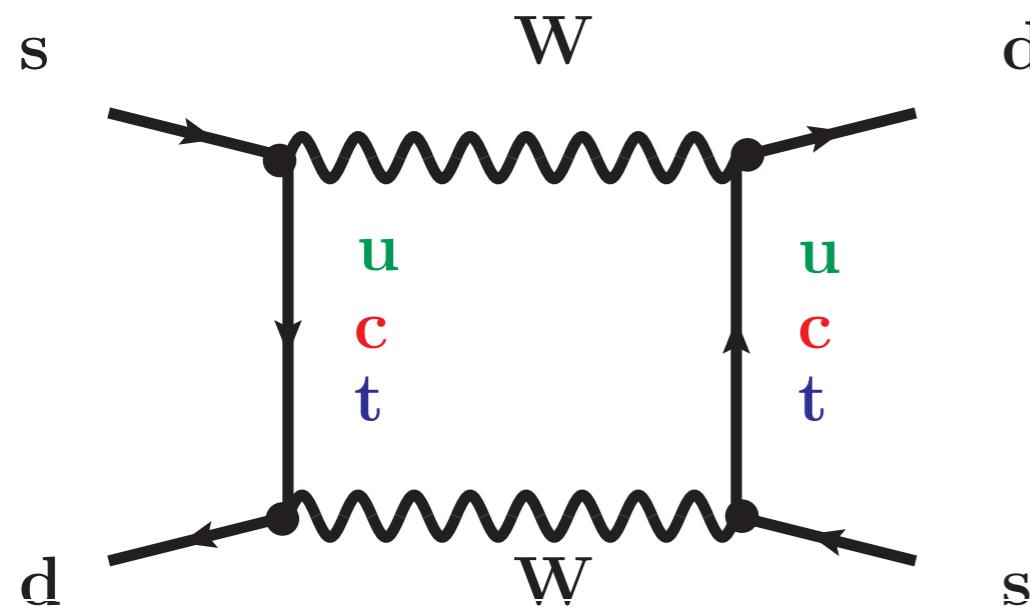
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Too deep we delved there, and woke the nameless fear.

[J. R. R. Tolkien, The Lord of the Rings]

$|\Delta S|=2$ Hamiltonian

$$H^{|\Delta S|=2} \propto \left[\lambda_c^2 \eta_{cc} S\left(\frac{m_c^2}{M_W^2}\right) + \lambda_t^2 \eta_{tt} S\left(\frac{m_t^2}{M_W^2}\right) + \lambda_c \lambda_t \eta_{ct} S\left(\frac{m_c^2}{M_W^2}, \frac{m_t^2}{M_W^2}\right) \right] \tilde{Q}^{|\Delta S|=2}$$



CKM parameters:

$$\lambda_i = V_{is}^* V_{id}$$

$$\lambda_u = -\lambda_c - \lambda_t$$

$$\text{Re } \langle K^0 | H^{|\Delta S|=2} | \bar{K}^0 \rangle \Rightarrow \Delta m_K \quad \text{Im } \langle K^0 | H^{|\Delta S|=2} | \bar{K}^0 \rangle \Rightarrow \epsilon_K$$

$|\Delta S|=2$ Hamiltonian

$$H^{|\Delta S|=2} \propto \left[\lambda_c^2 \eta_{cc} S\left(\frac{m_c^2}{M_W^2}\right) + \lambda_t^2 \eta_{tt} S\left(\frac{m_t^2}{M_W^2}\right) + \lambda_c \lambda_t \eta_{ct} S\left(\frac{m_c^2}{M_W^2}, \frac{m_t^2}{M_W^2}\right) \right] \tilde{Q}^{|\Delta S|=2}$$

$$\approx \lambda_{\text{Cabibbo}}^{10} \times \eta_{tt} \times \frac{m_t^2}{M_W^2}$$

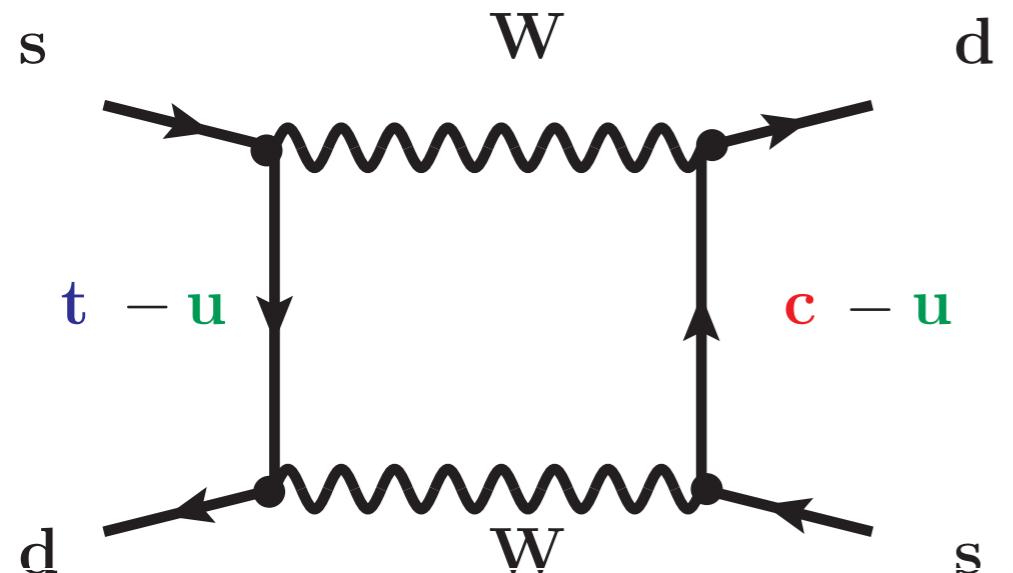
$\lambda_{\text{Cabibbo}} \approx 0.2$

Dominant contribution to ε_K ($\approx +74\%$)

$\eta_{tt} = 0.5765(65)$ at NLO QCD [Buras et al. '90]

$|\Delta S|=2$ Hamiltonian

$$H^{|\Delta S|=2} \propto \left[\lambda_c^2 \eta_{cc} S\left(\frac{m_c^2}{M_W^2}\right) + \lambda_t^2 \eta_{tt} S\left(\frac{m_t^2}{M_W^2}\right) + \lambda_c \lambda_t \eta_{ct} S\left(\frac{m_c^2}{M_W^2}, \frac{m_t^2}{M_W^2}\right) \right] \tilde{Q}^{|\Delta S|=2}$$



$$\approx \lambda_{\text{Cabibbo}}^6 \times \eta_{ct} \times \frac{m_c^2}{M_W^2} \log \frac{m_c^2}{M_W^2}$$

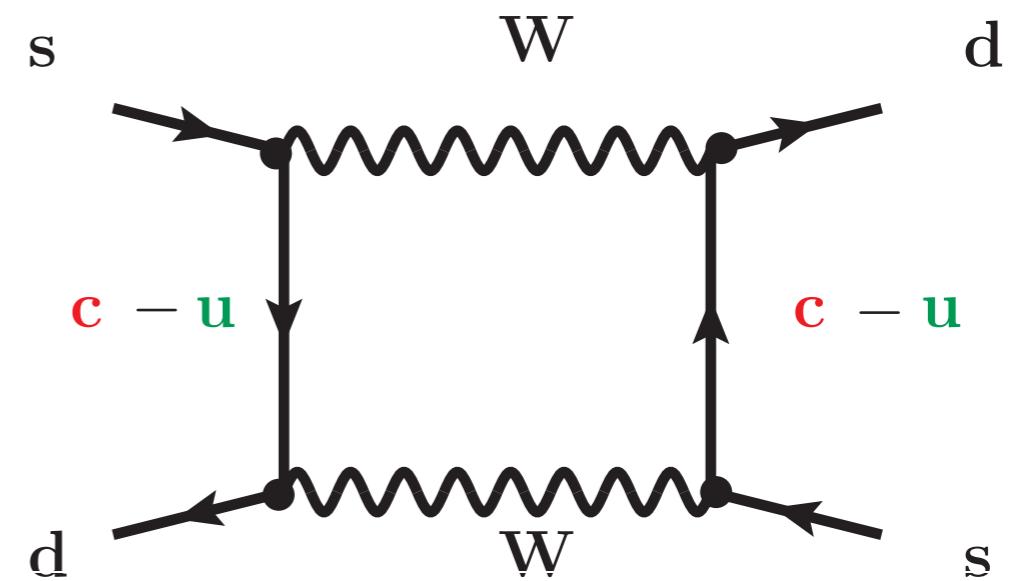
Contributes $\approx +40\%$ to ε_K (NLO)

$\eta_{ct} = 0.47(4)$ at NLO QCD [Herrlich, Nierste '96]

$|\Delta S|=2$ Hamiltonian

$$H^{|\Delta S|=2} \propto \left[\lambda_c^2 \eta_{cc} S\left(\frac{m_c^2}{M_W^2}\right) + \lambda_t^2 \eta_{tt} S\left(\frac{m_t^2}{M_W^2}\right) + \lambda_c \lambda_t \eta_{ct} S\left(\frac{m_c^2}{M_W^2}, \frac{m_t^2}{M_W^2}\right) \right] \tilde{Q}^{|\Delta S|=2}$$

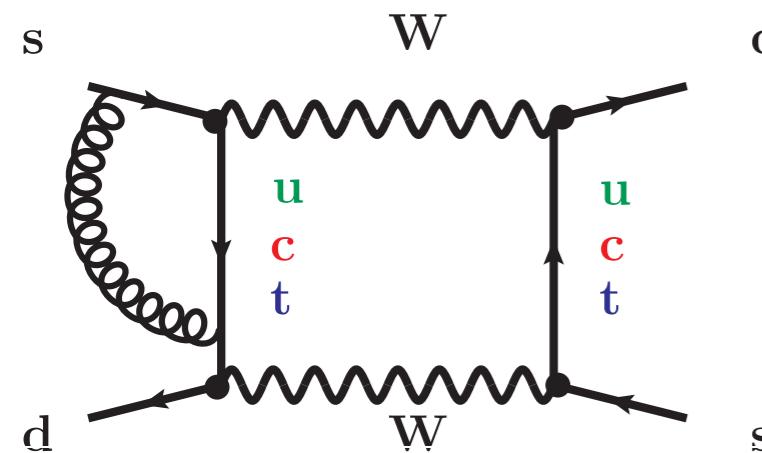
$$\approx \lambda_{\text{Cabibbo}}^6 \times \eta_{cc} \times \frac{m_c^2}{M_W^2}$$



Smallest contribution to ε_K ($\approx -14\%$ at NLO)

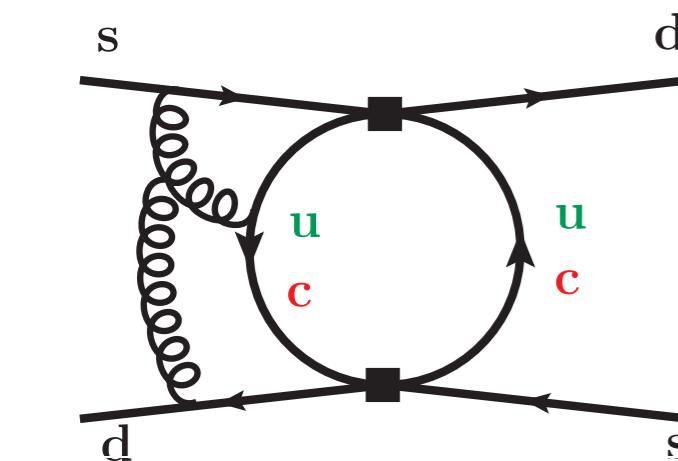
$\eta_{cc} = 1.40(35)$ at NLO QCD [Herrlich, Nierste '94, '03]

η_{ct} at NNLO: Calculation



Matching at M_W : initial condition

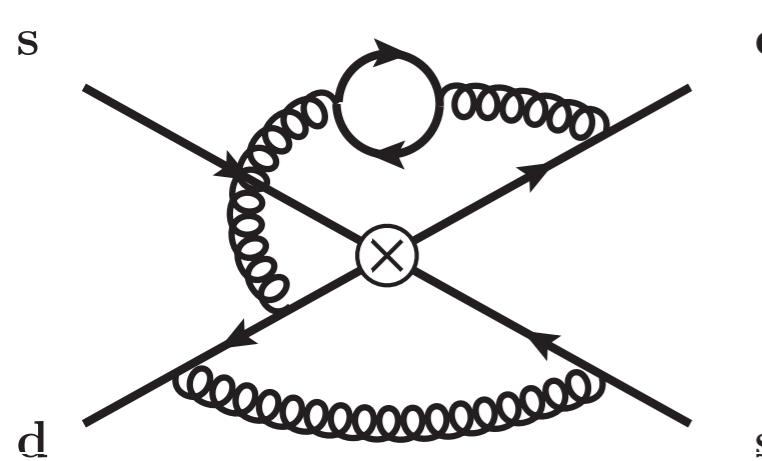
μ_{NP}



Running to m_c

μ_b

- $O(100\ 000)$ diagrams
- RGE for double insertions
- Include threshold corrections at m_b



Matching at m_c
to effective 3-flavour theory

μ_c

Λ_{QCD}

Computation of the diagrams

We had to compute $O(100\,000)$ one-, two-, and three-loop vacuum Feynman diagrams. We use

- **qgraf** [Nogueira '93], **MATAD** [Steinhauser '01]
- Our own **FORM** [Vermaseren '00] routines
- **FeynArts** [Hahn '01], **FIRE** [Smirnov '08], **Mathematica**

Two fully independent setups, renormalisation, numerics

η_{ct} at NNLO: Result...

$$\Rightarrow \eta_{ct} = 0.496$$

Corresponds to +7% shift!

η_{ct} at NNLO: Result...

$$\Rightarrow \eta_{\text{ct}} = 0.496 \pm ???$$

Corresponds to +7% shift!

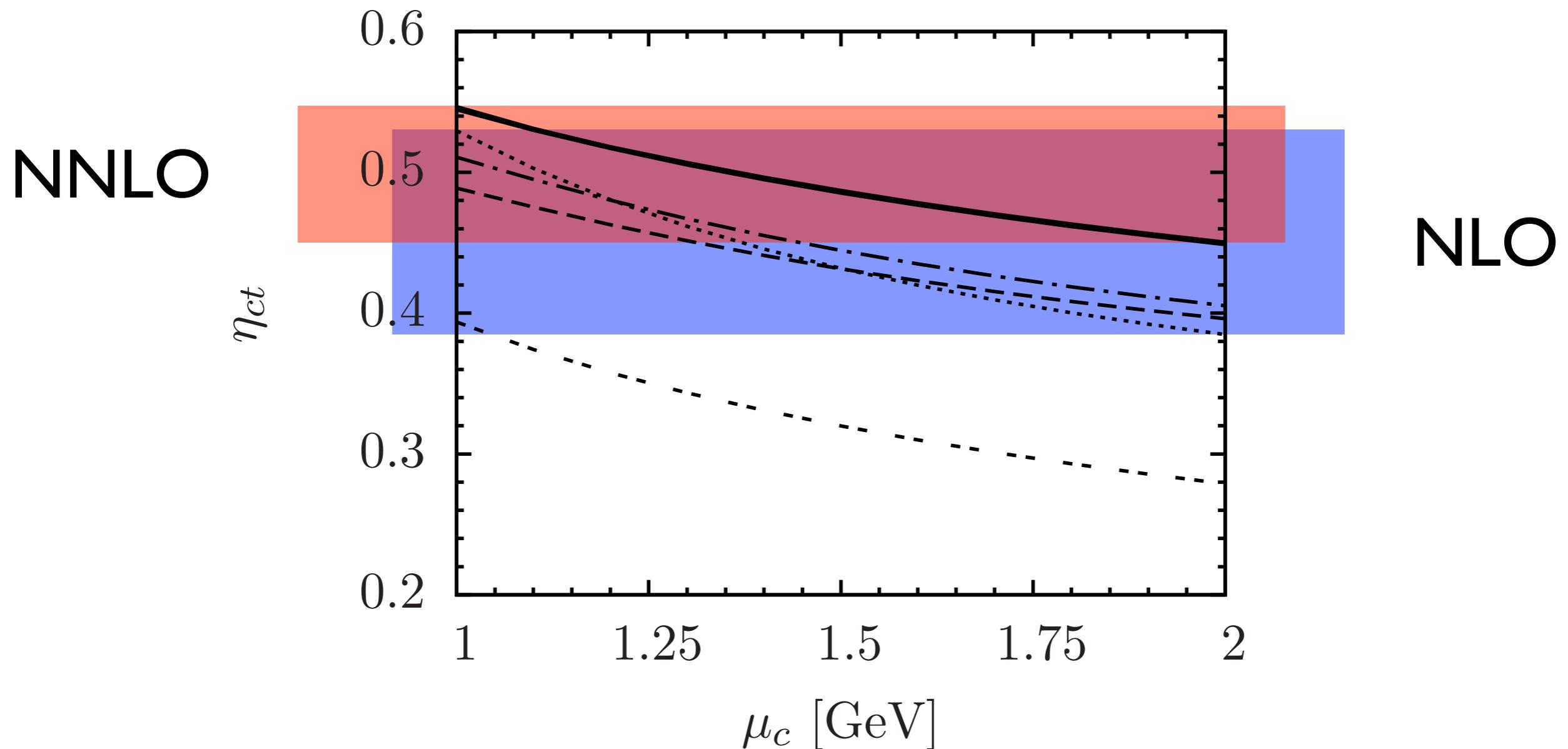
Theory uncertainty

Consequence of the truncation
of the perturbative expansion

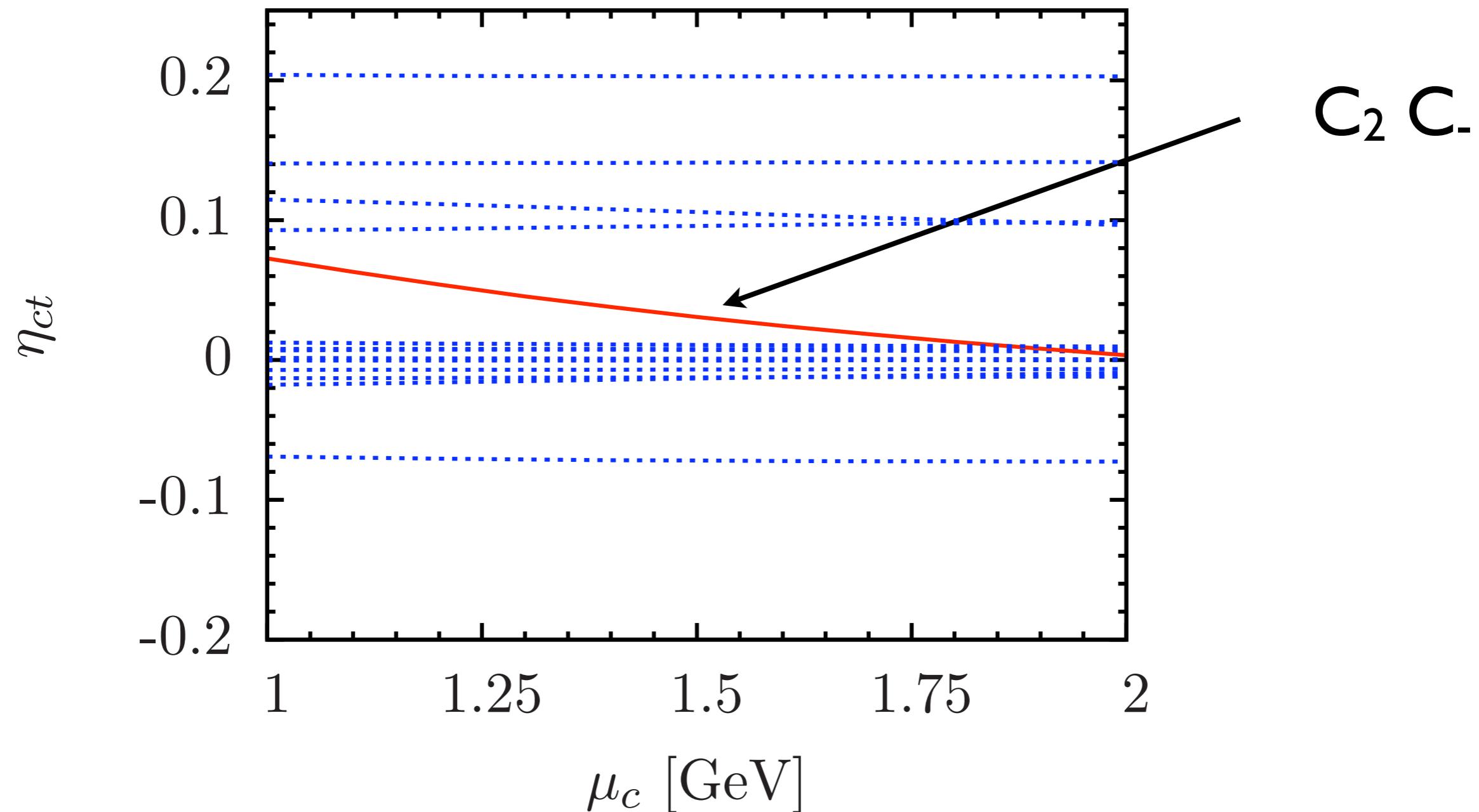
η_{ct} (and η_{cc}) is scale- and scheme independent!

- Study residual scale dependence
- Size of NNLO corrections

η_{ct} at NNLO: Scale Dependence



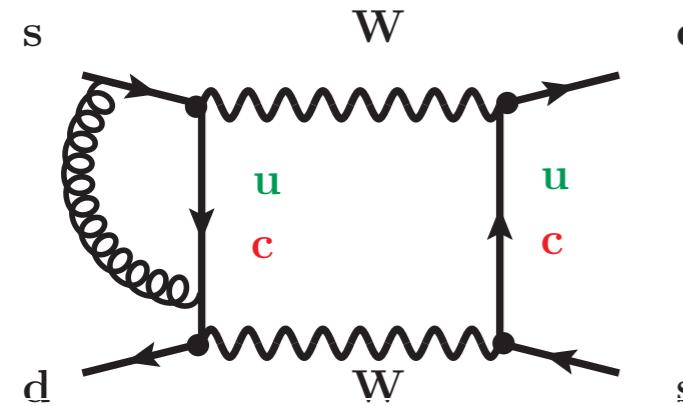
η_{ct} at NNLO: Scale Dependence



η_{ct} at NNLO: Result

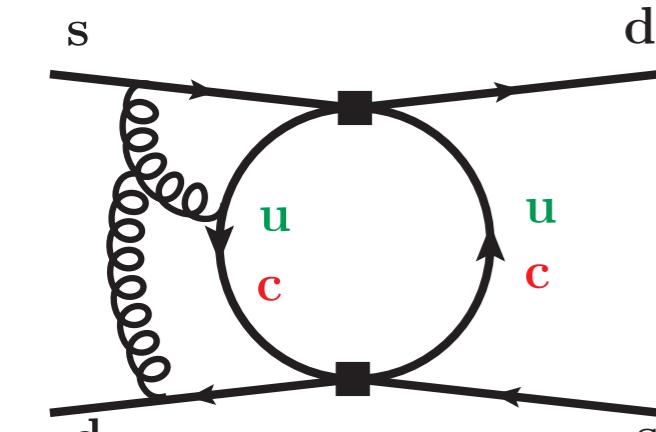
$$\Rightarrow \eta_{\text{ct}} = 0.496 \pm 0.047$$

η_{cc} at NNLO: Calculation



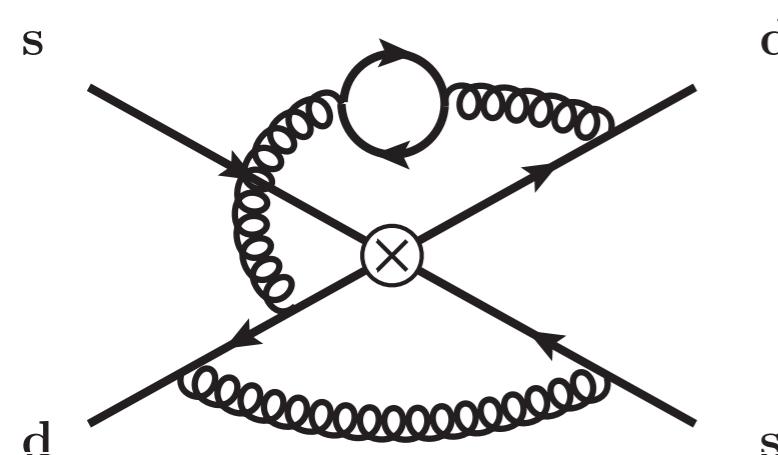
Matching at M_W : vanishes by GIM
[Witten '76]

μ_{NP}



Running to m_c :
only $\Delta S=1$ operators contribute

μ_b



Matching at m_c :

- Match to effective 3-flavour theory
- $O(100\ 000)$ 3-loop Feynman diagrams

μ_c

Λ_{QCD}

η_{cc} at NNLO: Result...

$$\begin{aligned}\langle Q_2 Q_2 \rangle = & \frac{69738523}{113400} + \frac{47407}{8505} \pi^2 - \frac{1733}{810} \pi^4 + \frac{1872}{35} \sqrt{3} \operatorname{Im} \operatorname{Li}_2((-1)^{1/3}) \\ & + 24 (\operatorname{Im} \operatorname{Li}_2((-1)^{1/3}))^2 - \frac{32}{27} \pi^2 \log(2)^2 + \frac{32}{27} \log(2)^4 + \frac{563}{18} \log \frac{\mu_c^2}{m_c^2} \\ & + \frac{32}{3} \pi^2 \log \frac{\mu_c^2}{m_c^2} + \frac{193}{3} \log^2 \frac{\mu_c^2}{m_c^2} + \frac{256}{9} \operatorname{Li}_4(1/2) - \frac{15145}{54} \zeta_3\end{aligned}$$

η_{cc} at NNLO: Result...

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$$\Rightarrow \eta_{cc} = 1.87$$

Corresponds to +36% shift!

η_{cc} at NNLO: Result...

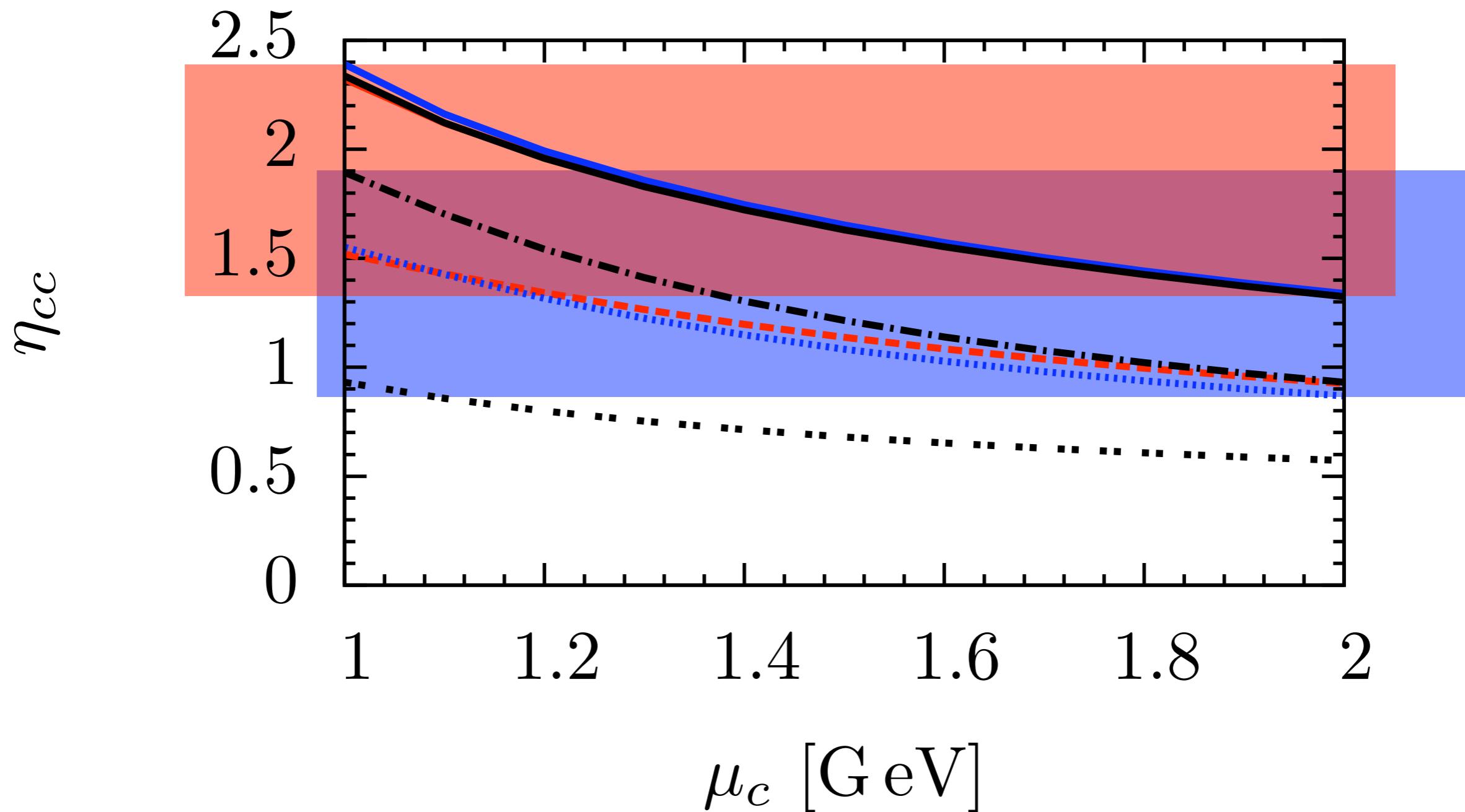
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$$\Rightarrow \eta_{cc} = 1.87 \pm ???$$

Corresponds to +36% shift!

η_{cc} at NNLO: Scale Dependence

Matching @ $mc(mc)$



η_{cc} at NNLO: Convergence

$$\begin{aligned}\eta_{cc}/\alpha_s^{(2/9)} = & 1 \\ & + \alpha_s (0.25 + 0.64 L_c) \\ & + \alpha_s^2 (1.20 + 0.44 L_c + 1.08 L_c^2)\end{aligned}$$

$$L_c = \log(m_c/M_W) = -4.14 \quad \alpha_s = \alpha_s(m_c) = 0.35$$

$$\begin{aligned}\eta_{cc}/\alpha_s^{(2/9)} = & 1 \\ & + (0.09 + 0.9) \\ & + (0.15 - 0.2 + 2.3)\end{aligned}$$

η_{cc} at NNLO: Uncertainty

- Divergence of perturbation series?
- Make use of RI-SMOM schemes, at NNLO?
- Compute charm on the lattice???

Preliminary prescription:
Add NNLO shift and scale uncertainty in quadrature

$$\eta_{cc} = 1.87 \pm 0.76$$

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„*Data! Data! Data!*“ he cried impatiently.

[Sherlock Holmes, The adventure of the copper beaches]

$|\varepsilon_K|$ - Numerics

using only experimental and lattice input:

$$|\varepsilon_K| \propto K_\varepsilon B_K |V_{cb}|^2 \xi_s \sin\beta$$

$$\times (|V_{cb}|^2 \xi_s \cos\beta \eta_{tt} S(x_t) + \eta_{ct} S(x_c, x_t) - \eta_{cc} S(x_c))$$

$ V_{cb} $	0.0406(13)
$\sin(2\beta)$	0.671(23)
ξ_s	1.243(28)
B_K	0.725(26)
K_ε	0.94(2)

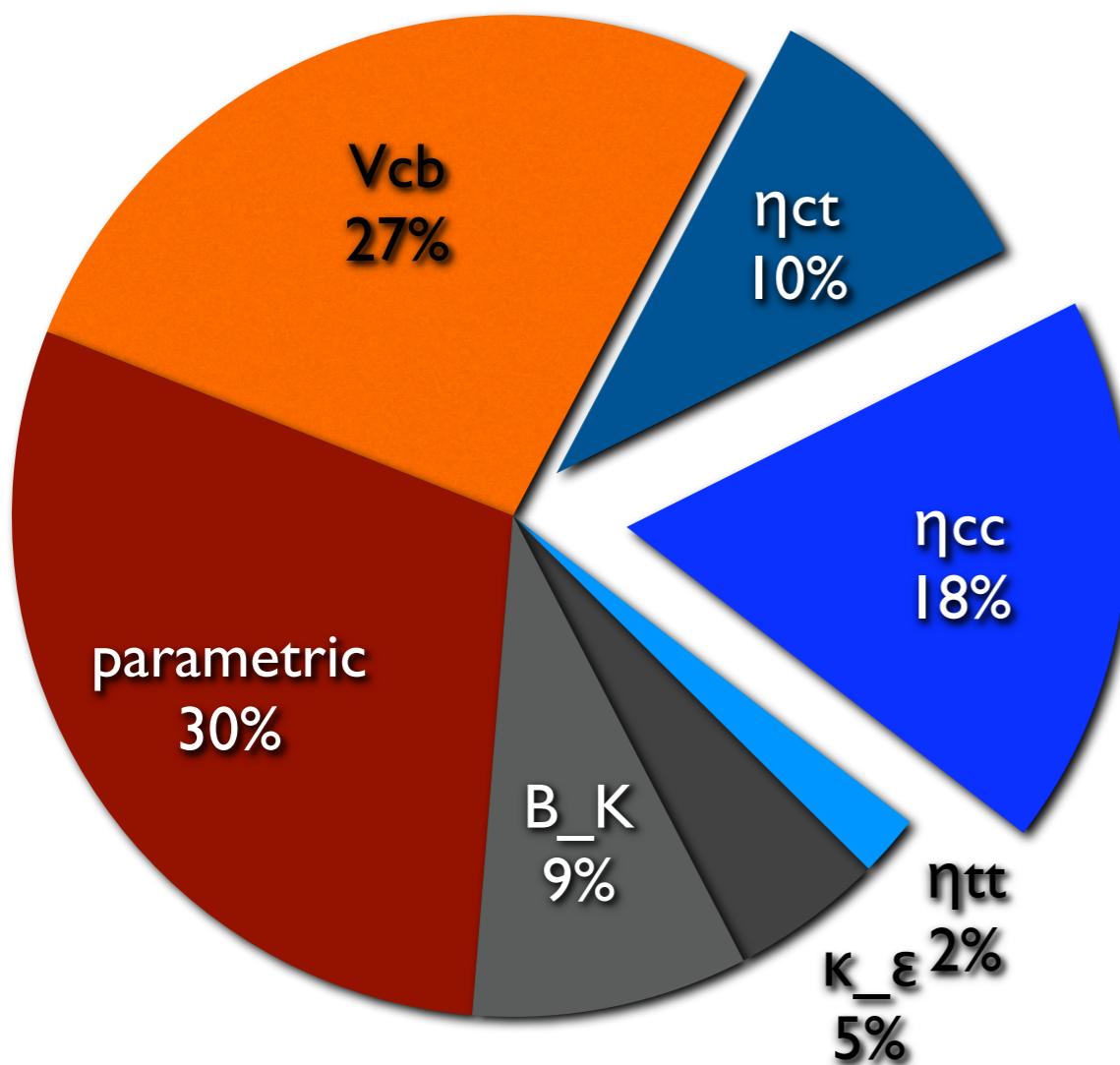
η_{ct}	0.496(47)
η_{cc}	1.87(76)
η_{tt}	0.5765(65)

$$\xi_s = \frac{F_{B_s} \sqrt{\hat{B}_s}}{F_{B_d} \sqrt{\hat{B}_d}} \quad x_q = \frac{m_q^2}{M_W^2}$$

$|\varepsilon_K|$ - Result & Error Budget

$$|\varepsilon_K| = 1.81(28) \times 10^{-3}$$

using $\eta_{cc} = 1.87(76)$
and $\eta_{ct} = 0.496(47)$



cf. $|\varepsilon_K| = 1.83(27) \times 10^{-3}$ (NLO)

cf. $|\varepsilon_K| = 1.53(23) \times 10^{-3}$
using $|V_{cb}| = 0.0387(11)$ (exclusive)

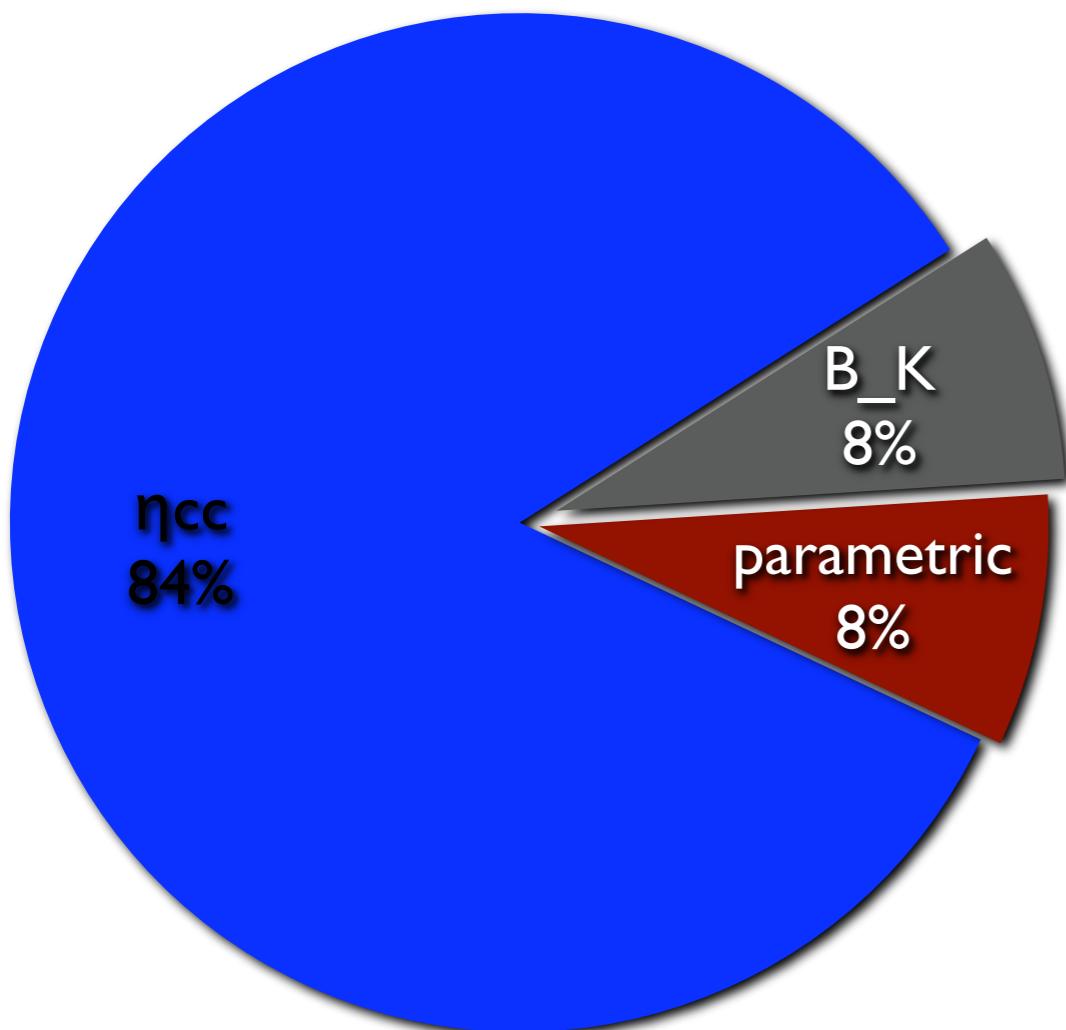
Experiment [PDG '10]:

$$|\varepsilon_K|^{\text{exp.}} \equiv 2.228(11) \times 10^{-3}$$

$\Delta M_K(\text{SD})$ -- Numerics

$$\Delta M_K(\text{SD}) = (3.0 \pm 1.2) \times 10^{-15} \text{ GeV} \text{ using } \eta_{cc} = 1.87(76)$$

cf. $\Delta M_K(\text{SD}) = (2.2 \pm 0.9) \times 10^{-15} \text{ GeV}$ using $\eta_{cc} = 1.38(53)$ (NLO!)



$$\Delta M_K(\text{SD}) = 3.483(6) \times 10^{-15} \text{ GeV}$$

[PDG '10]

$(86 \pm 34)\%$ short distance

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In five minutes you will say that it is all so absurdly simple.

[Sherlock Holmes, The adventure of the dancing men]

- ε_K is a powerful constraint on models of New Physics.
- Progress in non-perturbative calculations
(in particular B_K)
- NNLO calculation yields **+7% shift** of charm-top-quark contribution to ε_K , leading to $\eta_{ct} = 0.496(47)$
- NNLO calculation yields **+36% shift** of charm-quark contribution to ε_K , leading to $\eta_{cc} = 1.87(76)$
- SM prediction for ε_K is shifted by **$\approx -1\%$** to $\varepsilon_K = 1.81(28) \times 10^{-3}$.

We might already see new physics --
but first we need to understand the Standard Model!